

Astronomická spektroskopie

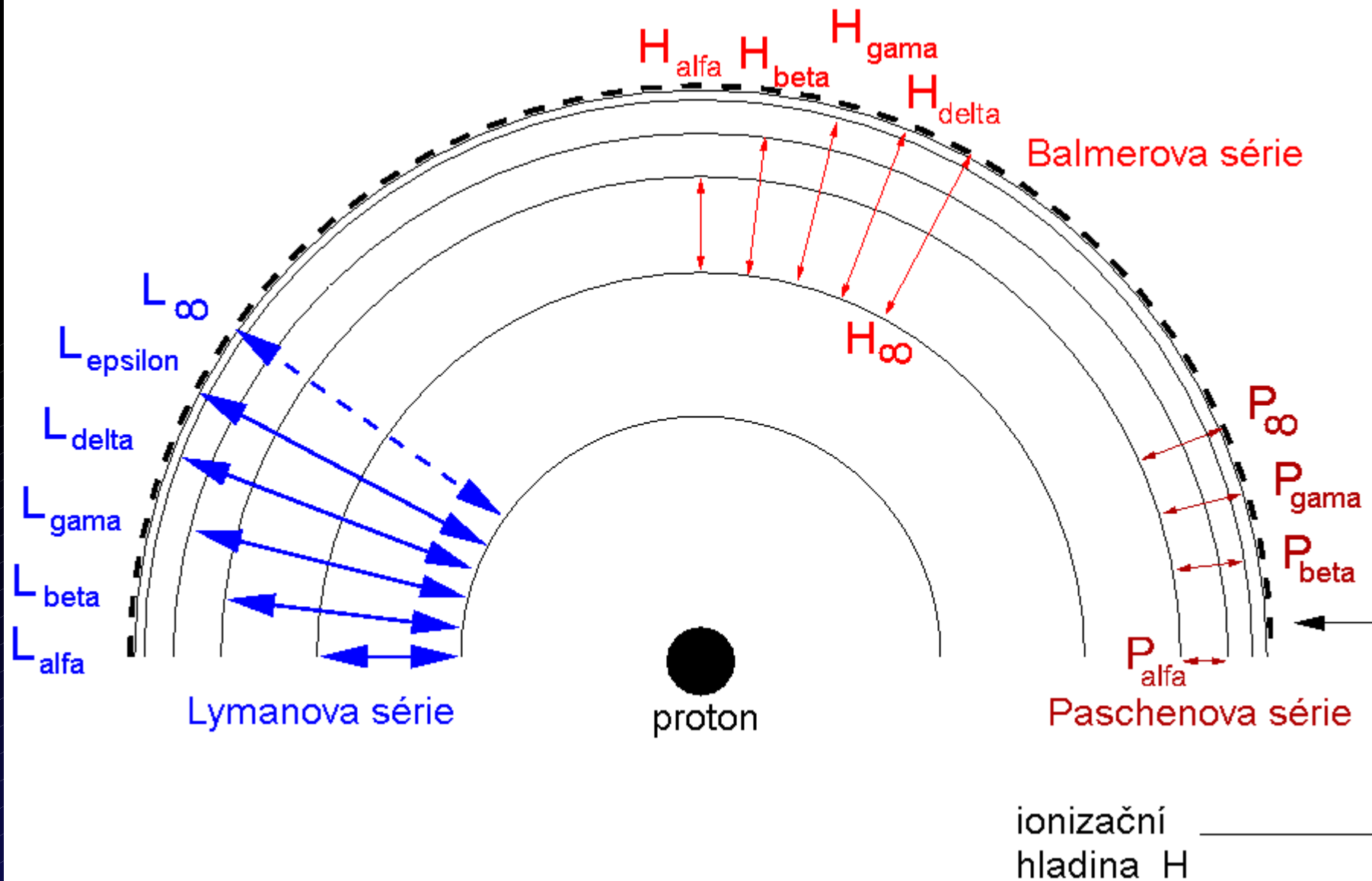
Základy a aplikace

Petr Škoda
Astronomický ústav AVČR

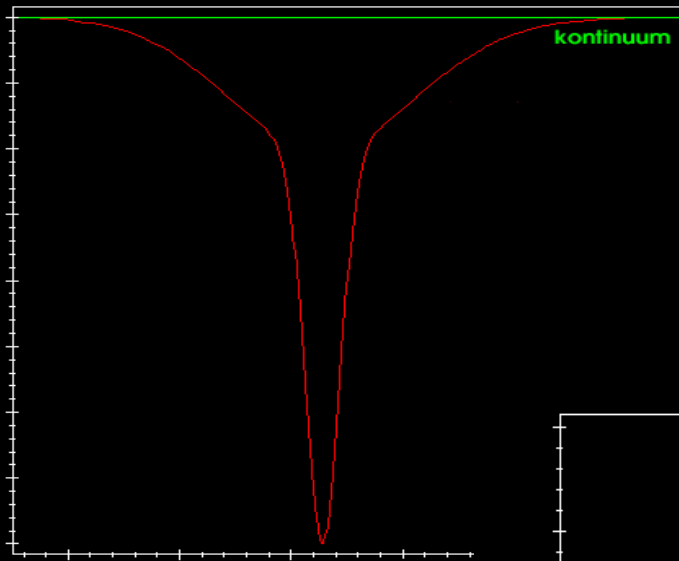
S použitím materiálu od J. A. Smithe (GEMINI) a M. Šlechty (ASU Ondřejov)

Workshop o amatérské spektroskopii, Valašské Meziříčí 5.5.2007

Vznik spekter

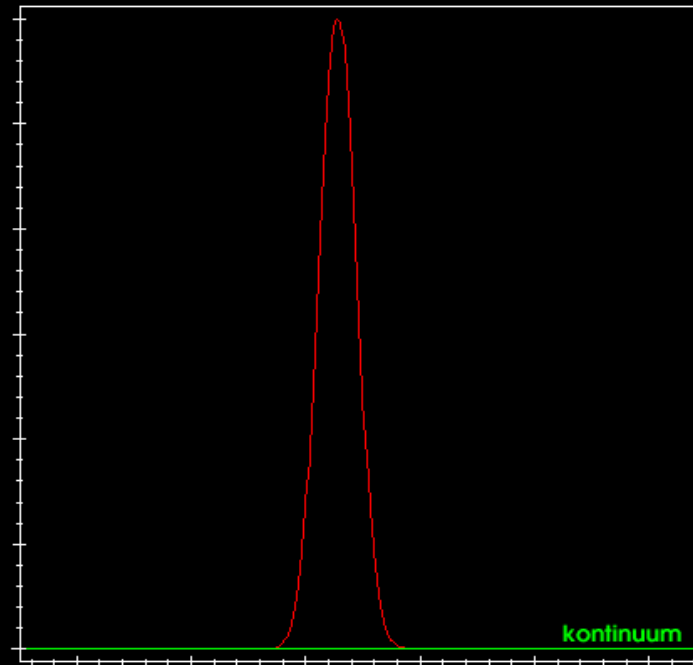


absorpční čára

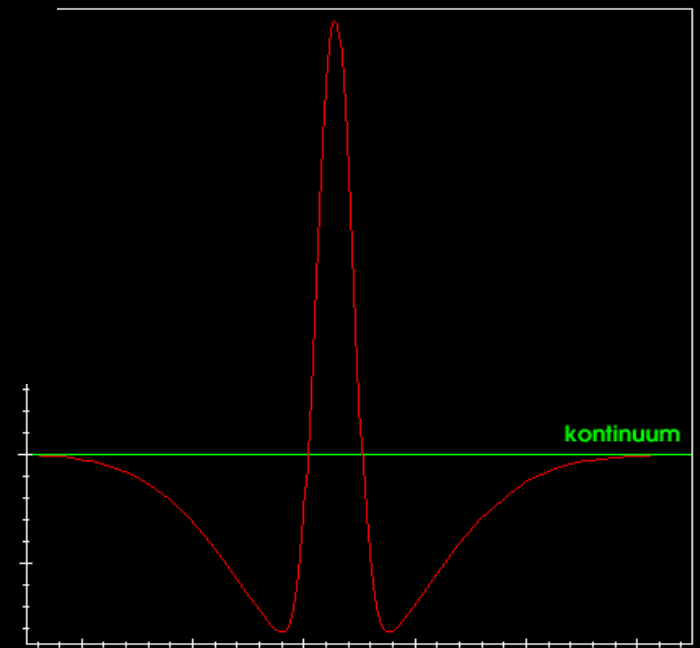


Typy čar

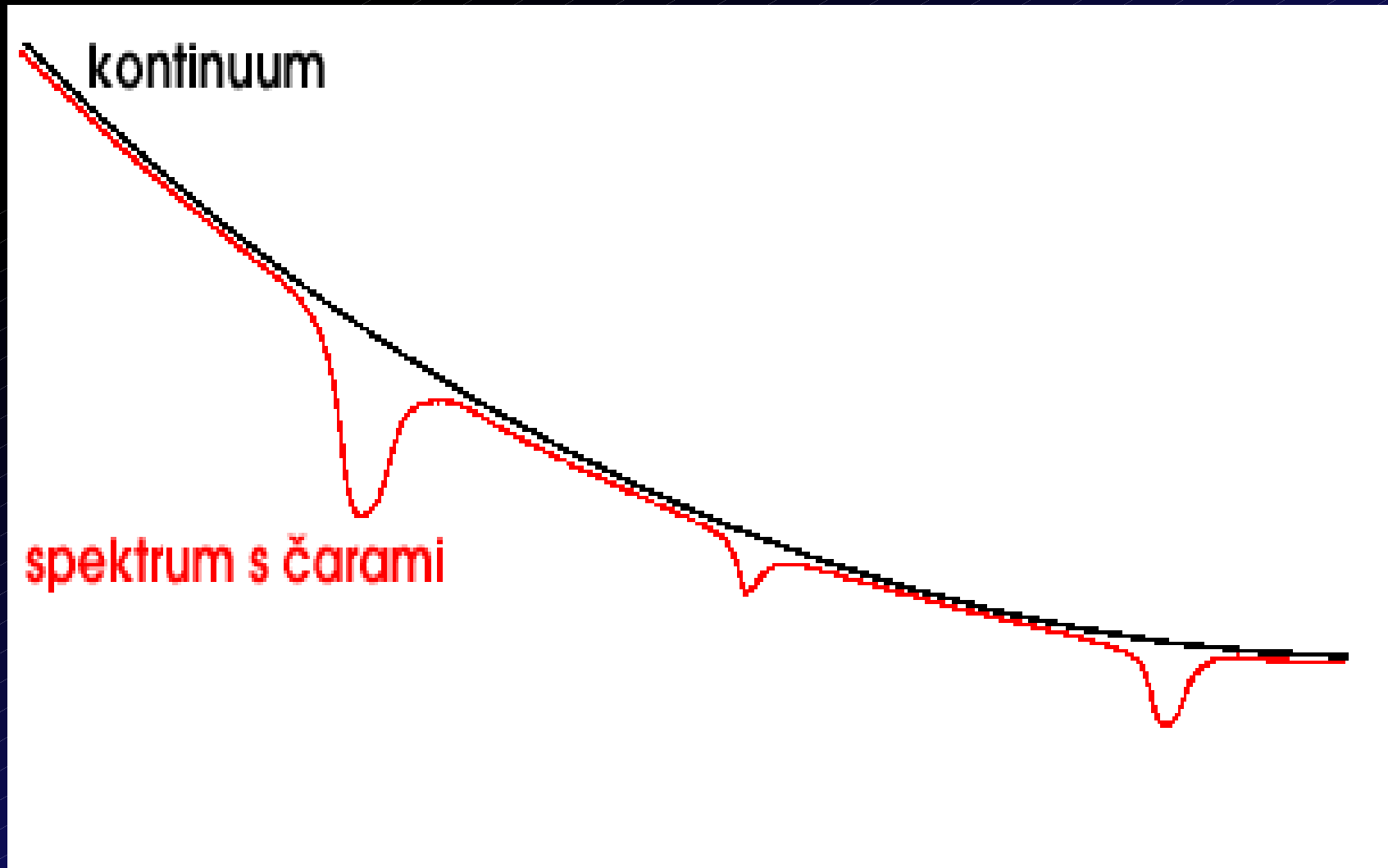
emisní čára (bez absorpčních křídel)



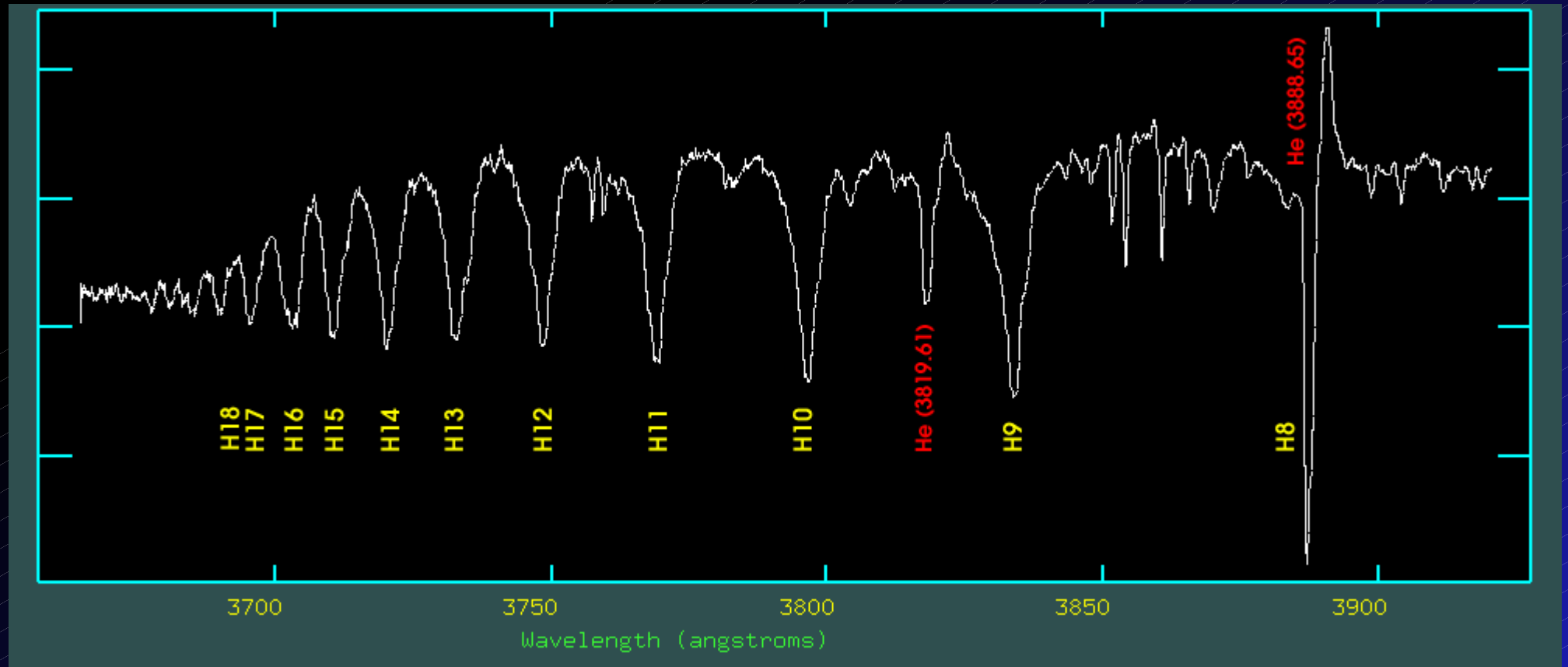
emisní čára (s absorpčními křídly)



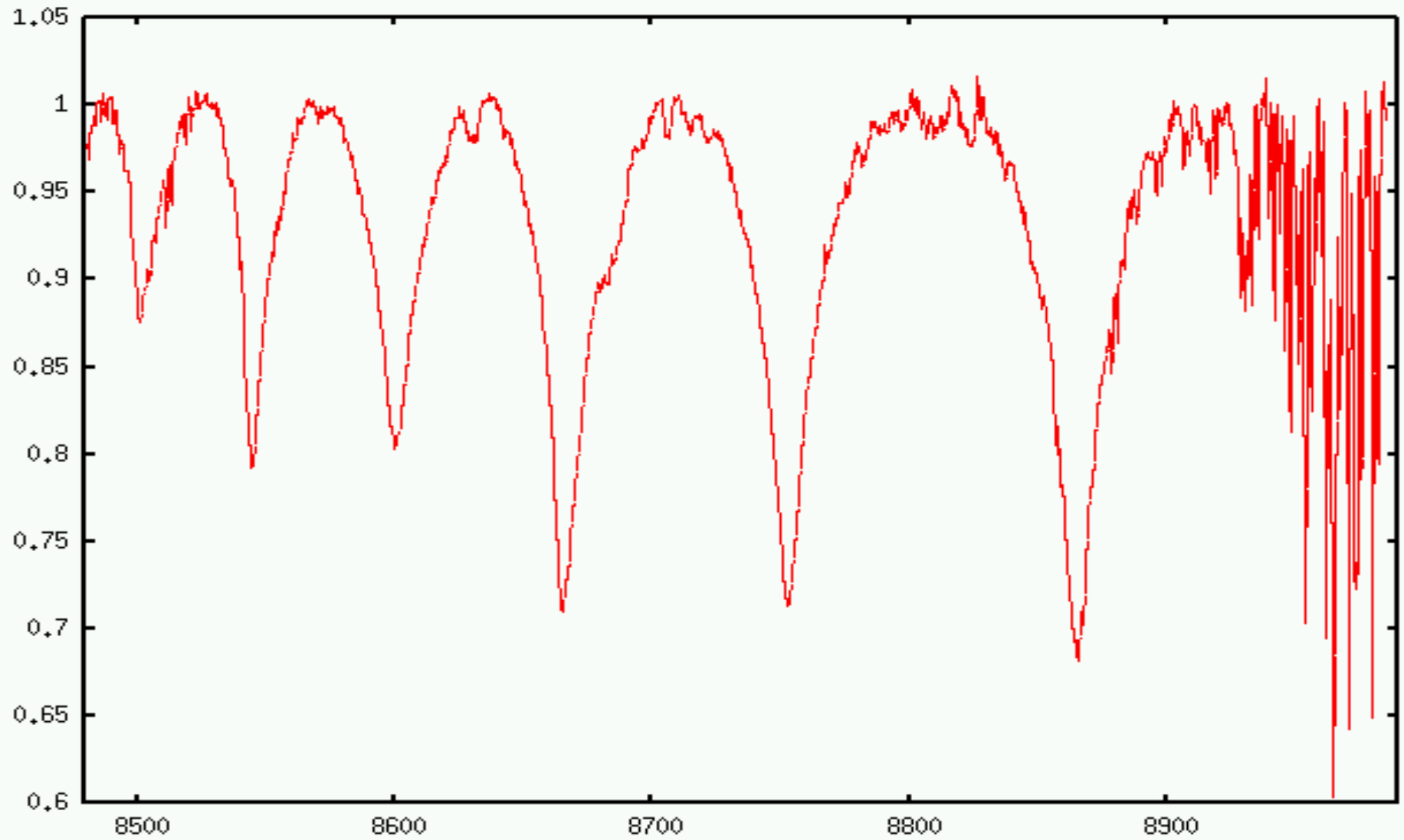
Tvar kontinua (Planck ?)



