

# Impact of the stellar evolution adopting m-CAK self-consistent winds

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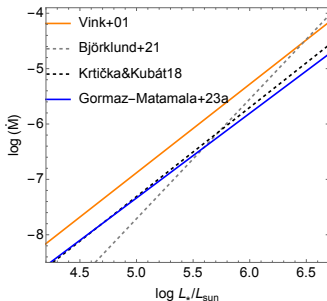


# Part I

## Self-consistent evolution models

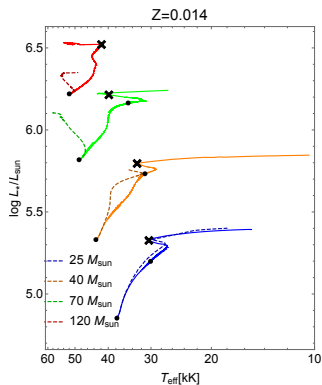
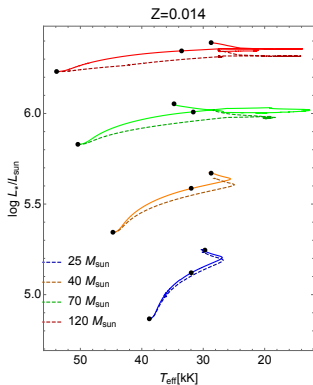
# $\dot{M}$ from self-consistent line-driven winds (Gormaz-Matamala et al. 2019,2021,2022a)

- New self-consistent  $\dot{M}$  based on the m-CAK theory (see talks from Curé, Araya and Figueroa-Tapia) are  $\sim 3$  lower than standard V01 recipe.
- Beyond m-CAK theory, wind hydrodynamics can be analytically solved by using the Lambert W-function.
- Wind solutions are unique, and thus provide a self-consistent value for the mass-loss rate  $\dot{M}$ .



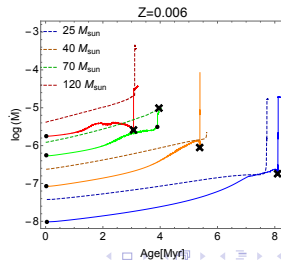
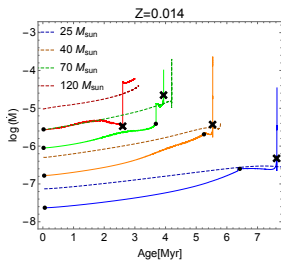
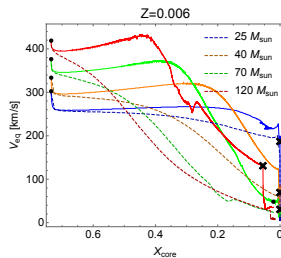
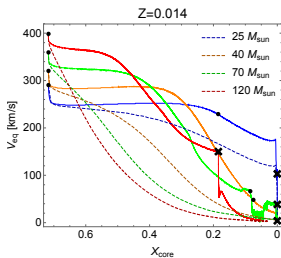
## New vs old evolution models (Gormaz-Matamala et al. 2022b,2023)

- New stellar evolution models adopting lower  $\dot{M}$  retain more mass and are brighter than old evolution models.
- Weaker winds for evolution also affect the rotation and the mixing of the stellar structures.



Old models from Ekström+2012 (dashed lines)

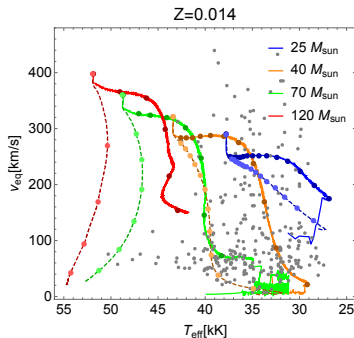
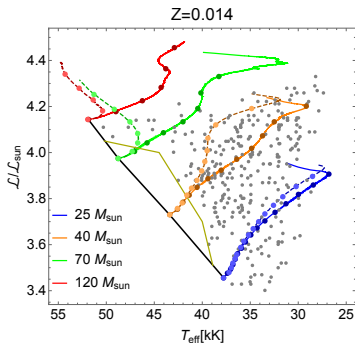
## Evolution of equatorial rotation



## Properties of the stellar models at the end of the main sequence.

$M_{\text{zams}}$	$Z$	$v_{\text{rot,ini}}$	mass loss	End of Main Sequence							
				$t_{\text{MS}}$	$M_*$	$R_*$	$v_{\text{rot,surf}}$	$Y_{\text{surf}}$	N/C	N/O	
70.0	0.014	360.0	$\dot{M}_{\text{sc}}$	3.9	60.4	26.9	3.9	0.63	18.2	5.6	
			$\dot{M}_{\text{Vink}}$	4.2	42.7	18.7	7.1	0.89	73.3	69.5	
70.0	0.006	376.0	$\dot{M}_{\text{sc}}$	4.0	64.6	21.7	26.2	0.51	8.0	2.4	
			$\dot{M}_{\text{Vink}}$	3.9	57.3	15.5	37.0	0.69	34.0	13.4	
40.0	0.014	321.0	$\dot{M}_{\text{sc}}$	5.5	36.4	23.6	37.4	0.46	5.5	1.5	
			$\dot{M}_{\text{Vink}}$	5.8	31.9	19.2	9.9	0.55	13.9	3.9	
40.0	0.006	334.0	$\dot{M}_{\text{sc}}$	5.4	38.2	18.5	71.7	0.37	3.7	1.0	
			$\dot{M}_{\text{Vink}}$	5.6	36.2	16.2	30.6	0.44	6.1	1.6	
25.0	0.014	291.0	$\dot{M}_{\text{sc}}$	7.6	24.3	16.0	113.0	0.32	2.6	0.6	
			$\dot{M}_{\text{Vink}}$	8.0	23.6	15.9	68.4	0.34	3.3	0.8	
25.0	0.006	302.0	$\dot{M}_{\text{sc}}$	8.1	24.6	16.7	185.0	0.33	4.3	0.8	
			$\dot{M}_{\text{Vink}}$	7.7	24.3	14.1	159.0	0.31	3.4	0.7	

## Comparison with rotational surveys from Holgado+22



Self-consistent evolution models (solid lines) explain better the lack of fast rotators ( $v_{\text{rot}} \gtrsim 150 \text{ km s}^{-1}$ ) at  $T_{\text{eff}} \gtrsim 42.5 \text{ kK}$  and the abundance of stars with  $v_{\text{rot}} \lesssim 150 \text{ km s}^{-1}$  and  $T_{\text{eff}}$  between  $\sim 27$  and  $40 \text{ kK}$ .

# Highlights

- Self-consistent stellar winds gives lower  $\dot{M}$  than the original values, thus generating important changes in the stellar evolution of massive stars.

$$\dot{M} \propto Z^{0.4 + \frac{15.75}{M_*/M_\odot}},$$

compare with Gräfener&Hamann08 for WNL and Krtička&Kubát24 for OB-stars.

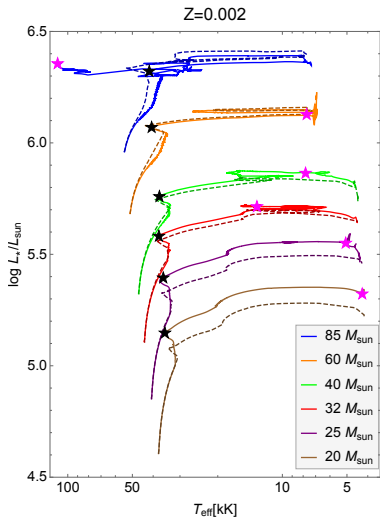
- Rotational velocities predicted by the self-consistent evolution models are in agreement with the most recent observational surveys.
- Implications of new evolution models are wide, and can be extended for lower metallicities and more massive stars.