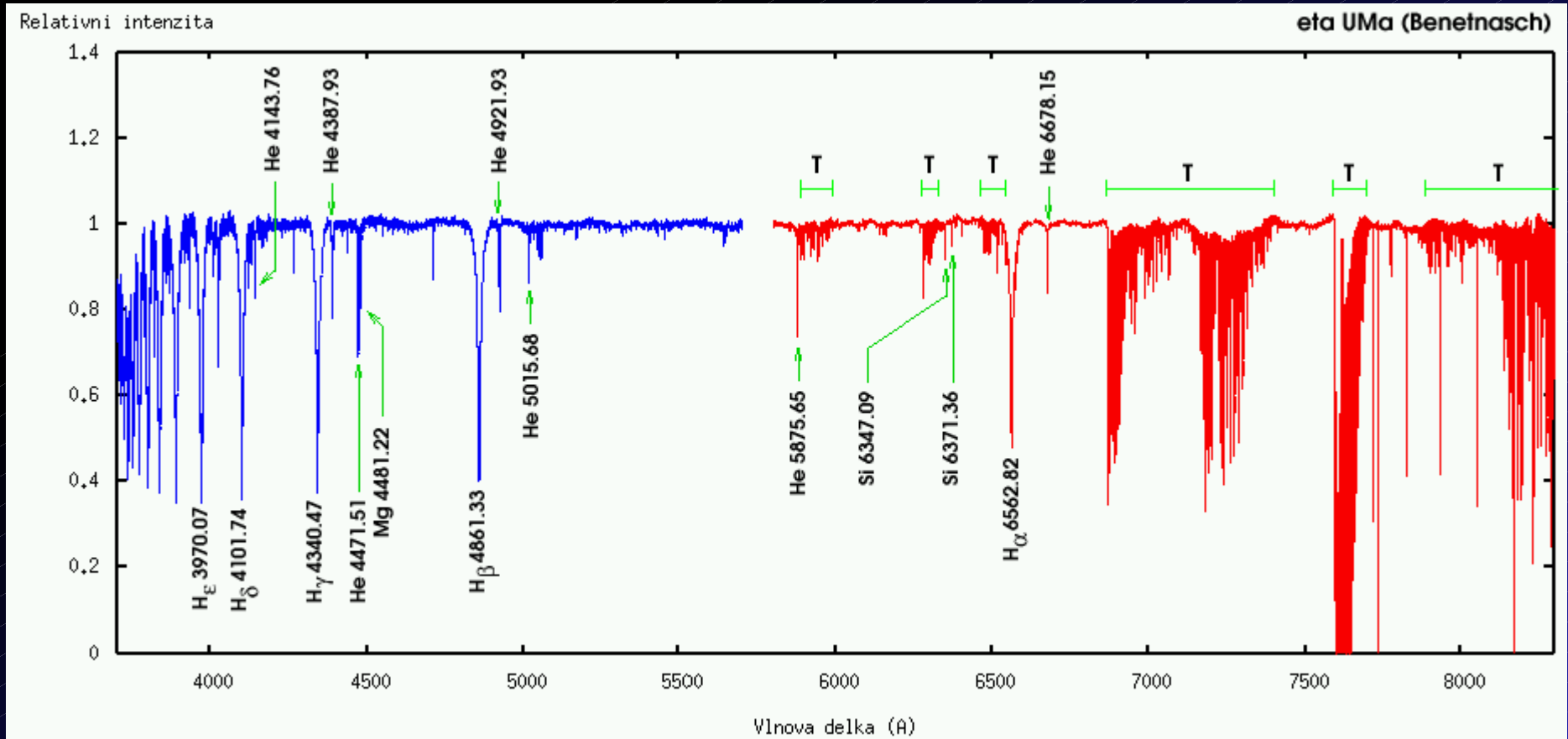
Balmerova hrana



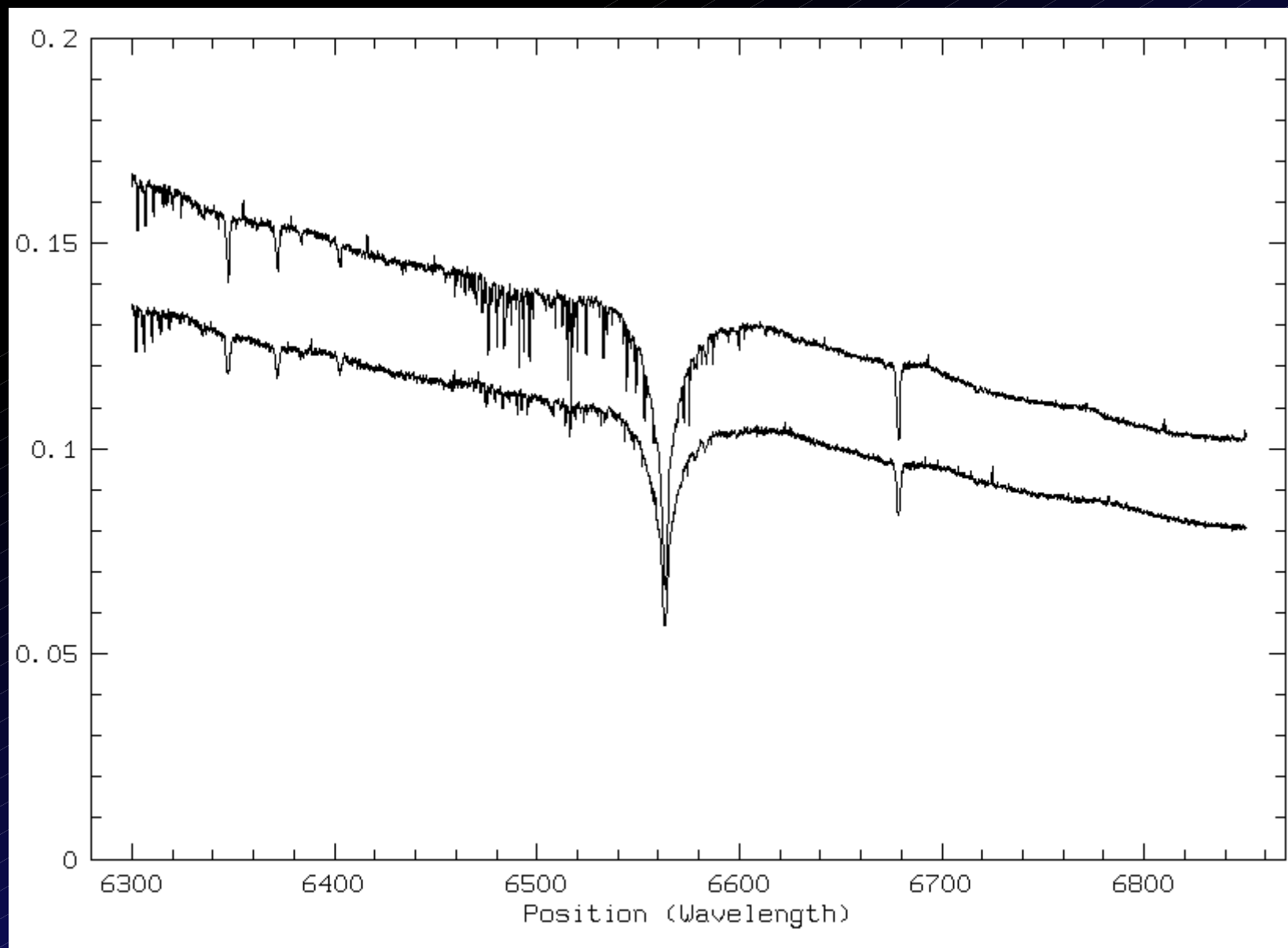
Paschenova hrana



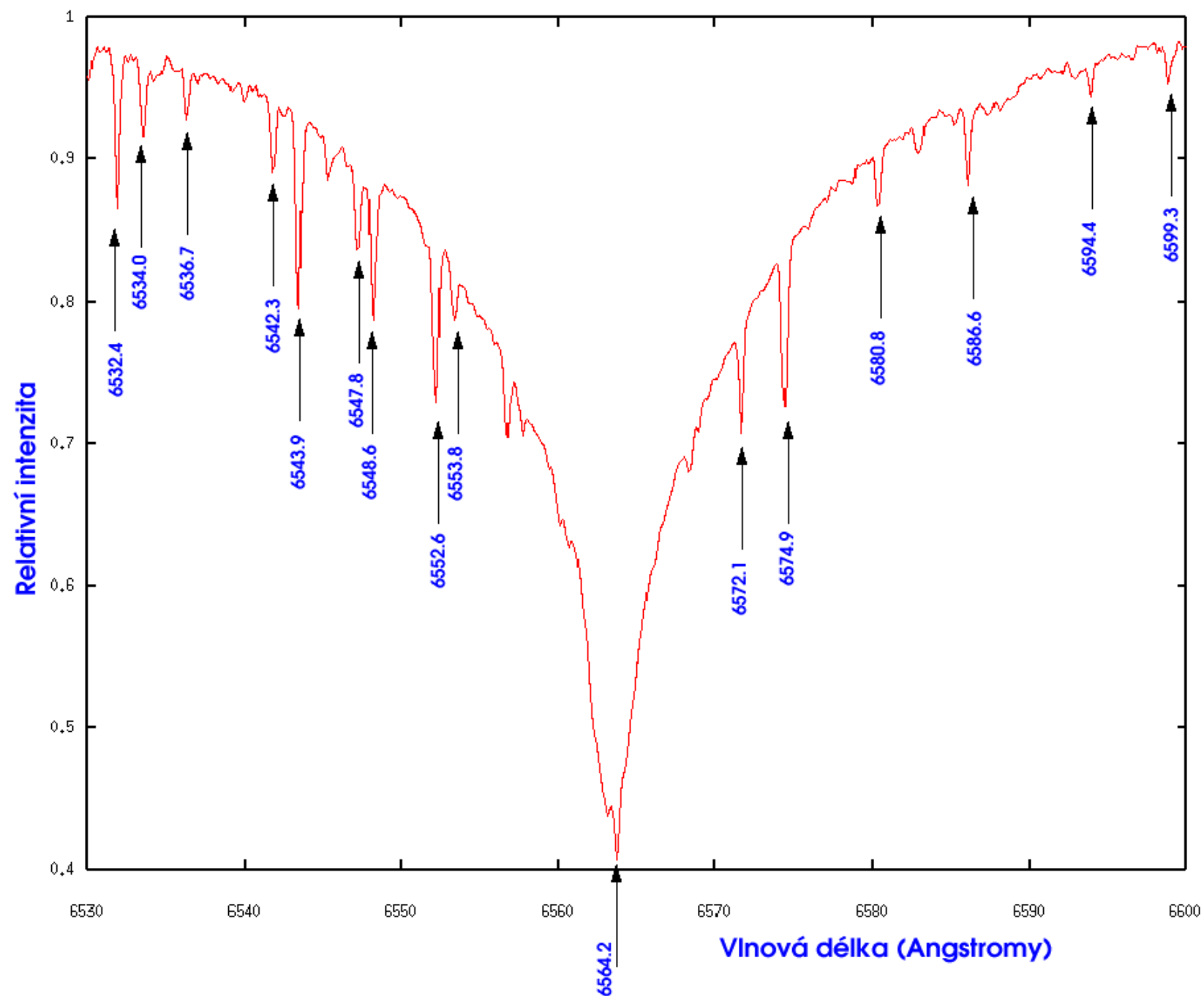
Typické čáry horkých hvězd



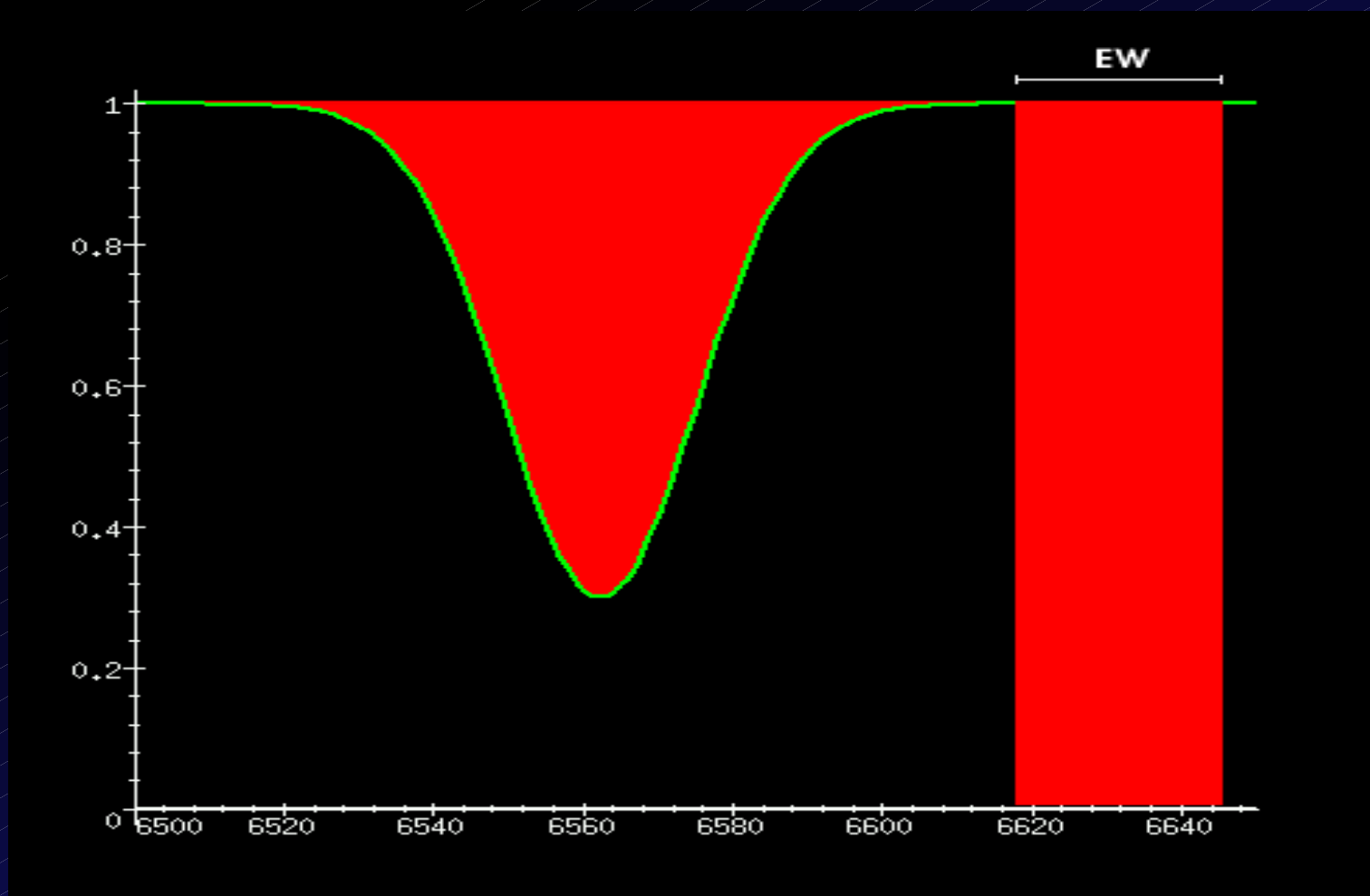
Telurické čáry – vlhkost



Polohy telurických čar v oblasti čáry vodíku H alfa

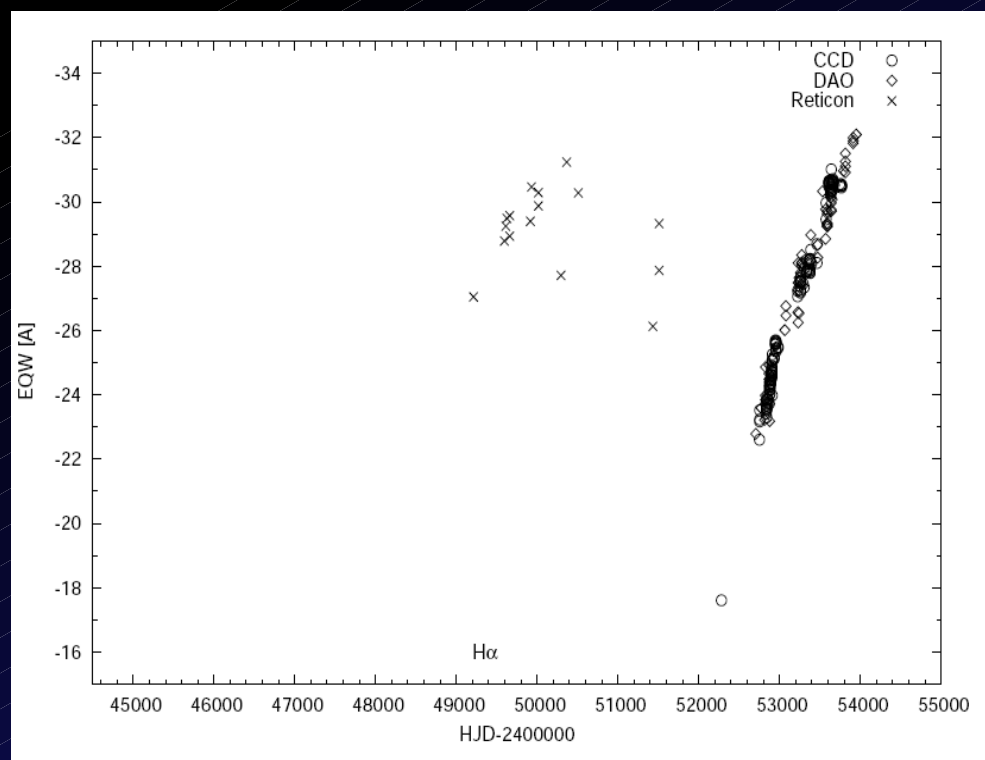


Ekvivalentní šířka

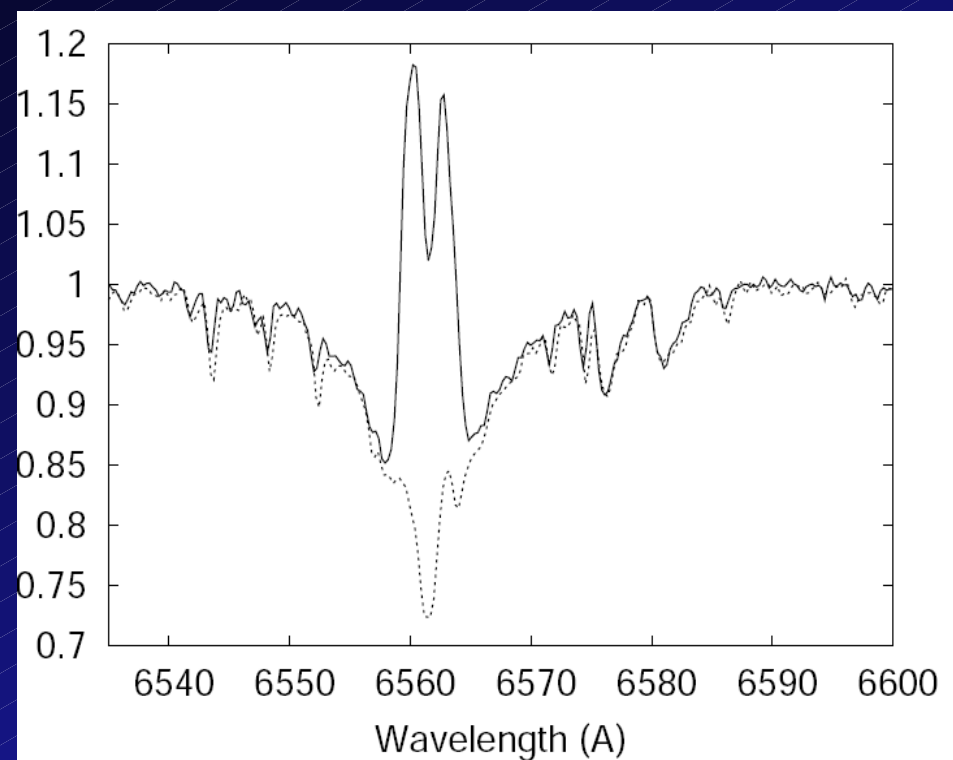


Changes of EW in Time

- Moments in general (asteroseismology – mode identification)
- Emission lines – negative EW
- Estimates of expanding shell
- Sensitive to continuum placement – shallow lines
- Abundances - check normalization – different lines same ion

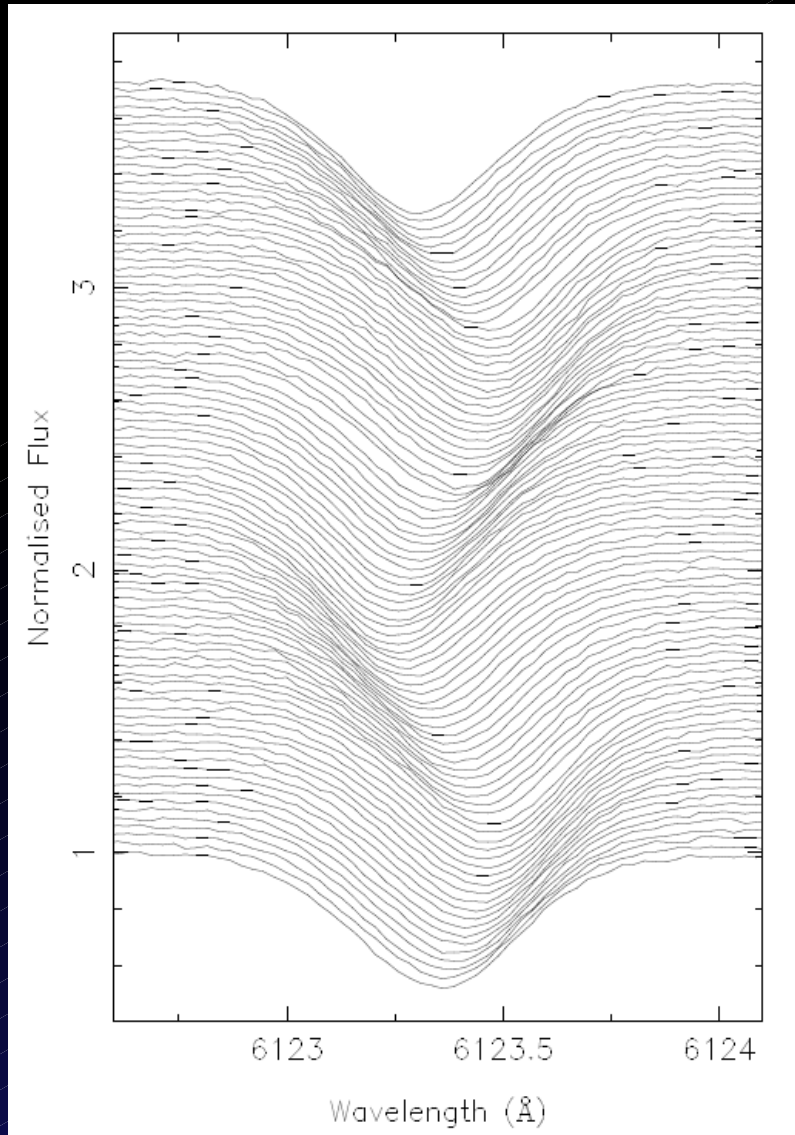


Omi Cas : Koubsky et al. 2004

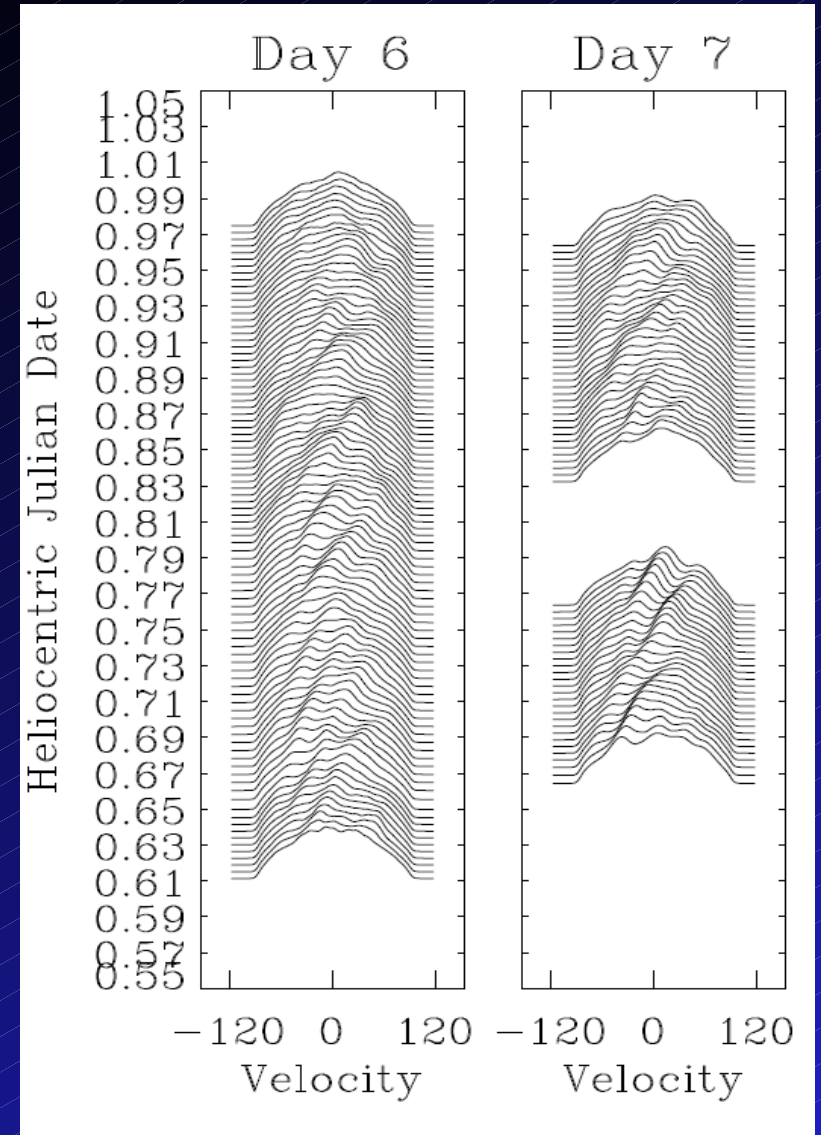


HD6226 : Slechta and Skoda 2004

Measured Pulsations

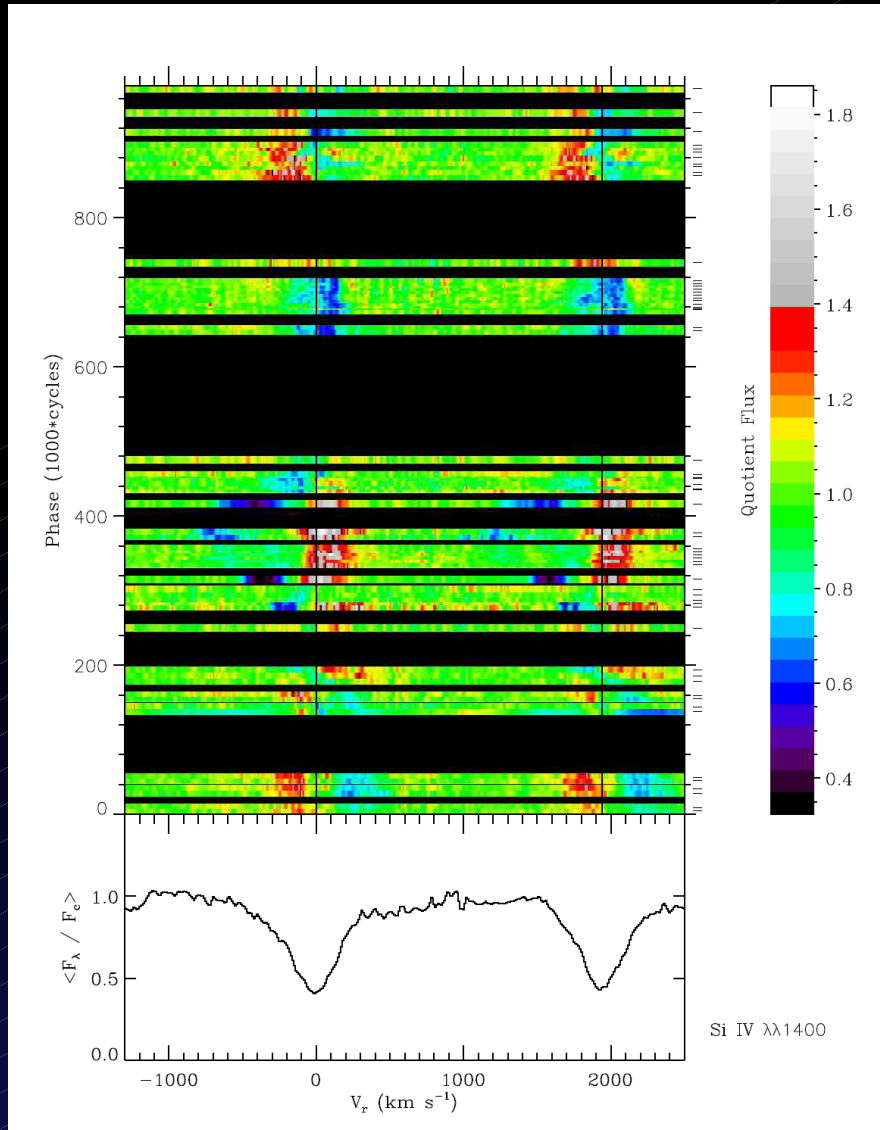


Rho Pup – del Sct type



Eps Cep - del Sct type

Dynamic Spectra



Interactive features, color cuts, LUT

Multiple lines at the same time

Quotient or Differential

X axis : λ or RV

Y axis : phase or time

For study of LPV (asteroseismology, winds)

Requires

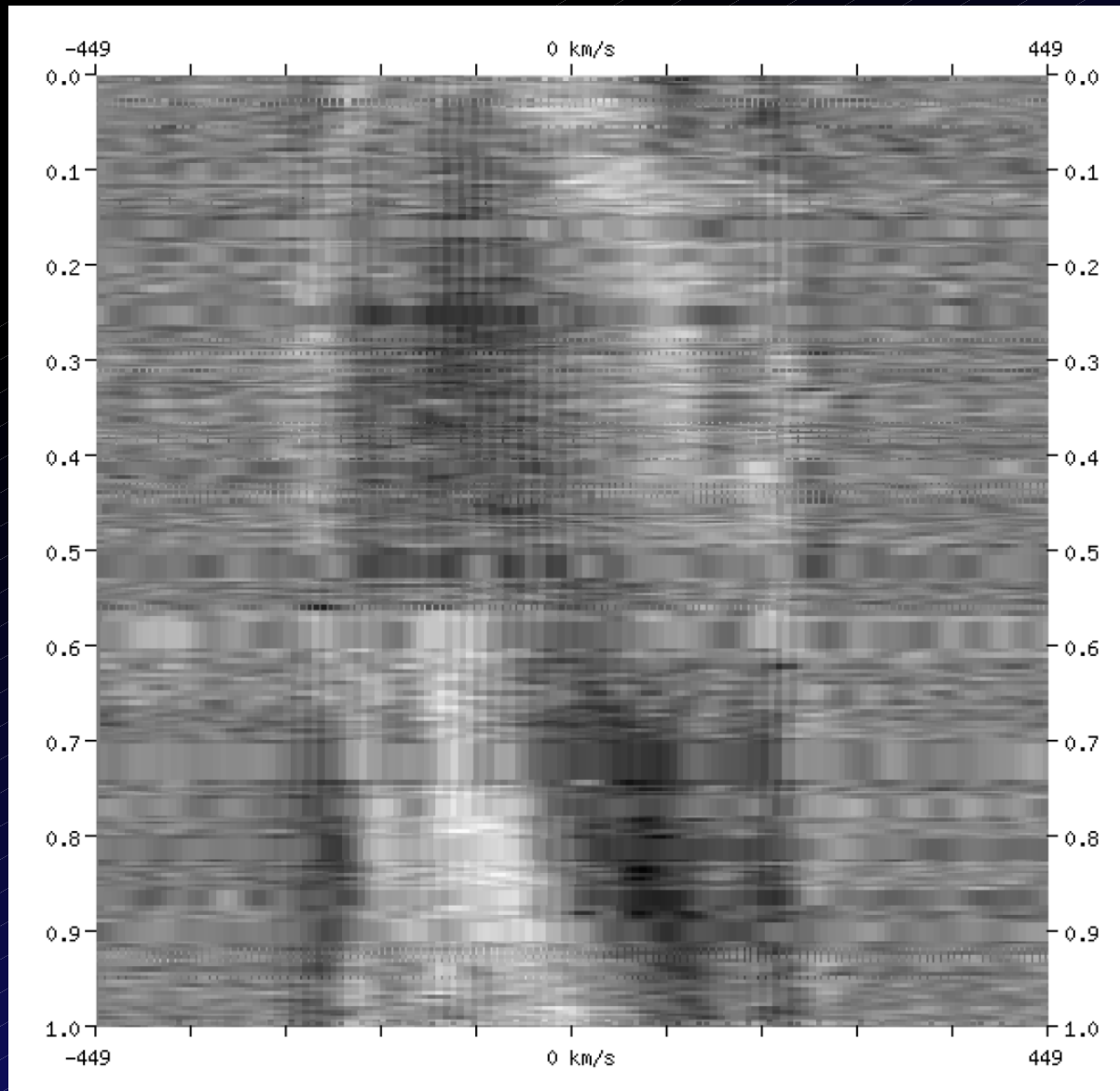
time (JD) - winds

period (see Period analysis) - phase (LPV)

change of template (average, median)

removing bad data (interactive overplotting)

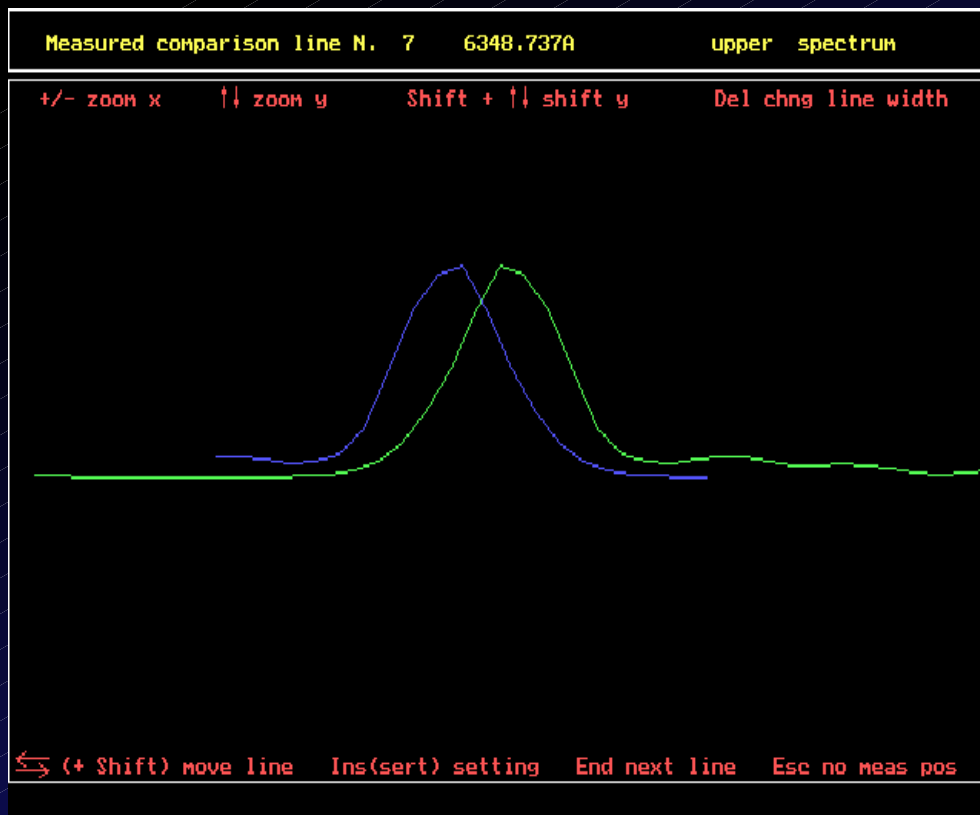
Dynamic Spectra



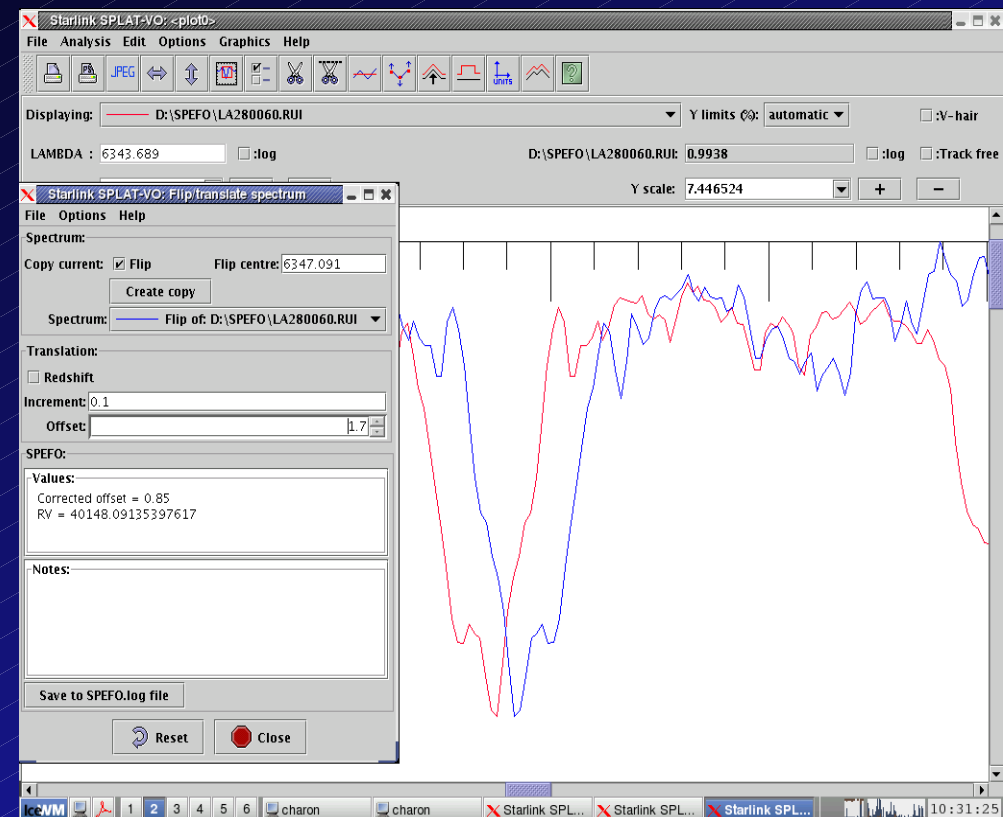
Netolický 2004

RV by Mirroring

Shift until best match of direct and flipped profile - interactive
Complicated profiles (Be)
Adjustable region of interest (wing/core)
Needs reference line position