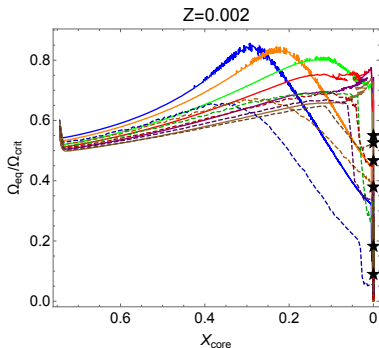
# Part II

## Implications: SMC massive stars

## Massive stars at SMC metallicity (Gormaz-Matamala et al. 2024)



Lower  $\dot{M}$ , but larger  $v_{\text{rot}}$  (and thus more efficient rotational mixing).

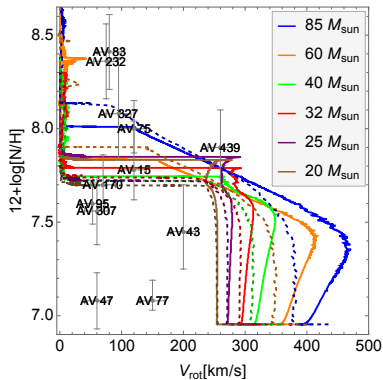
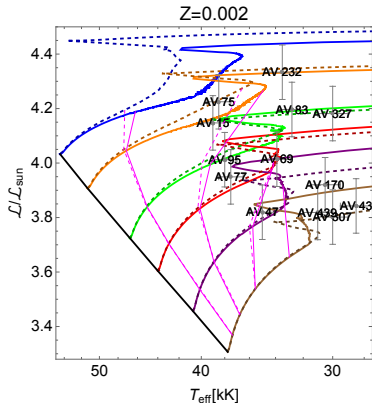


## Rotation for massive stars at SMC metallicity

$M_{\text{zams}}$ [ $M_{\odot}$ ]	wind recipe	H-core burning							He-core burning					
		$\langle v_{\text{rot}} \rangle$ [ $\text{km s}^{-1}$ ]	$v_{\text{rot, final}}$	$\tau_{\text{MS}}$ [Myr]	$M_{\text{fin}}$ [ $M_{\odot}$ ]	$Y_{\text{surf}}$	N/C mass fraction	N/O	$\tau_{\text{He}}$ [Myr]	$M_{\text{fin}}$ [ $M_{\odot}$ ]	$Y_{\text{surf, fin}}$ mass fraction	N/C mass fraction	N/O	$R_{\text{max}}$ [ $R_{\odot}$ ]
85	$\dot{M}_{\text{Sc}}$	323.	87.	3.6	80.1	0.47	5.7	2.3	0.33	47.4	0.60	0.3	0.1	831.0
85	$\dot{M}_{\text{Vink}}$	319.	26.	3.6	77.2	0.54	9.8	3.8	0.31	51.36	0.83	132.2	89.0	882.4
60	$\dot{M}_{\text{Sc}}$	355.	111.	4.2	57.3	0.38	3.5	1.3	0.36	33.14	0.81	65.5	91.3	910.7
60	$\dot{M}_{\text{Vink}}$	317.	85.	4.2	56.3	0.41	4.7	1.6	0.35	39.45	0.74	64.0	33.8	744.8
40	$\dot{M}_{\text{Sc}}$	323.	210.	5.4	39.1	0.33	3.2	1.0	0.46	30.24	0.63	27.0	8.4	1280.7
40	$\dot{M}_{\text{Vink}}$	300.	125.	5.3	38.7	0.31	3.2	0.9	0.45	29.16	0.57	20.0	5.0	1288.3
32	$\dot{M}_{\text{Sc}}$	299.	258.	6.6	31.5	0.34	4.5	1.1	0.55	24.53	0.61	31.6	7.7	1106.1
32	$\dot{M}_{\text{Vink}}$	282.	177.	6.4	31.2	0.31	4.1	0.9	0.54	24.68	0.58	28.2	5.5	1103.6
25	$\dot{M}_{\text{Sc}}$	273.	285.	8.6	24.8	0.36	7.4	1.2	0.66	20.04	0.44	12.0	1.7	984.2
25	$\dot{M}_{\text{Vink}}$	262.	132.	7.7	24.7	0.29	4.7	0.9	0.69	22.77	0.32	6.4	1.0	987.3
20	$\dot{M}_{\text{Sc}}$	252.	277.	10.5	19.9	0.34	8.8	1.1	0.87	18.88	0.36	10.6	1.3	899.4
20	$\dot{M}_{\text{Vink}}$	248.	194.	9.3	19.9	0.28	4.7	0.8	0.95	19.25	0.23	5.7	0.9	793.0