SPEFO- Oscilloscopic



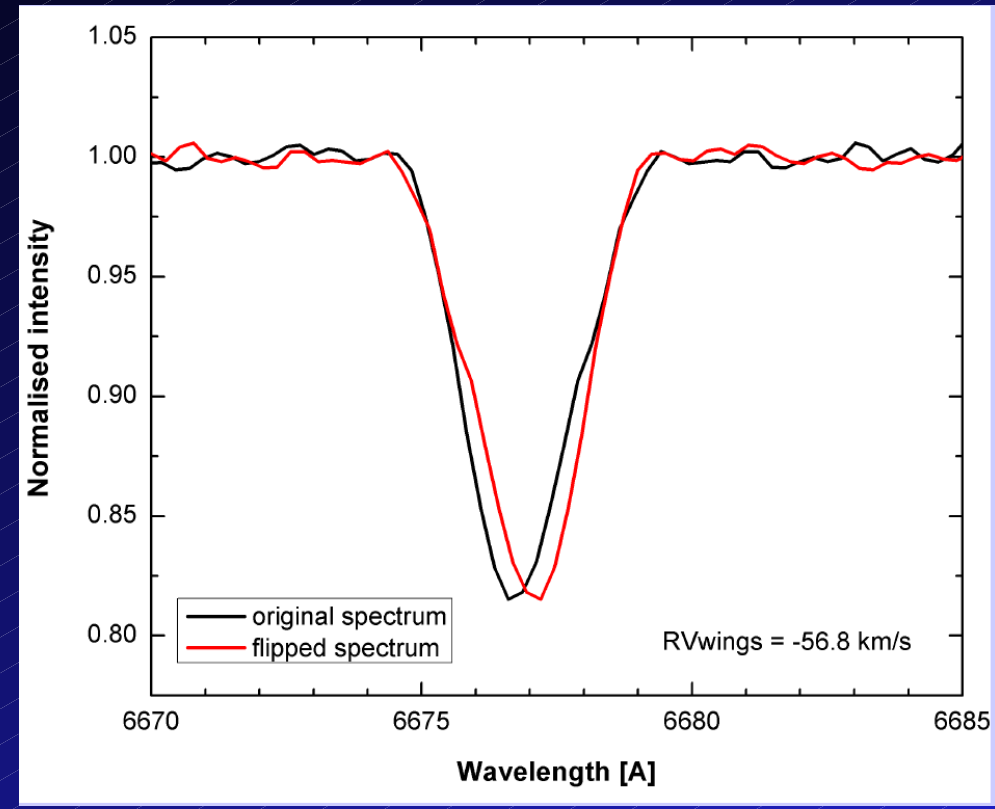
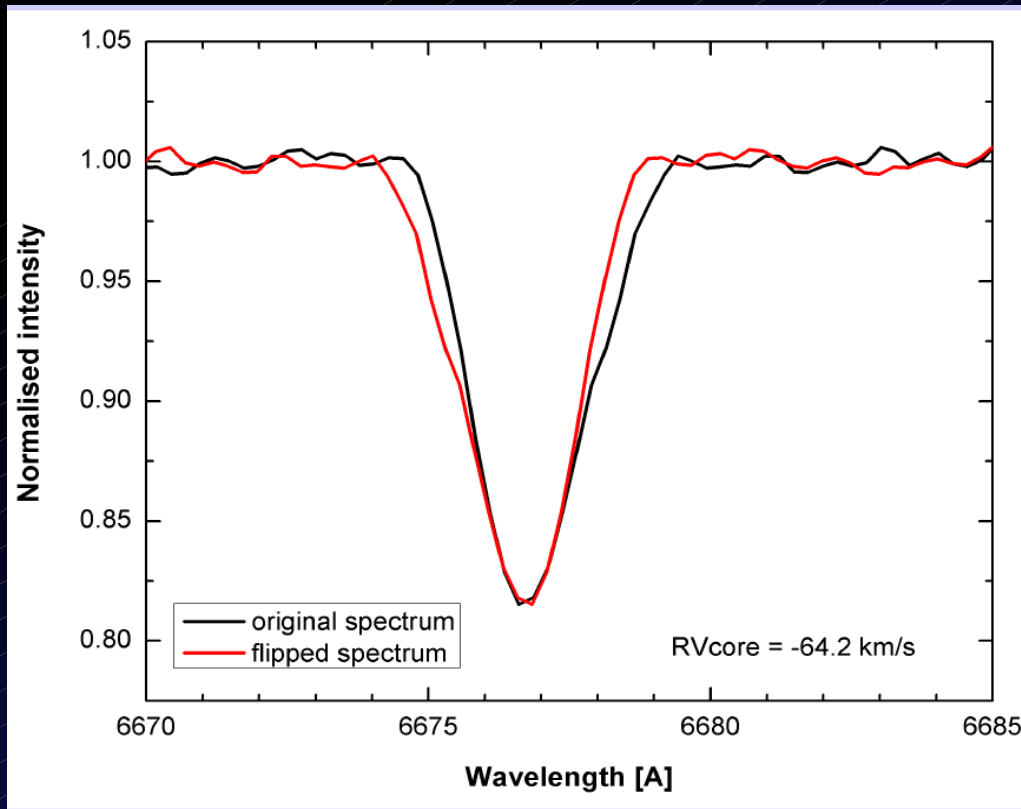
SPLAT

Mirroring Method

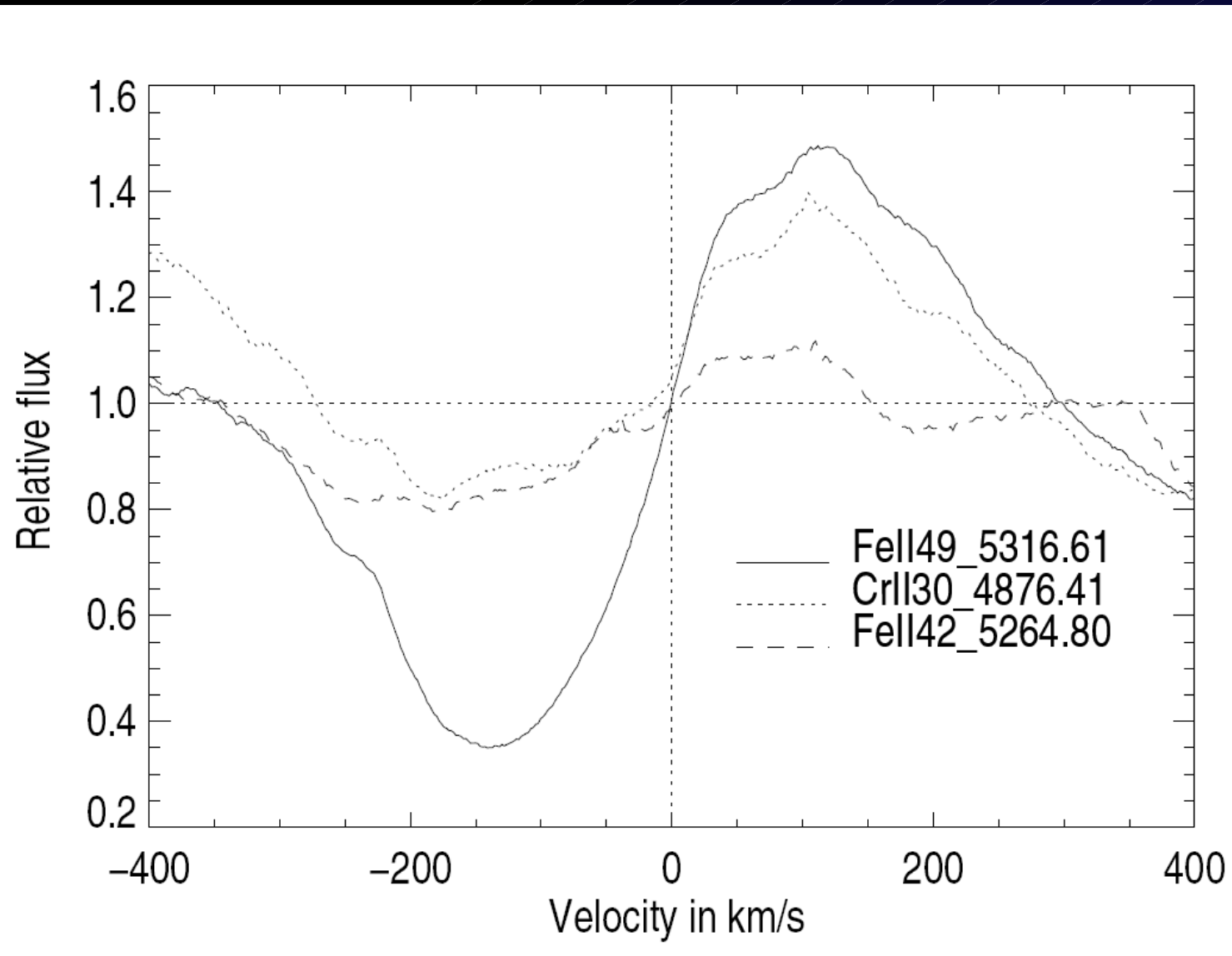
Separate match of core from match of wings – where in depth ?

Different physics (shells, shears, winds)

Asymetry – how to handle ?



Different Lines overplotted RV scale



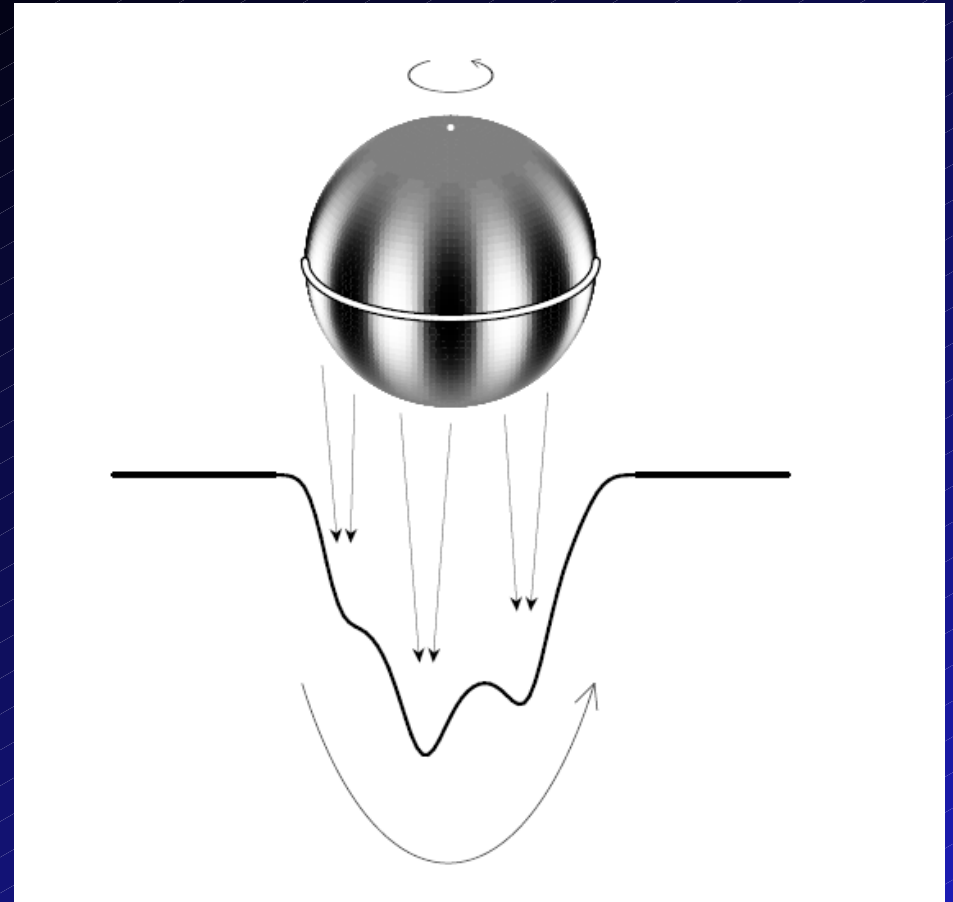
Doppler Imaging - Bumps

Vogt & Penrod -80s

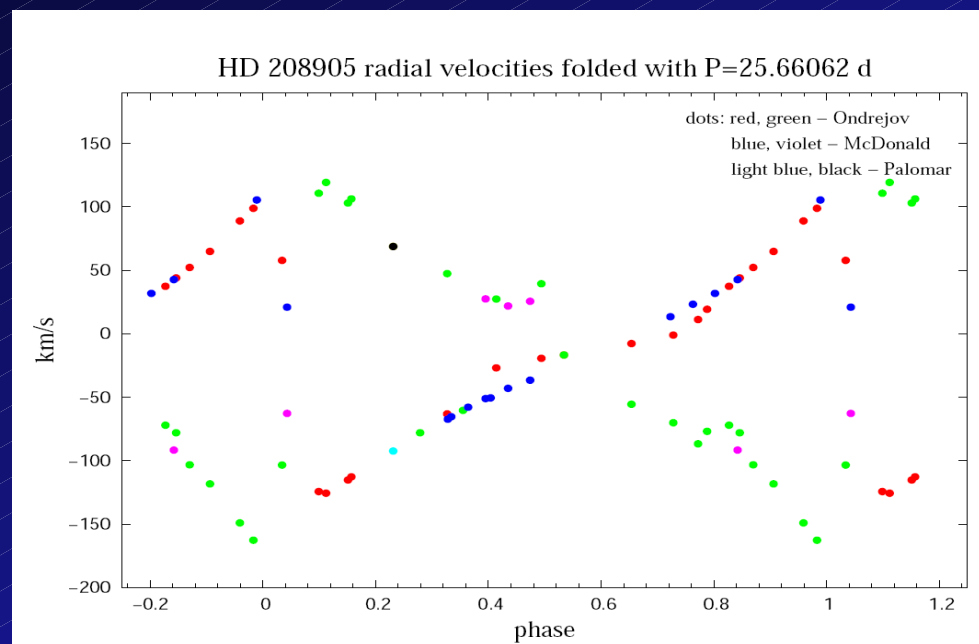
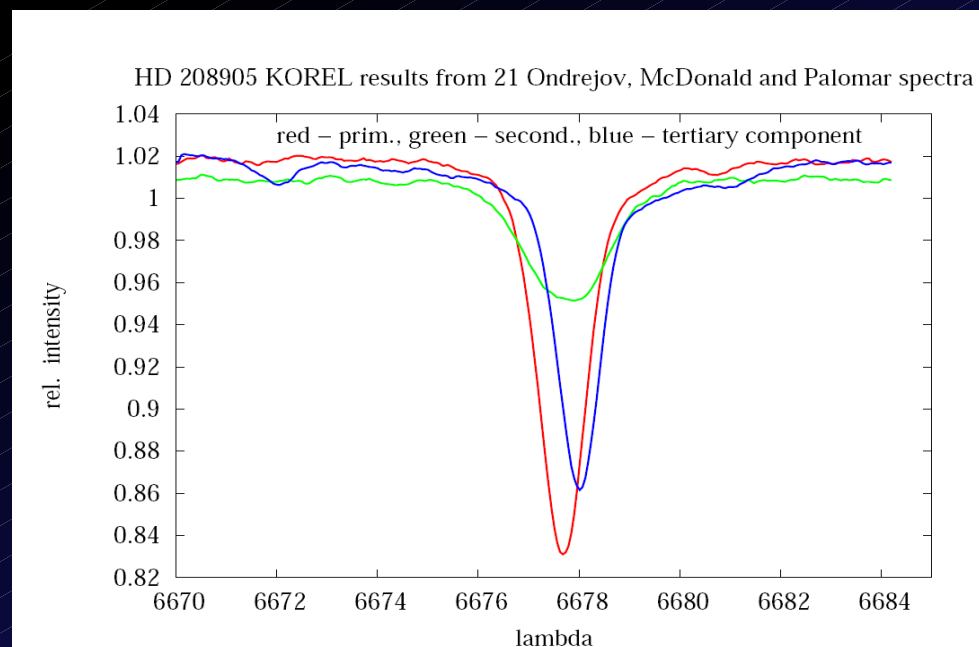
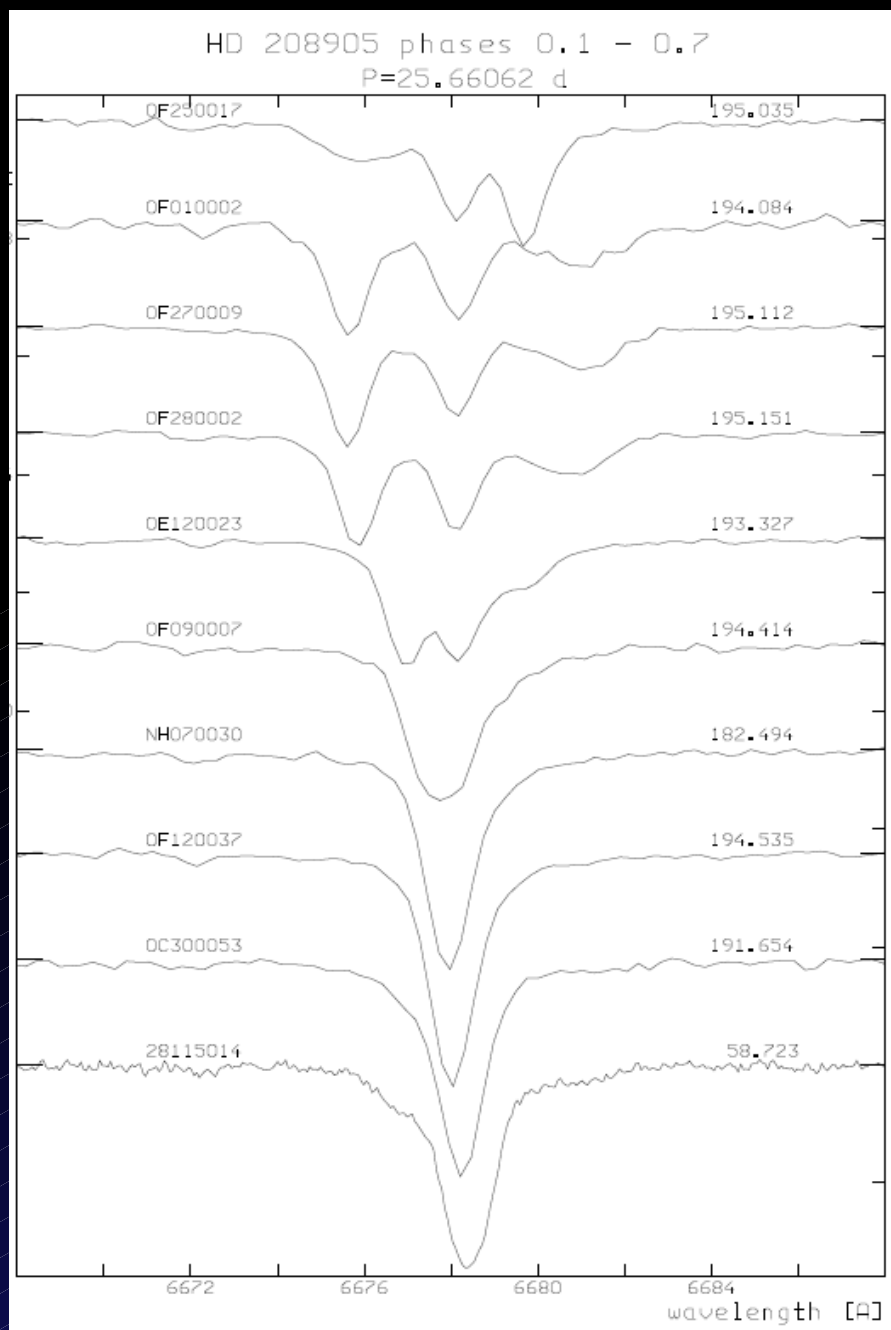
Zet Oph

Ball VOGT

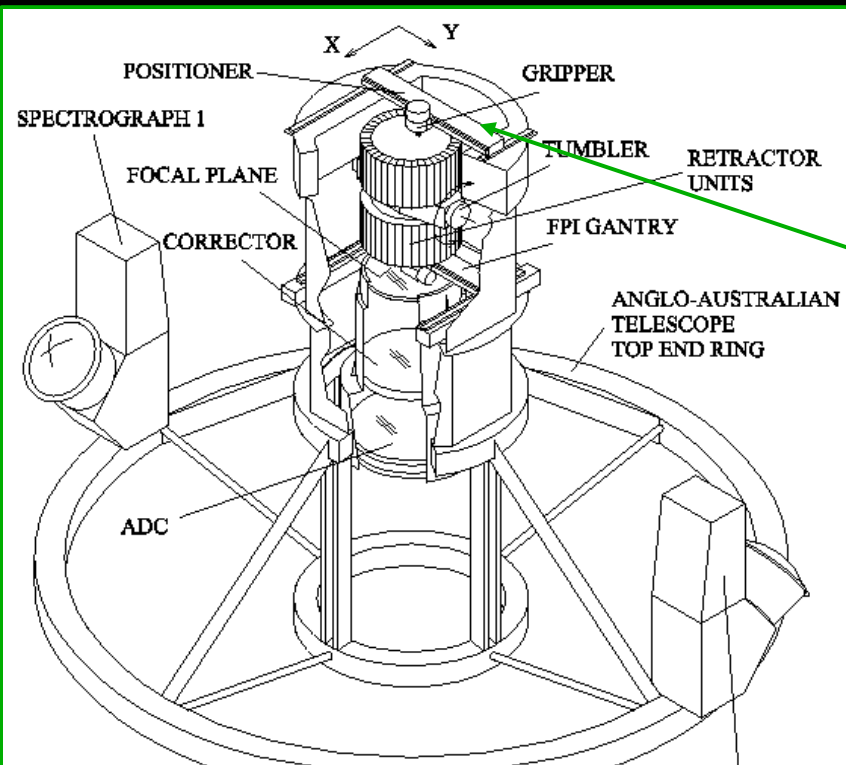
- Requires high SNR
 - ($>300-500$)
 - Perfect rotation coverage
 - Artefacts otherwise
 - NRP or solar spots
- Doppler Tomography
 - Accretion jets in Algols
 - Orbits in RV phase space



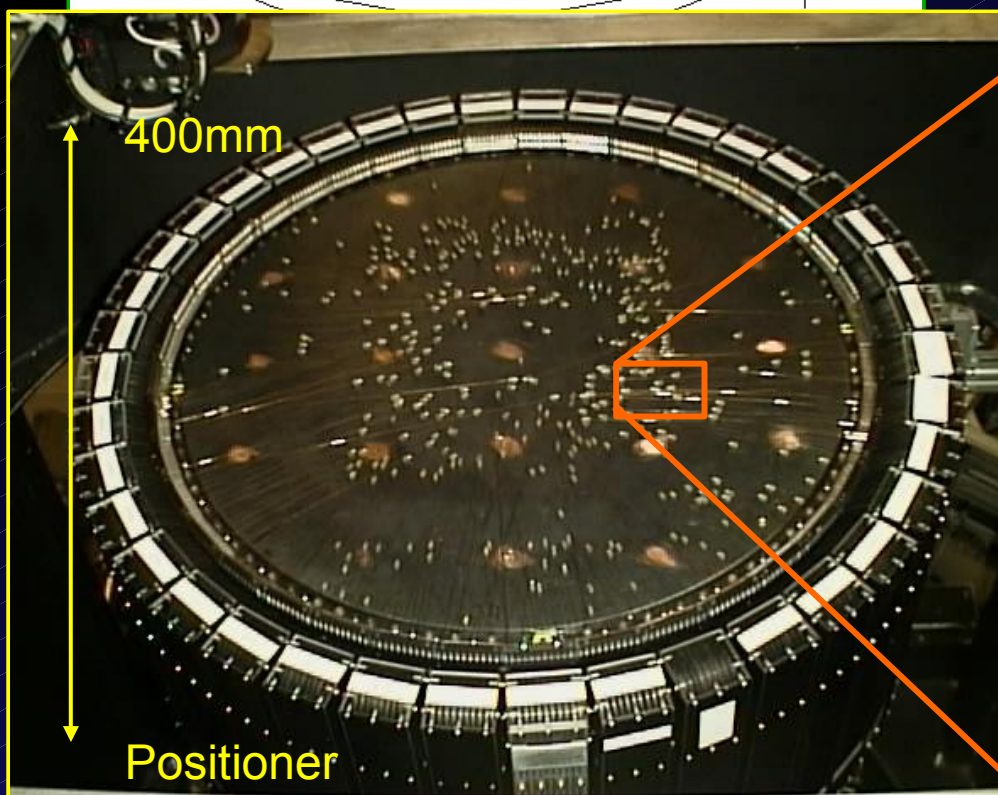
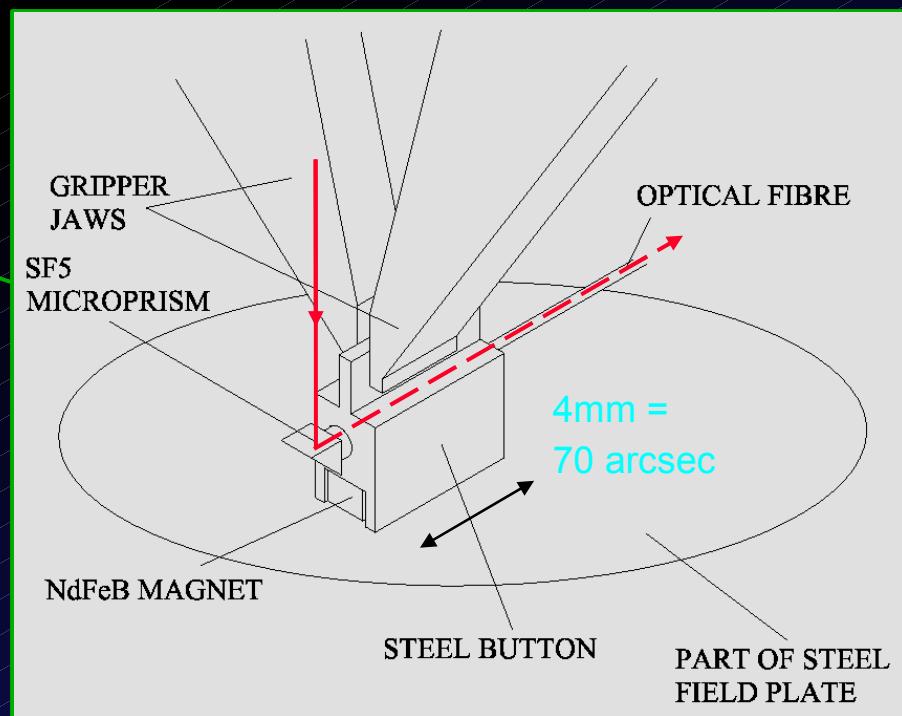
Spectra Disentangling in Fourier Space - KOREL



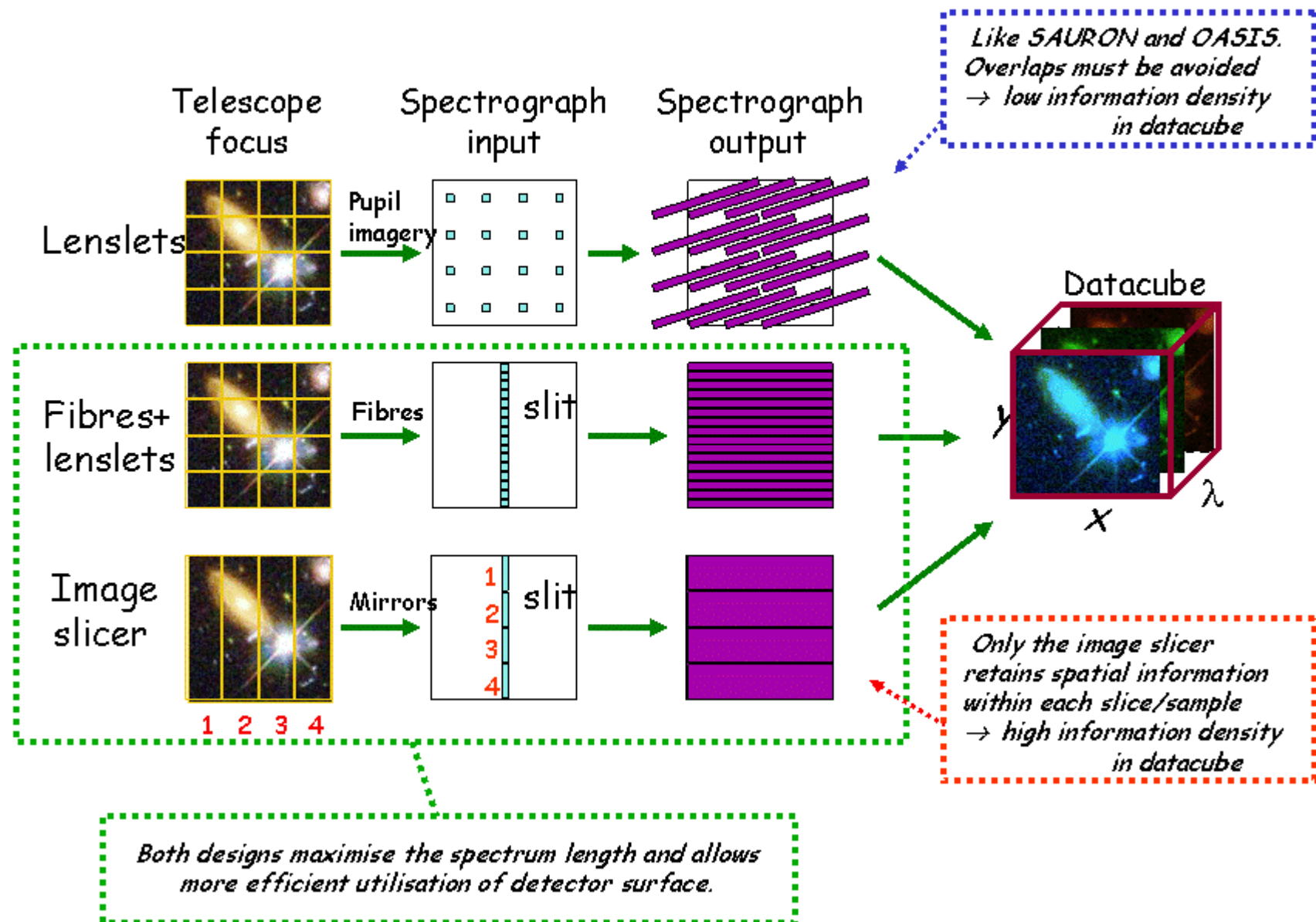
HD208905: Koubsky et al. 2006



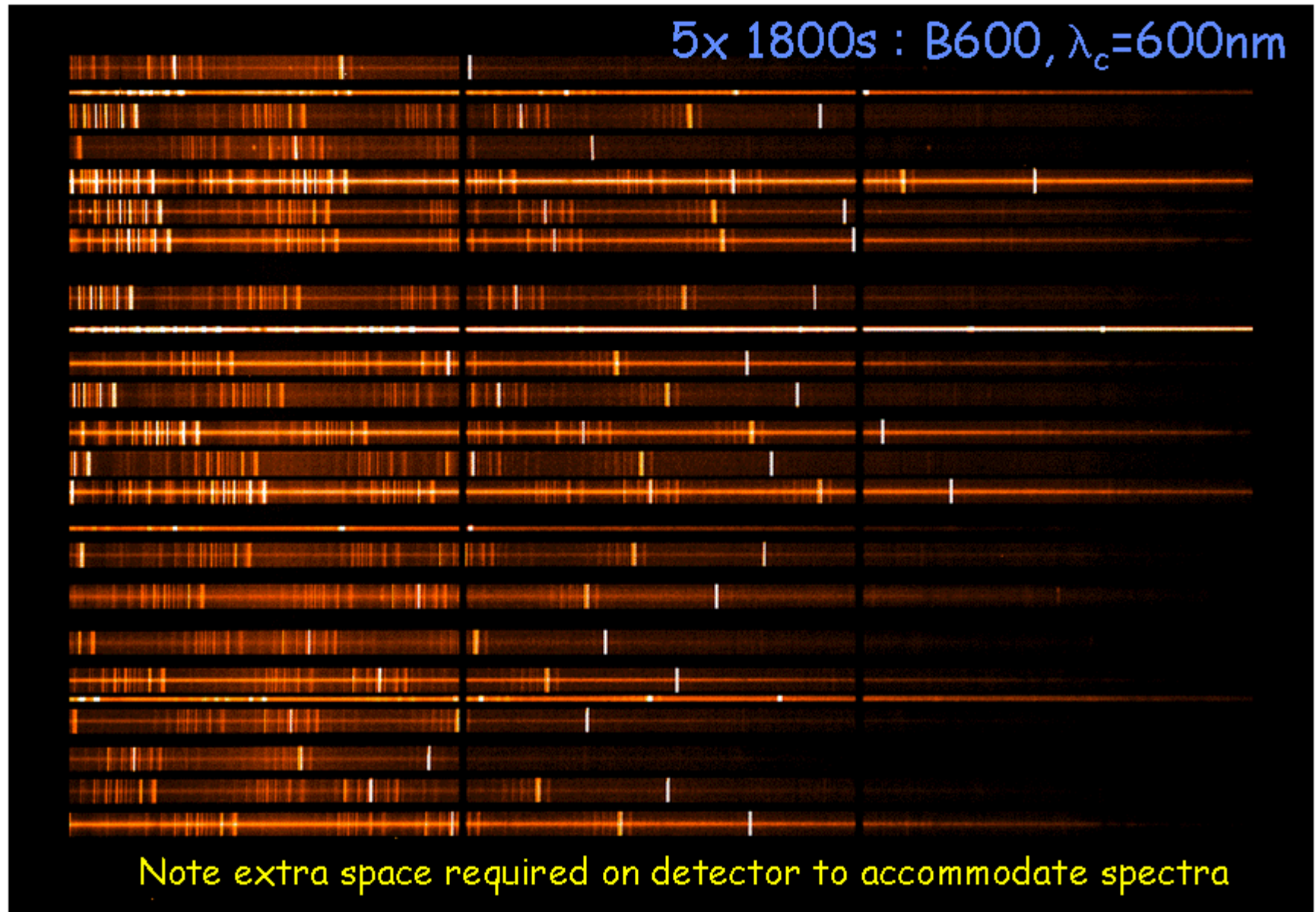
2dF



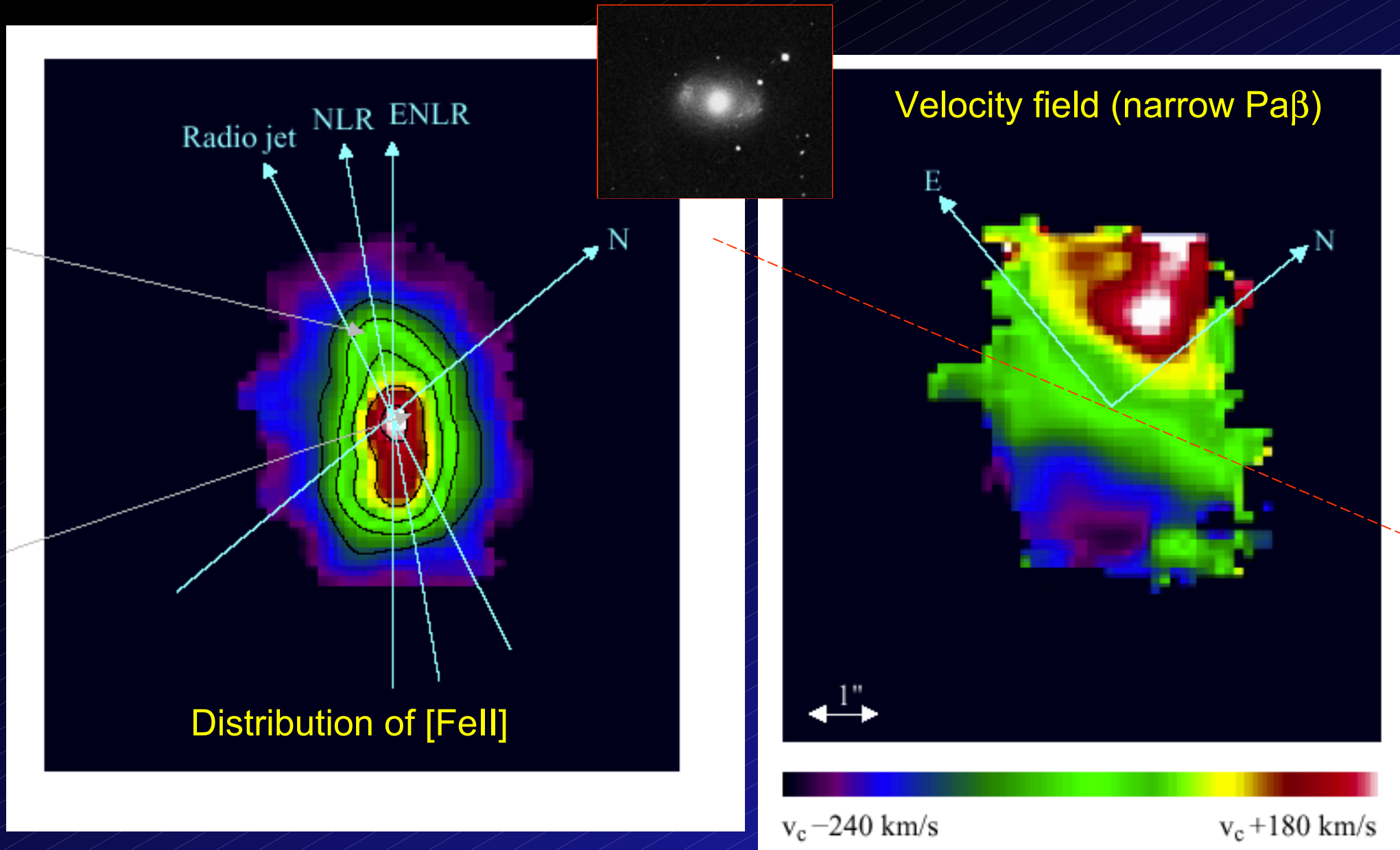
Techniques of IFS



GMOS multislit example

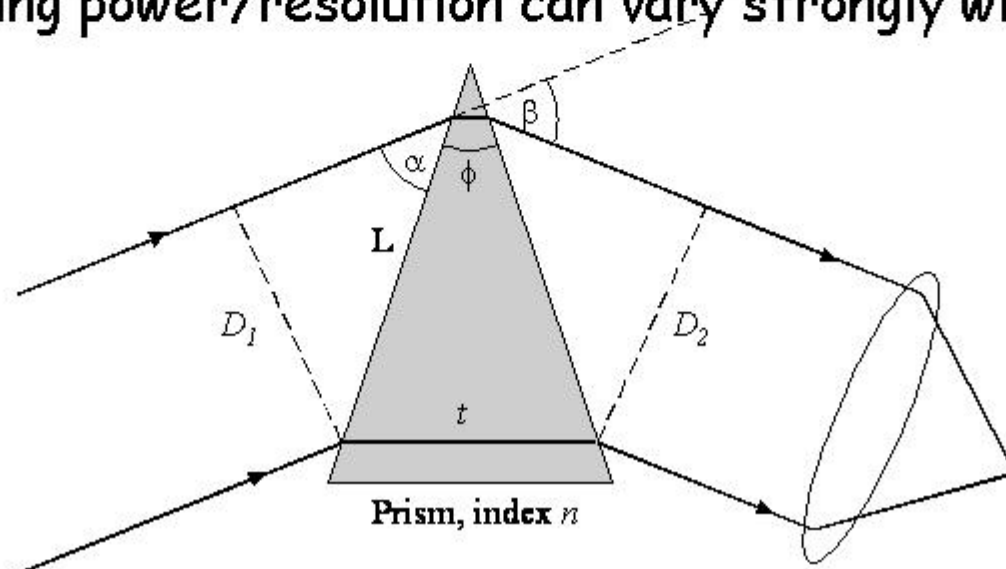


Dissecting active galaxies



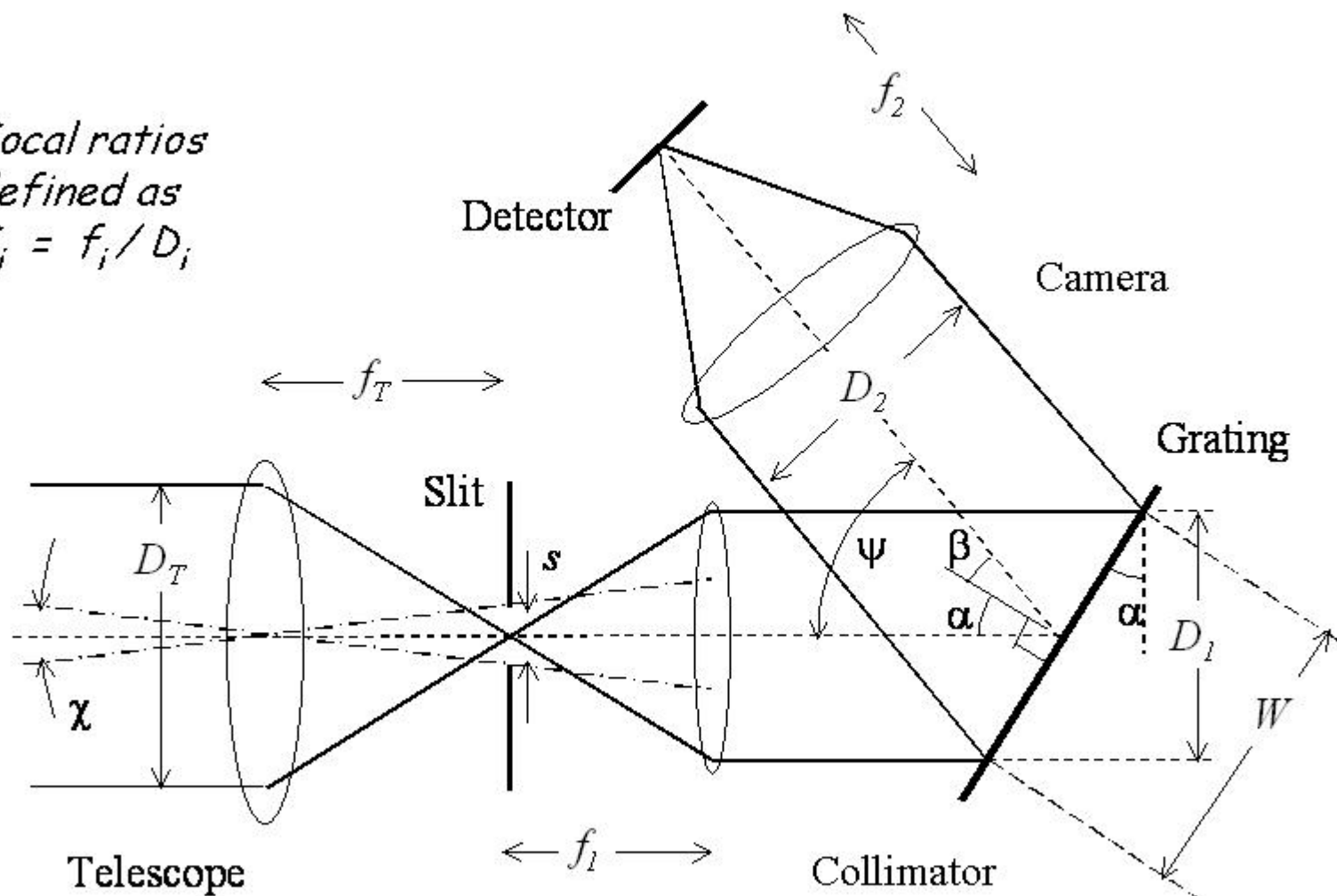
Prisms

- Useful where only low resolving power is required
- Advantages:
 - simple - no rulings! (but glass must be of high quality)
 - multiple-order overlap not a problem - only one order!
- Disadvantages:
 - high resolving power not possible
 - resolving power/resolution can vary strongly with λ