Applications for studies analysing rotation and chemical enrichment in SMC  
(Gómez-González et al. in prep.)

## SMC O-type giants and supergiants from Bouret+21



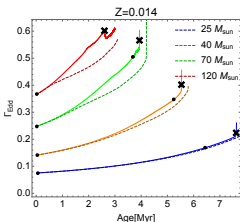
Analyses of theoretical chemical enrichment must be extended for BSG and RSG (see talks from Venero, Sánchez-Arias, and poster from González-Torà).

## Part III

# Implications for stellar evolution close to the Eddington limit

# Stellar winds in the proximity to the Eddington limit

- Very massive stars: born with  $M_{\text{zams}} \gtrsim 100 M_{\odot}$ , with large convective cores.
- Strong stellar winds make VMS become WR (WNh) stars during their main sequence phase (Crowther+10).
- The  $\dot{M} \sim \Gamma_e$  relationship changes its slope over a certain  $\Gamma_{e,\text{trans}}$  (Vink+11; Vink&Gräfenner12).
- "Proximity" to Eddington limit?  $\Gamma_e \gtrsim 0.3$ ?  
 $\Gamma_e \gtrsim 0.7$ ?



## Sabhanit+23

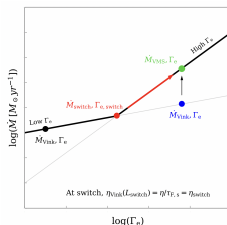


Figure 4. Schematic representation of the new VMS mass-loss framework where the condition to switch to a high- $\Gamma_e$ , optically-thick wind is based on wind efficiency parameters. A step-by-step method is described in the text.

**Table 1.** For given CAK force multiplier  $\alpha$  we list expected transition Eddington parameters and mass-loss rate dependence for  $\Gamma_e \ll 1$  and  $\Gamma_e \rightarrow 1$ .

$\alpha$	$\Gamma_{e, \text{trans}}$	$\dot{M} \propto \Gamma_e^{-1/\alpha+1/2}$	$\dot{M} \propto 1/(1 - \Gamma_e)^{(1-\alpha)/\alpha+2}$
0.3	0.479	$\Gamma_e^{-3.83}$	$(1 - \Gamma_e)^{-4.3}$
0.4	0.473	$\Gamma_e^{-3.0}$	$(1 - \Gamma_e)^{-3.5}$
0.5	0.468	$\Gamma_e^{-2.5}$	$(1 - \Gamma_e)^{-3.0}$
0.6	0.464	$\Gamma_e^{-2.17}$	$(1 - \Gamma_e)^{-2.7}$

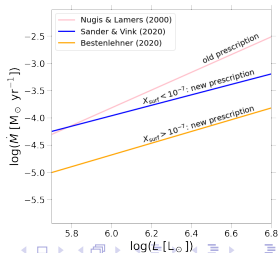
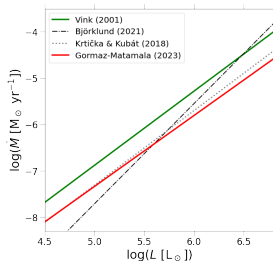
## Bestenlehner20



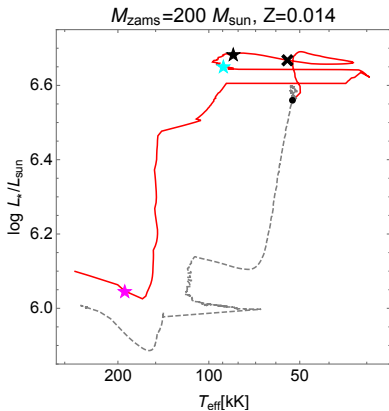
# New mass loss prescription for thin and thick winds