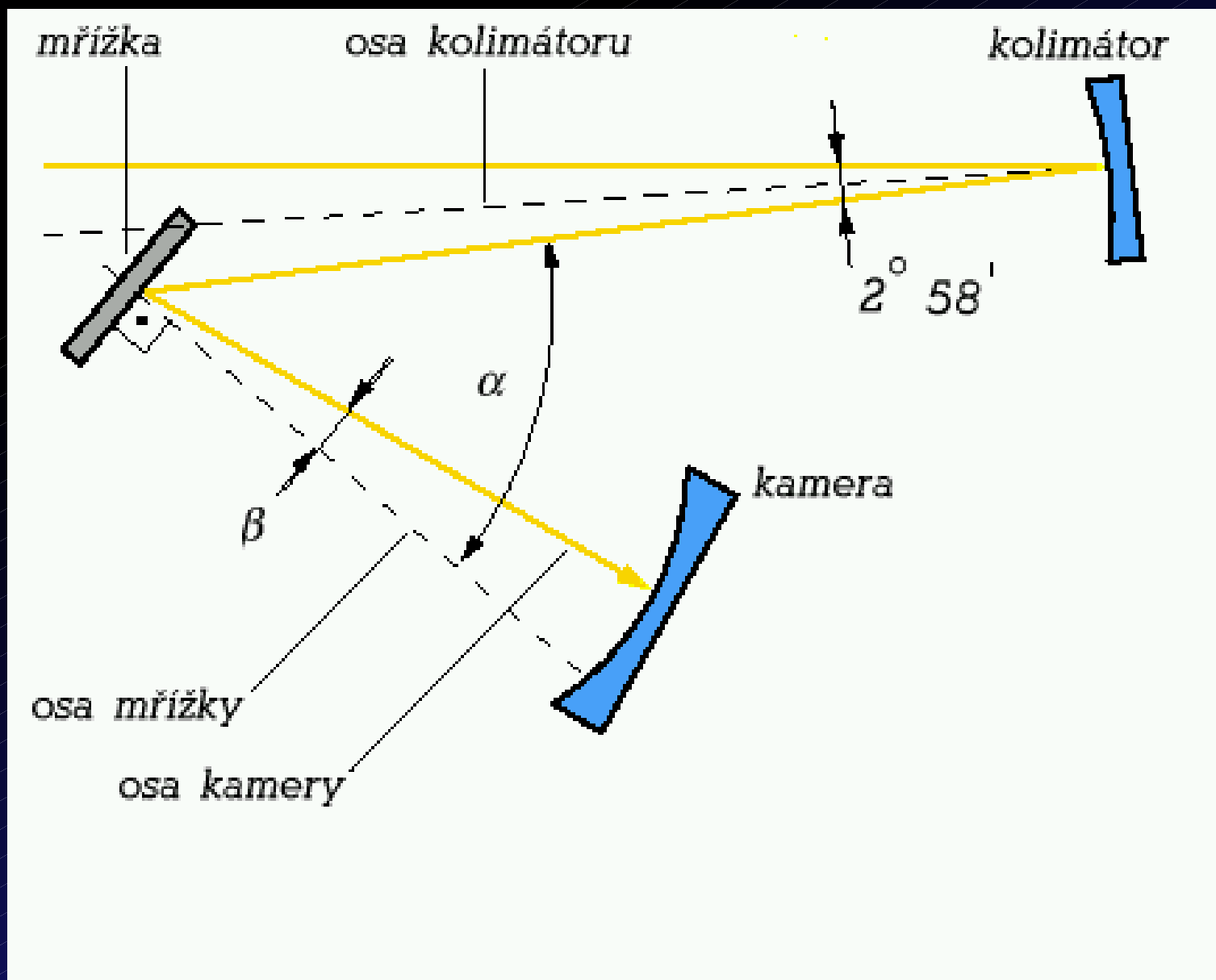


Generic spectrograph layout

Focal ratios
defined as
 $F_i = f_i / D_i$



Coudé spektrograf 2m v Ondřejově

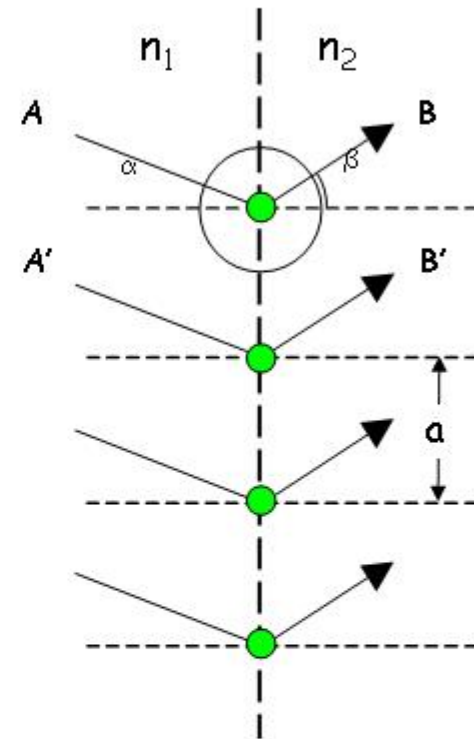


Grating equation

- Interference condition:
 $\Rightarrow$ path difference between AB and $A'B'$

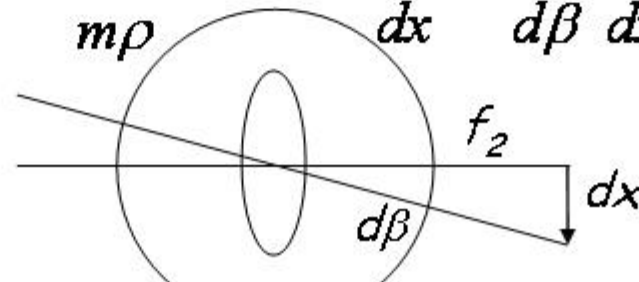
- Grating equation:

$$m\rho\lambda = n_1 \sin \alpha + n_2 \sin \beta \quad \text{where} \quad \rho = \frac{1}{a}$$



- Dispersion:

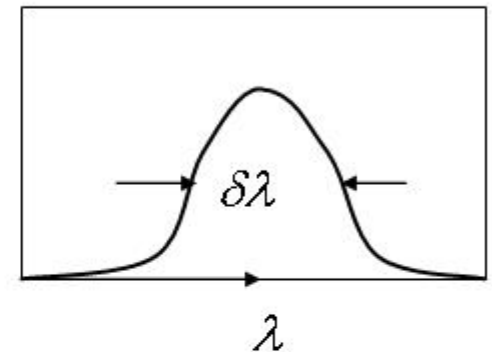
$$\frac{d\lambda}{d\beta} = \frac{\cos \beta}{m\rho} \quad \frac{d\lambda}{dx} = \frac{d\lambda}{d\beta} \frac{d\beta}{dx} = \frac{\cos \beta}{m\rho f_2}$$



"Spectral resolution"

- Terminology (sometimes vague!)

- Wavelength resolution $\delta\lambda$
- Resolving power $R = \lambda / \delta\lambda$



- Classically, in the diffraction limit,

Resolving power = total number of rulings
x spectral order

I.e. $R^* = m\rho W$

Total grating length

- But in most practical cases for astronomy ($\chi < \lambda/D_T$), the resolving power is determined by the width of the slit, so $R < R^*$

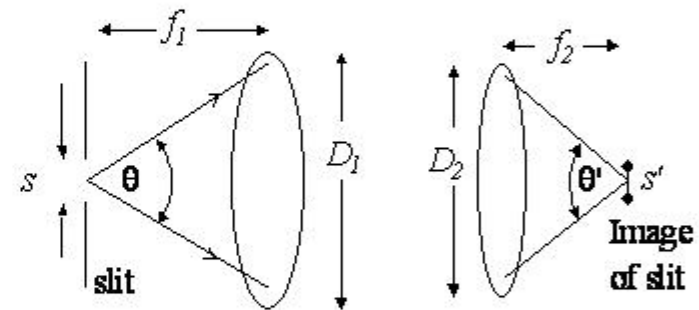
Spectral resolution

- Spectral resolution: $\delta\lambda = \left(\frac{d\lambda}{dx}\right) s' = \frac{\cos \beta}{m\rho f_2} s'$ Image of slit on detector

- Projected slit width:
Conservation of *Etendue* ($nA\Omega$)

□

$$\square \Rightarrow \theta s = \theta' s' \Rightarrow s' = s \frac{\theta}{\theta'} = s \frac{F_2}{F_1}$$



$$\Rightarrow \delta\lambda = \left(\frac{d\lambda}{dx}\right) s' = \frac{\cos \beta}{m\rho f_2} s \frac{F_2}{F_1} = \frac{s D_1 \cos \beta}{m\rho D_2 f_1}$$

Camera focal length

Resolving power

- Illuminated grating length: $W = \frac{D_2}{\cos \beta}$
- Spectral resolution (width) $\delta\lambda = \frac{s}{m\rho F_1 W}$
 - Collimator focal ratio
- Resolving power:
 - expressed in laboratory terms $R \equiv \frac{\lambda}{\delta\lambda} = \frac{m\rho\lambda F_1 W}{s}$
 - Physical slitwidth
 - expressed in astronomical terms $R = \frac{m\rho\lambda W}{\chi D_T}$
 - Grating length
 - Angular slitwidth
 - Telescope size

since $s = \chi f_T$ and $\frac{f_T}{D_T} = \frac{f_1}{D_1} = F_T = F_1$

Size of spectrograph must scale with telescope size

Importance of slit width

- Width of slit determines:
 - Resolving power (R) since $R\chi = \text{constant}$
 - Throughput (η)
- Hence there is always a tradeoff between *throughput* and *spectral information*
- Function $\eta(\chi)$ depends on *Point Spread Function* (PSF) and profile of extended source
 - generally $\eta(\chi)$ increases slower than χ^{+1} whereas $R \propto \chi^{-1}$
so ηR *maximised at small χ*
- Signal/noise also depends on slit width
 - throughput ($\rightarrow$ signal)
 - wider slit admits more sky background ($\rightarrow$ noise)

Blazing

- Diffracted intensity:

Interference pattern

$$I = \left(\frac{\sin^2 N\phi}{\sin^2 \phi} \right) \left(\frac{\sin^2 \theta}{\theta^2} \right)$$

Single slit diffraction

Φ = phase difference between adjacent rulings

θ = phase difference from centre of one ruling to its edge

→ Shift envelope peak to $m=1$

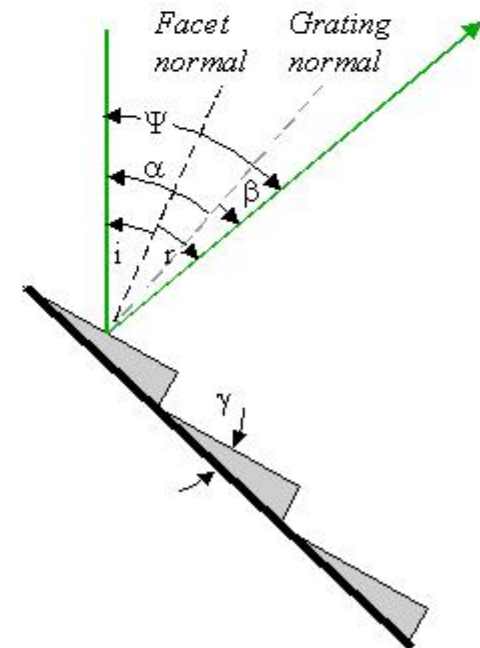
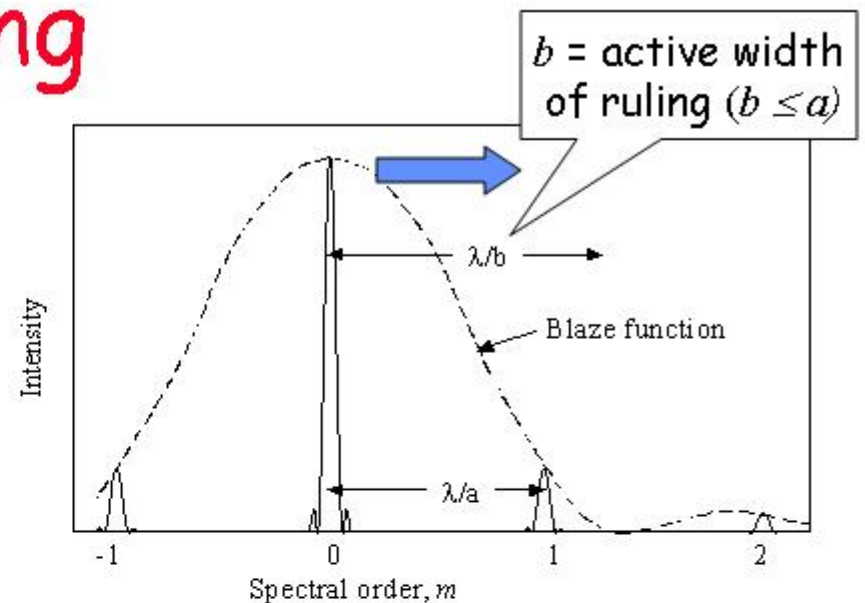
- Blaze condition

specular reflection off grooves:

$$\alpha + \beta = 2\gamma \quad \text{also} \quad \Psi = \alpha - \beta$$

$$\Rightarrow \rho m \lambda_B = \sin \alpha + \sin \beta = 2 \sin \gamma \cos \frac{\Psi}{2}$$

$$\sin x + \sin y = 2 \sin \frac{x+y}{2} \cos \frac{x-y}{2}$$



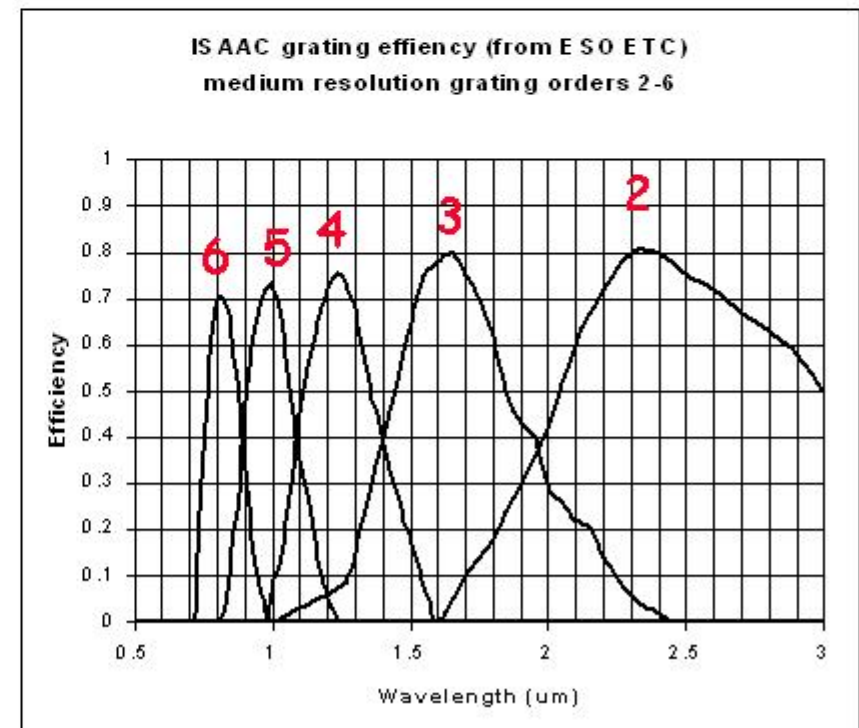
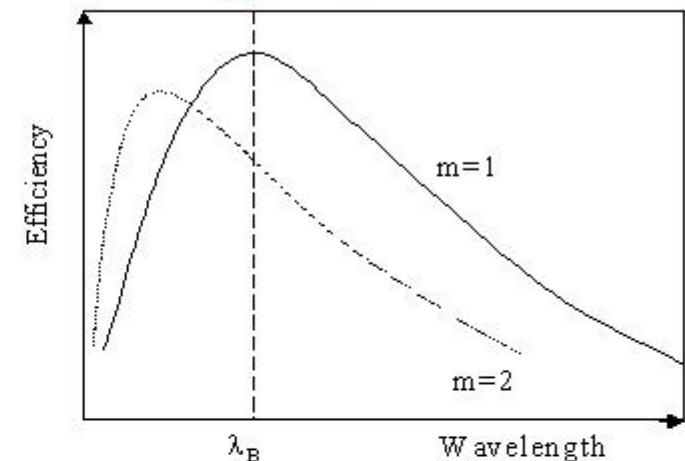
Efficiency vs wavelength

- Approximation valid for $a > \lambda$
- $\lambda_{max}(m) = \lambda_B(m=1)/m$
- Rule-of-thumb:
40.5% x peak at

$$\lambda_+ = \frac{2m\lambda_B}{2m-1} \text{ and } \lambda_- = \frac{2m\lambda_B}{2m+1}$$

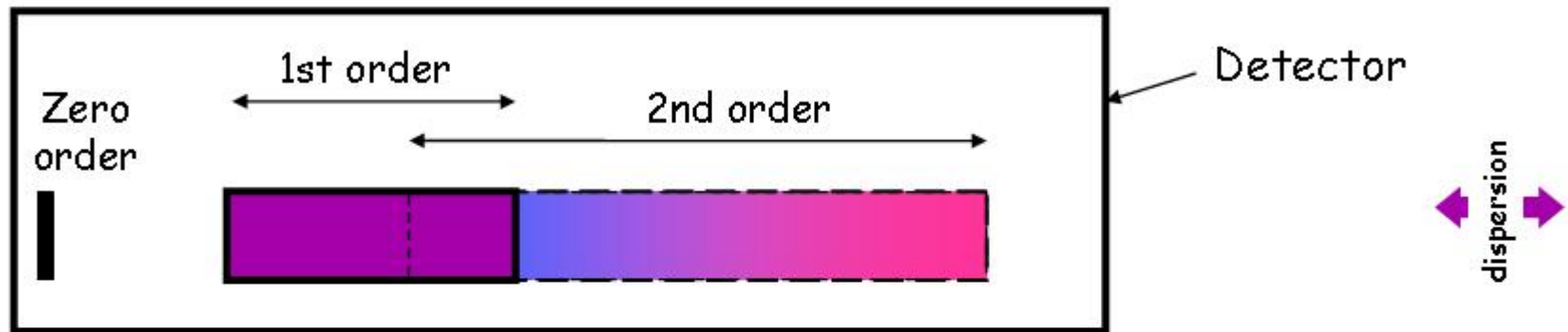
$$\Rightarrow \lambda_+ - \lambda_- \approx \frac{\lambda_B}{m} \text{ (large } m)$$

- Sum over all orders < 1
 - reduction in efficiency with increasing order



(See: Schroeder, *Astronomical Optics*)

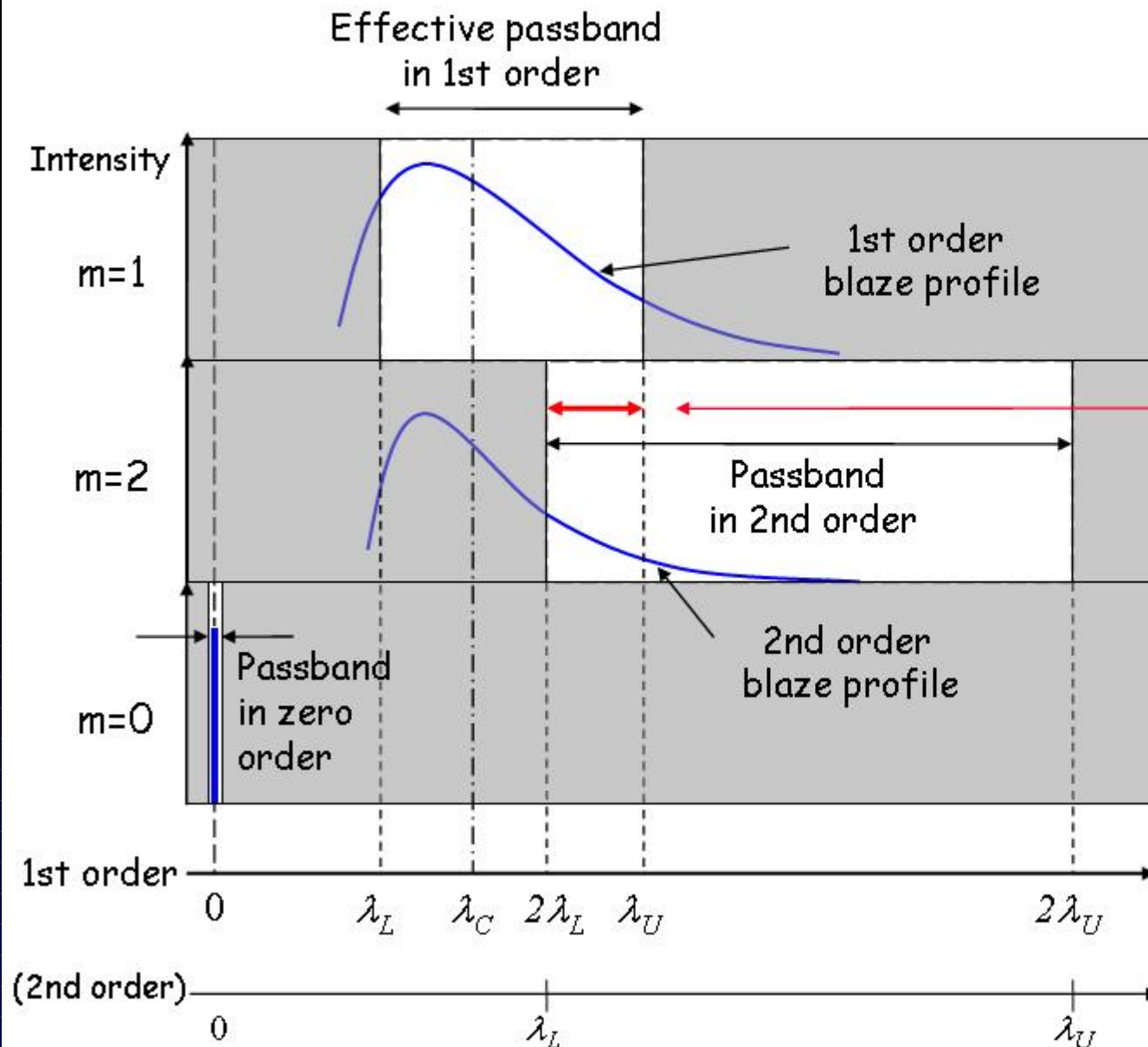
Order overlaps



To eliminate overlap between 1st and 2nd order

- Limit wavelength range incident on detector using *passband* filter or *longpass* ("order rejection") filter acting with long-wavelength cutoff of optics or detector (e.g. 1100nm for CCD)
- Optimum wavelength range is 1 octave (then $2\lambda_L = \lambda_U$)
- *Zero order may be a problem in multiobject spectroscopy*

Order overlaps



Don't forget
higher orders!

First and second
orders overlap!

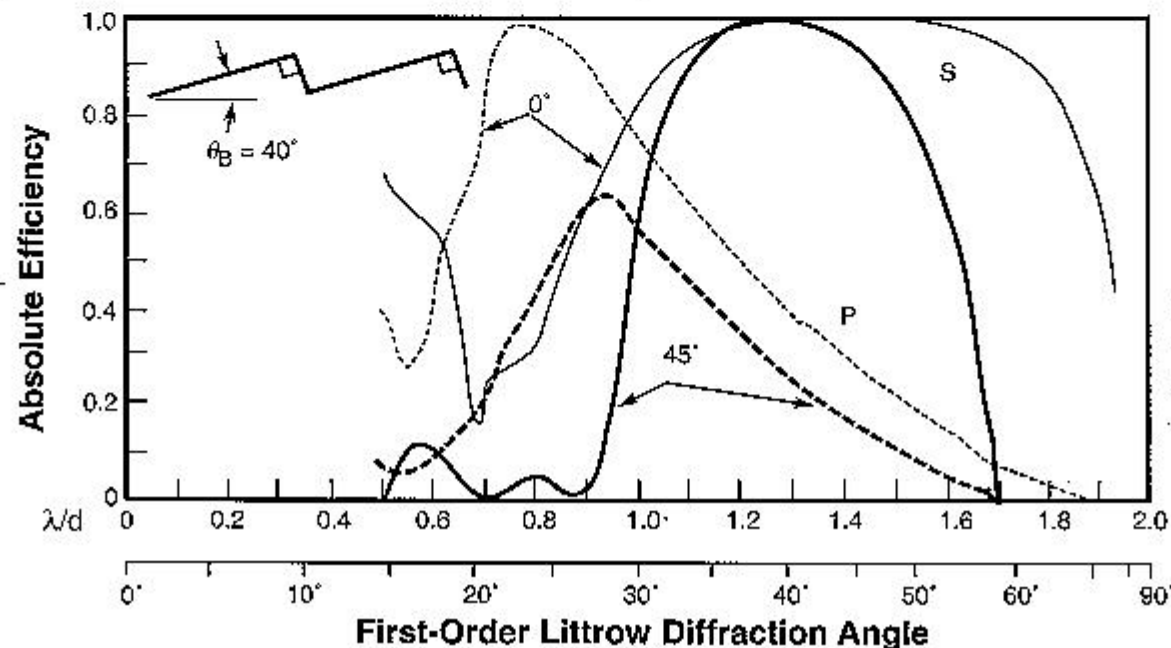
Zero order
matters
for MOS

Wavelength in first
order marking
position on detector
in dispersion direction
(if dispersion $\sim$ linear)

Efficiency - semi-empirical (contd)

Different regimes for blazed (triangular) grooves

$\gamma < 5^\circ$	obeys scalar theory, little polarisation effect ($P \approx S$)
$5 < \gamma < 10^\circ$	S anomaly at $\rho\lambda \approx 2/3$, P peaks at lower $\rho\lambda$ than S
$10 < \gamma < 18^\circ$	various S anomalies
$18 < \gamma < 22^\circ$	anomalies suppressed, $S \gg P$ at large $\rho\lambda$
$22 < \gamma < 38^\circ$	strong S anomaly at P peak, S constant at large $\rho\lambda$
$\gamma > 38^\circ$	S and P peaks very different, efficient in Littrow only



NOTE
Results apply to
Littrow only

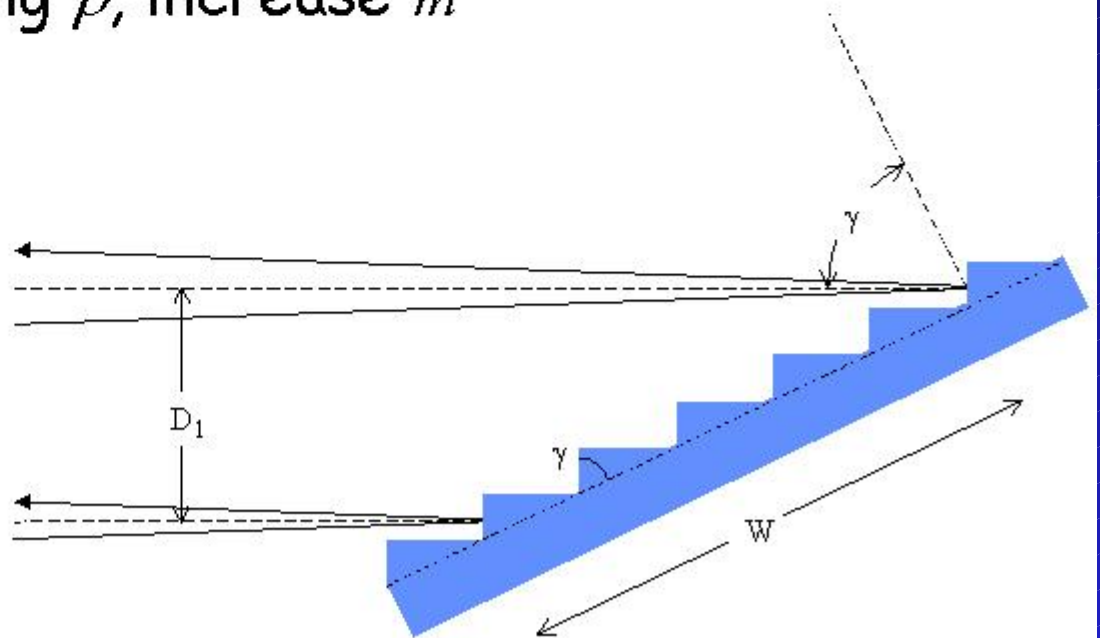
From: *Diffraction Grating Handbook*,
C. Palmer, Thermo RGL,
(www.gratinglab.com)

$\rho\lambda \longrightarrow$

$\alpha=\beta \longrightarrow$

Echelle gratings

- Obtain very high R ($> 10^5$) using very long grating
- In Littrow: $R = \frac{m\rho\lambda W}{\chi D_T} = \frac{2D_1}{\chi D_T} \tan \gamma$ Groove angle
- Maximising γ requires large $m\rho$ since $m\rho\lambda = 2\sin\gamma$
- Instead of increasing ρ , increase m
- *Echelle* is a coarse grating with large groove angle
- R parameter = $\tan\gamma$
(e.g $R2 \Rightarrow \gamma = 63.5^\circ$)



Multiple orders

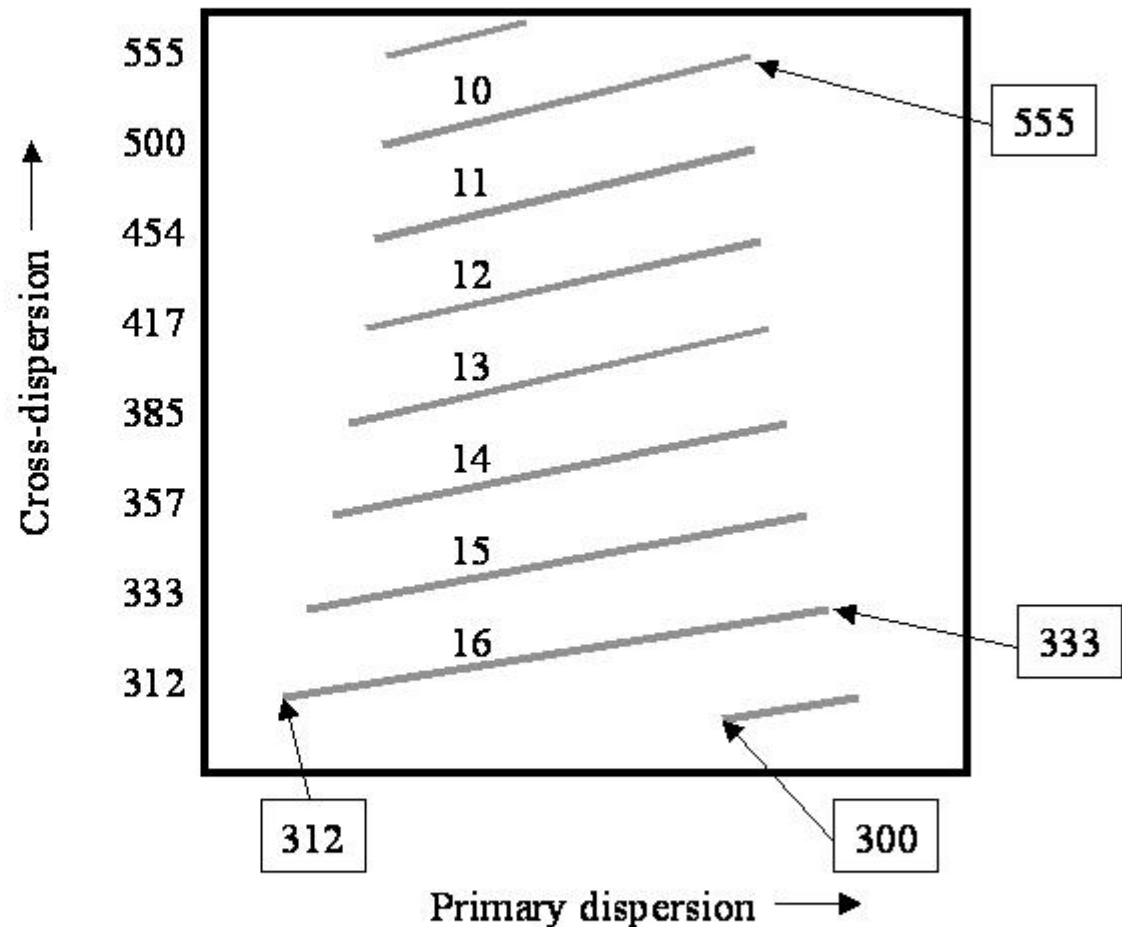
- Many orders to cover desired $\lambda\lambda$:
Free spectral range

$$\Delta\lambda = \lambda/m$$

- Orders lie on top of each other:

$$\lambda(m) = \lambda(n) \times (n/m)$$

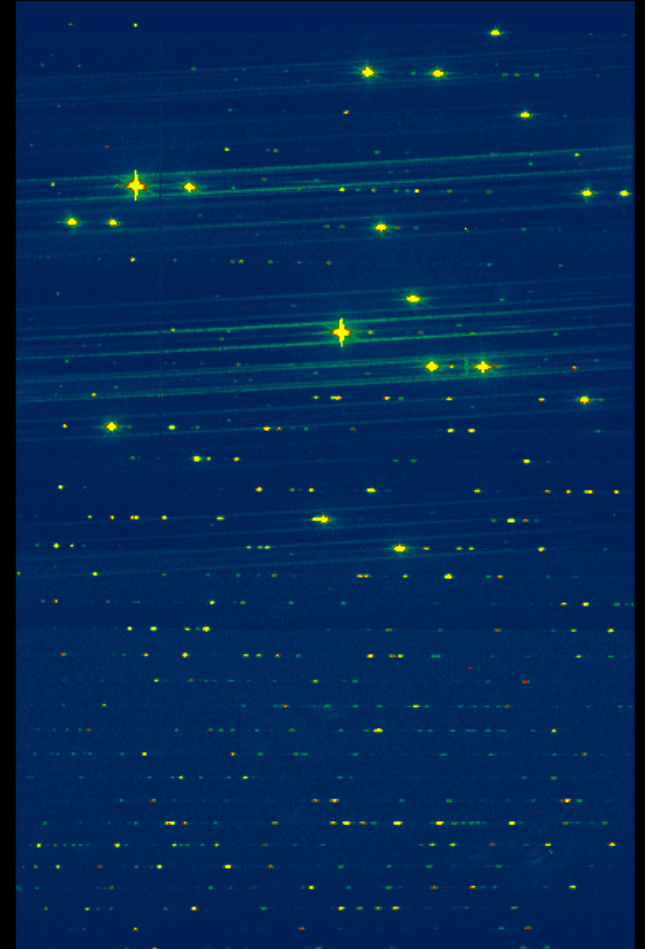
- Solution:
 - use narrow passband filter to isolate one order at a time
 - cross-disperse to fill detector with many orders at once