- Optically thin winds ( $X_{\text{surf}} \geq 0.3$  &  $\Gamma_e \leq 0.5$ )  
 ⇒ Gormaz-Matamala et al. (2023)
- H-rich optically thick winds ( $\Gamma_e > 0.5$  &  $X_{\text{surf}} \geq 10^{-7}$ )  
 ⇒ Bestenlehner (2020)
- H-poor optically thick winds ( $X_{\text{surf}} < 10^{-7}$ )  
 ⇒ Sander & Vink (2020)

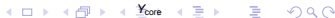
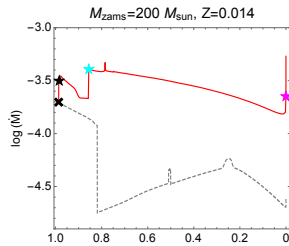
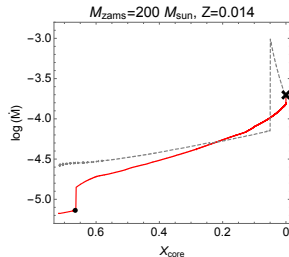
Full scheme in Romagnolo et al. (2024).



# Evolution for $M_{\text{zams}} = 200 M_{\odot}$ , $\Omega = 0.4$ , $\alpha_{\text{ov}} = 0.1$ .

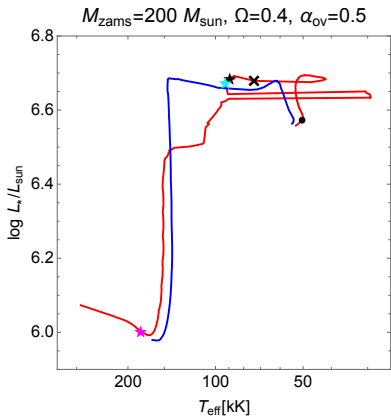


Grey dashed model from Yusof+13.

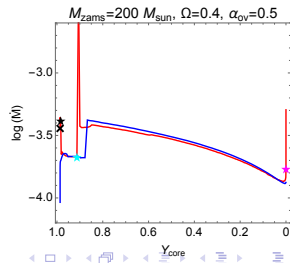
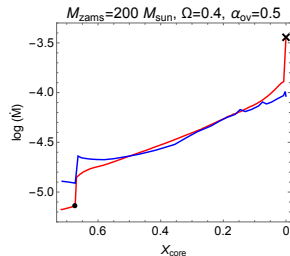




$\alpha_{OV} = 0.5$ , Ledoux criterion, and Tayler-Spruit dynamo.



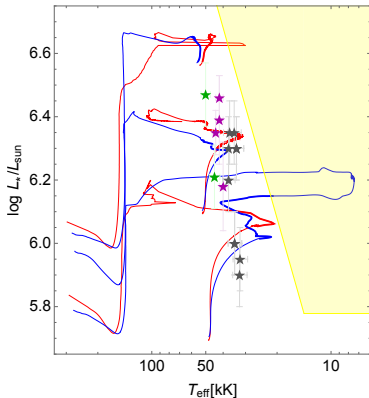
GENEC (red) vs Mesa (blue) models



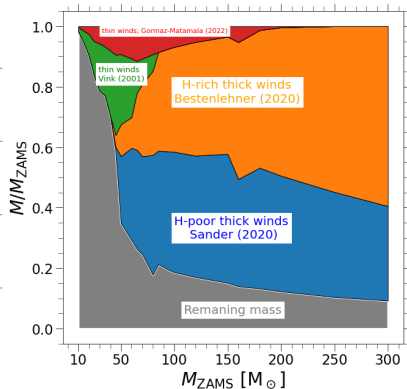
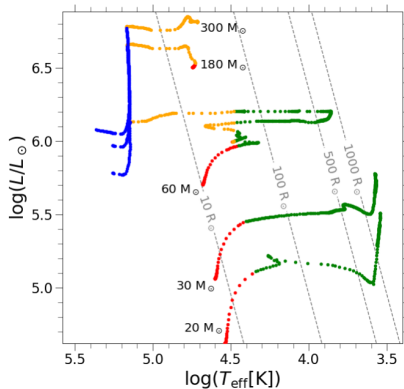
Existence of WNh evolved from  $M_{\text{zams}} = 60 M_{\odot}$ 

- In evolution models, transition between optically thin winds (OB-type stars) and thick winds (WNh stars) is done at H-core burning stage once some threshold on  $X_{\text{surf}}$  is reached: 0.3 for GENEC, 0.4 for Mesa.
- However, there is spectroscopical evidence of WNh with  $X_{\text{surf}} \gg 0.3$  (Martins+08, Crowther+10, Hamann+19, Martins 2023).
- New evolution models with GENEC and MESA accurately predicts WNh stars from  $M_{\text{zams}} \gtrsim 60 M_{\odot}$ .

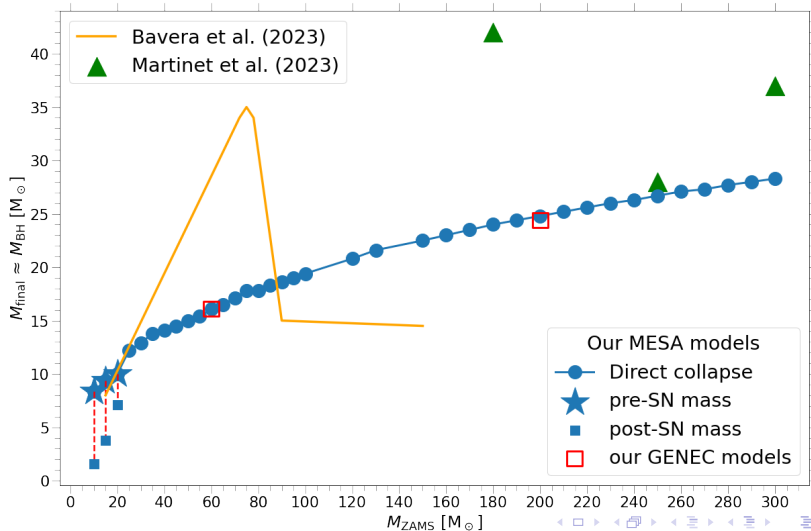
Gormaz-Matamala, Romagnolo & Belczynski (2024, in prep.)



## New BH masses for the MW (Romagnolo et al. 2024)



## Final masses from new winds



# Conclusions

Evolution adopting new self-consistent winds:

- predict stronger chemical enrichment for  $\sim 20 - 40 M_{\odot}$  stars at SMC metallicity, because of a lower loss of angular momentum.
- predict  $M_{\text{bh,max}} \simeq 28 M_{\odot}$  for the Milky Way ( $Z = 0.014$ ), far larger than  $M_{\text{Cyg X-1}} = 21.2 \pm 2.2 M_{\odot}$  and lower than  $M_{\text{Gaia BH3}} = 32.7 \pm 0.7 M_{\odot}$  (formed from a metal poor MW massive star).
- predict WNh stars (WR winds,  $X_{\text{surf}} \geq 0.3$ ) to be descendant of stars with  $M_{\text{zams}} \gtrsim 60 M_{\odot}$ .

## Summary and future work

Based in our new evolution models:

- New evolution tracks for lower ( $Z \ll 0.002$ ) metallicities, plus diagnostics over XShootU sources.
- Evaluate the expected luminosities, abundances and mass losses for Ofpe/WR stars in the GC.
- Extend the transition between thin (O-type) to thick (WNh) winds for VMS based on  $\Gamma_e = 0.5$  to lower metallicities (Tehrani+19, Martins&Palacios22).
- BH/NS masses beyond the MW.



¡Muchísimas gracias!

Thanks a lot!

Dziękuję bardzo!

Muito obrigado!