Cross dispersion may use prisms
or low dispersion grating

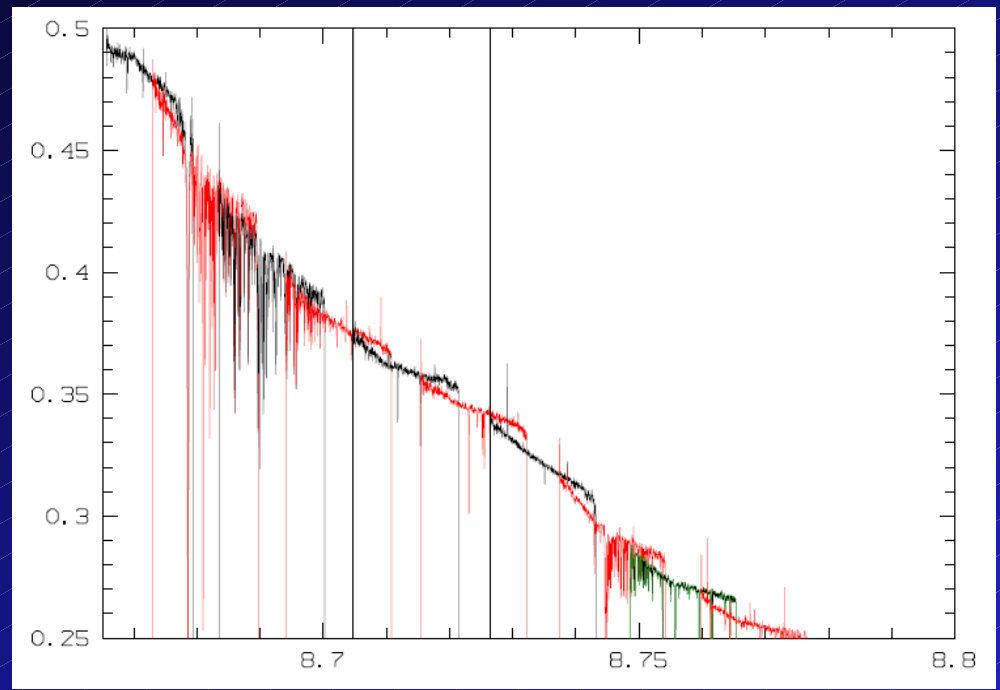
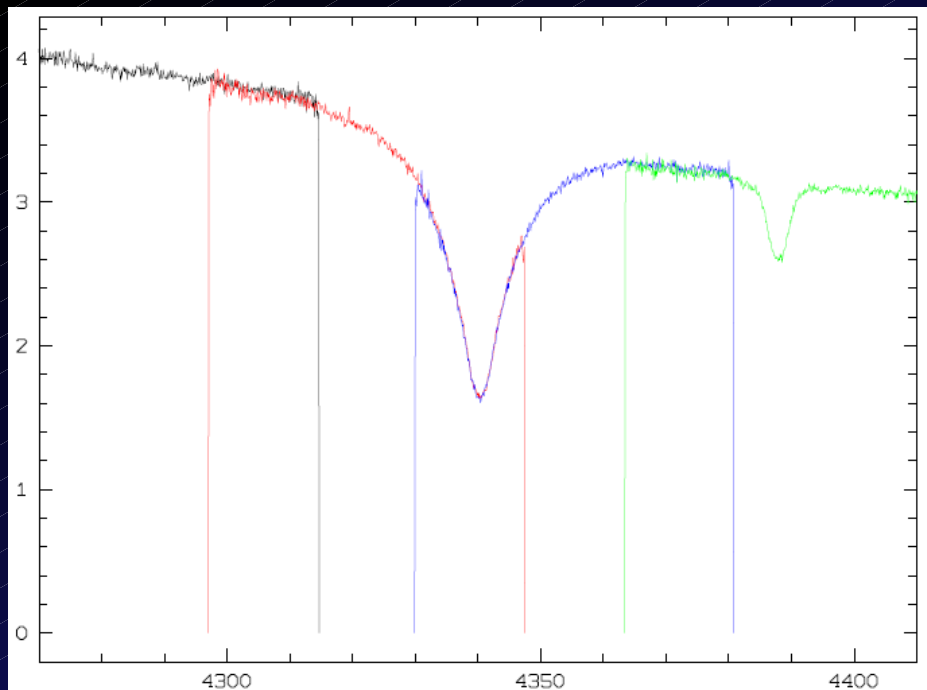
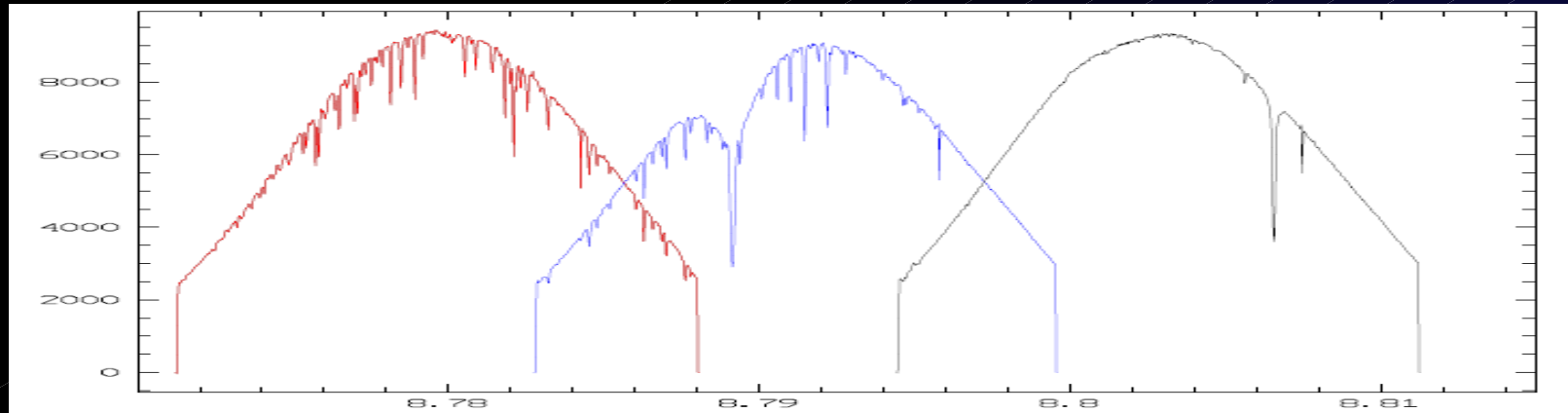


CHANL: 0 START: 2.0,1.0
FRAME: ra5921.at1 END: 750.0,1179.0
CUTS: 0.0,1700.0 MIN,MAX: 0.0,0.0
CURSO: 261, 215, 254

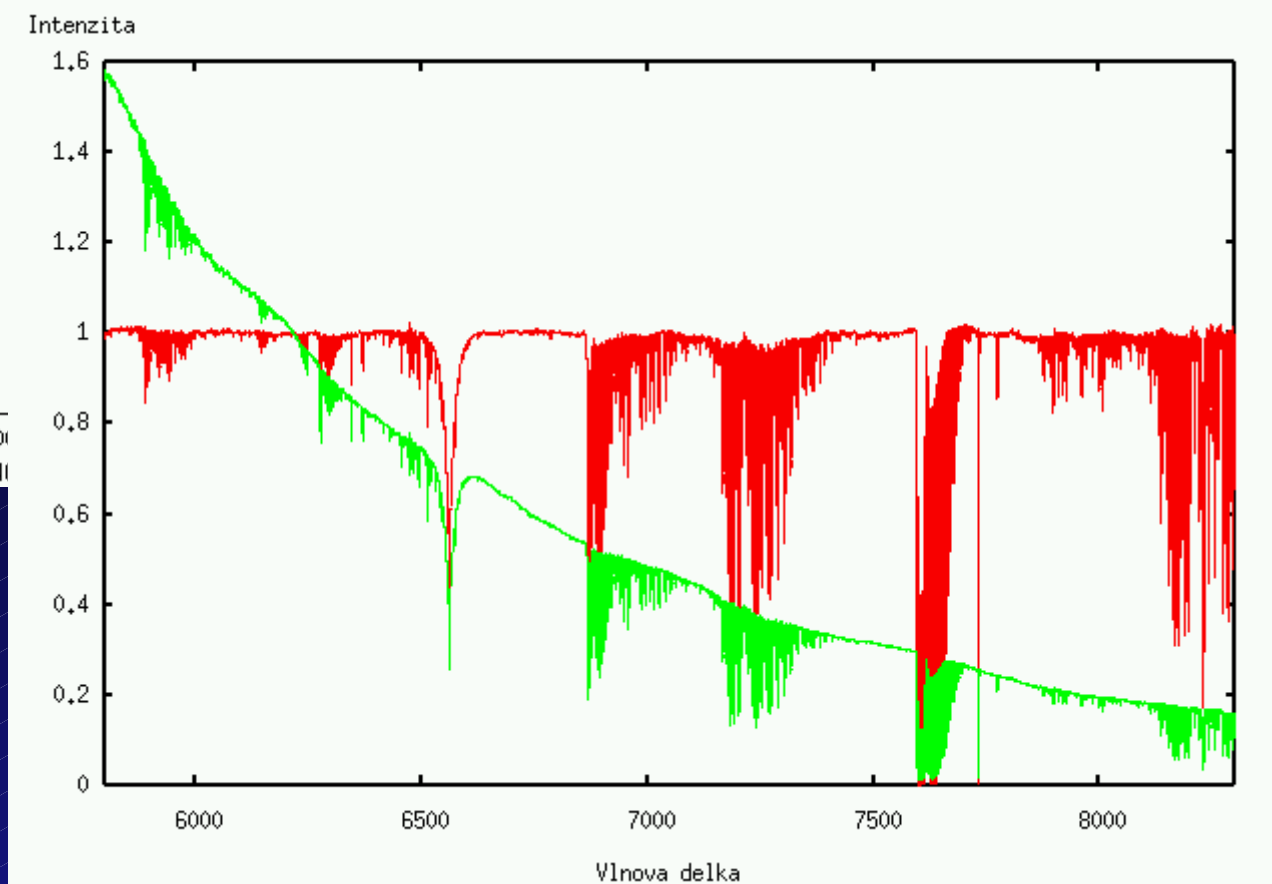
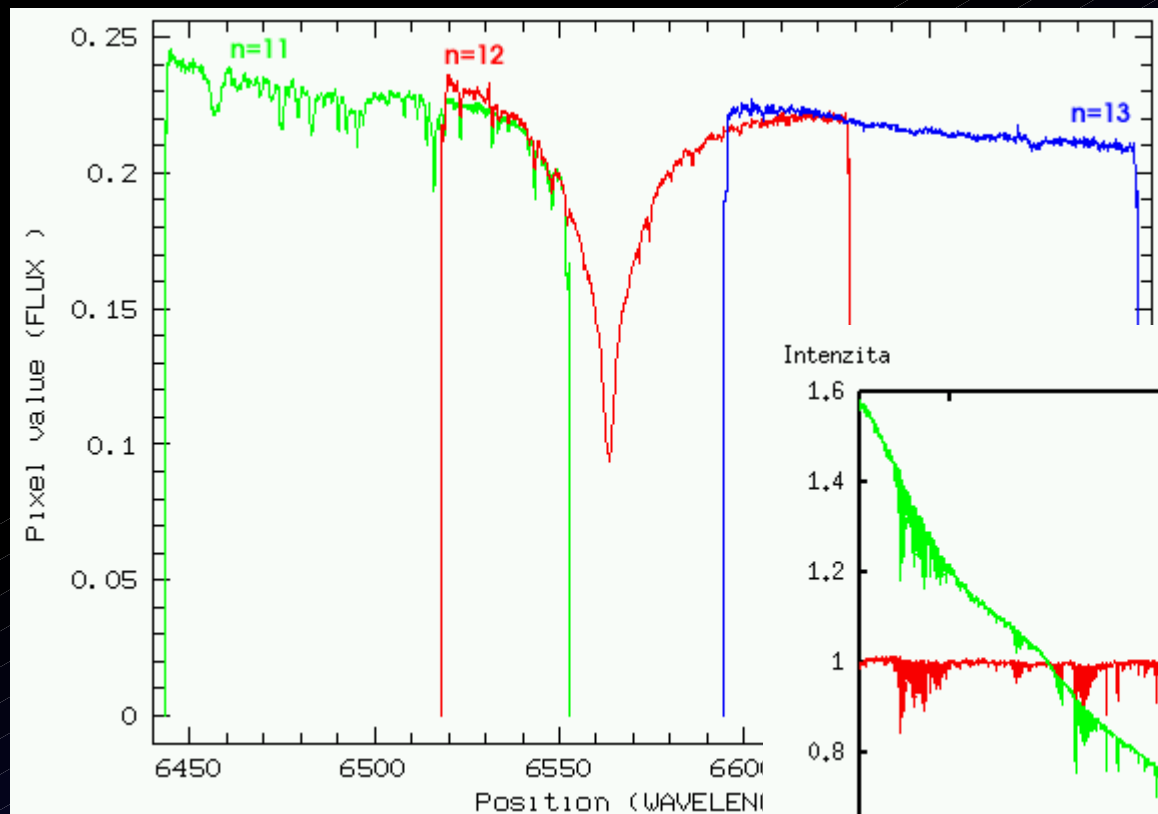


CHANL: 0 START: 2.0,1.0
FRAME: ra5915.at1 END: 750.0,1179.0
CUTS: 0.0,120.0 MIN,MAX: 0.0,0.0
CURSO: 261, 215, 254

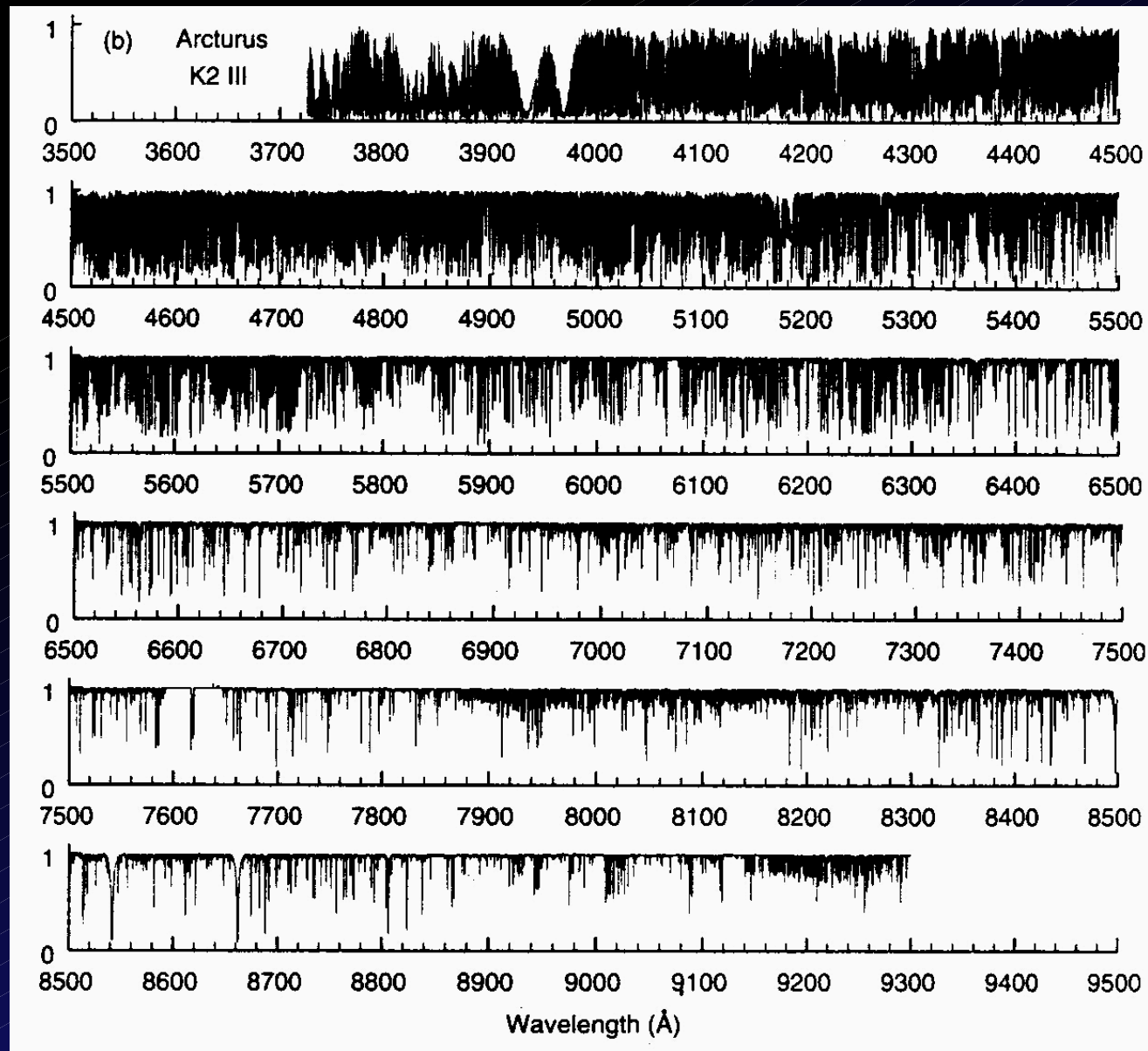
Echelle Spectra Problems



Echelle – spojování řádů



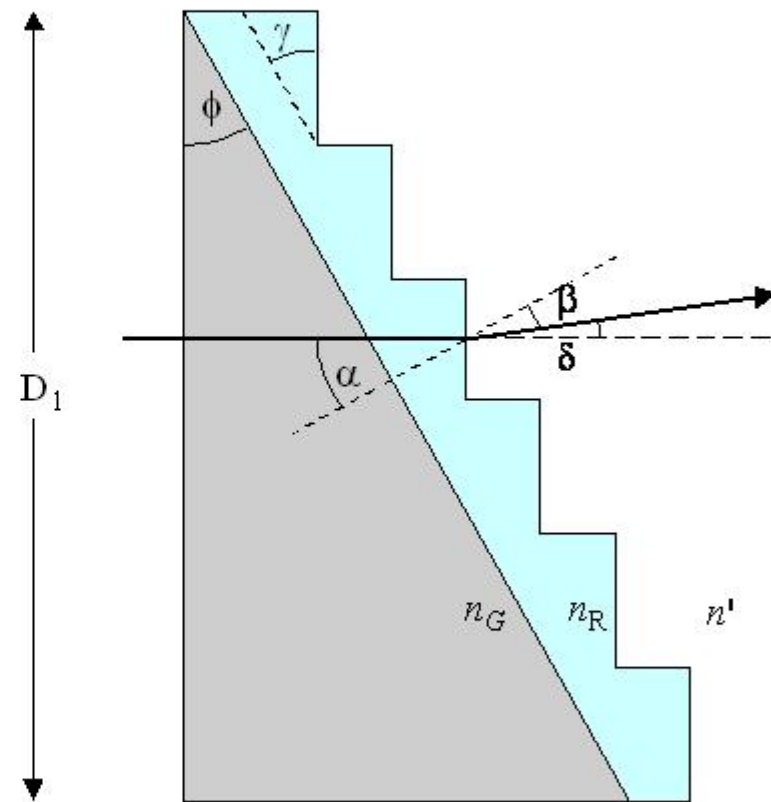
Merged Echelle



Arcturus: Hinkle et al.

Grisms

- Transmission grating attached to prism
- Allows in-line optical train:
 - simpler to engineer
 - quasi-Littrow configuration - no variable anamorphism
- Inefficient for $\rho > 600/\text{mm}$ due to groove shadowing and other effects

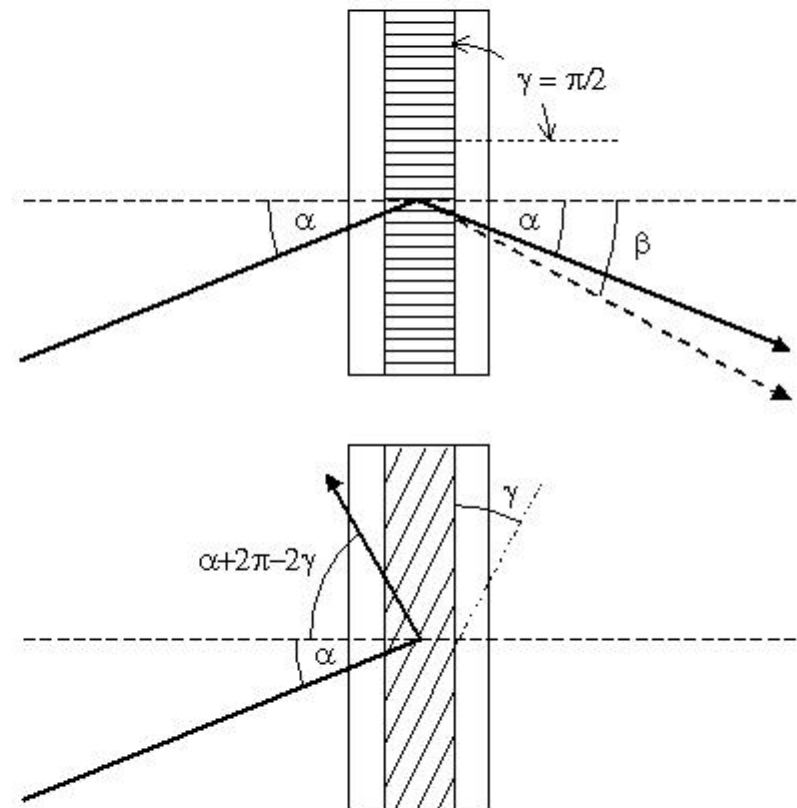


Volume Phase Holographic gratings

- So far we have considered *surface relief* gratings
- An alternative is *VPH* in which refractive index varies harmonically throughout the body of the grating: $n_g(x, z) = n_g + \Delta n_g \cos[2\pi\rho_g(x \sin \gamma + z \cos \gamma)]$
- Don't confuse with '*holographic*' gratings (SR)
- Advantages:
 - Higher peak efficiency than SR
 - Possibility of very large size with high ρ
 - Blaze condition can be altered (*tuned*)
 - Encapsulation in flat glass makes more robust
- Disadvantages
 - Tuning of blaze requires *bendable spectrograph!*
 - Issues of wavefront errors and cryogenic use

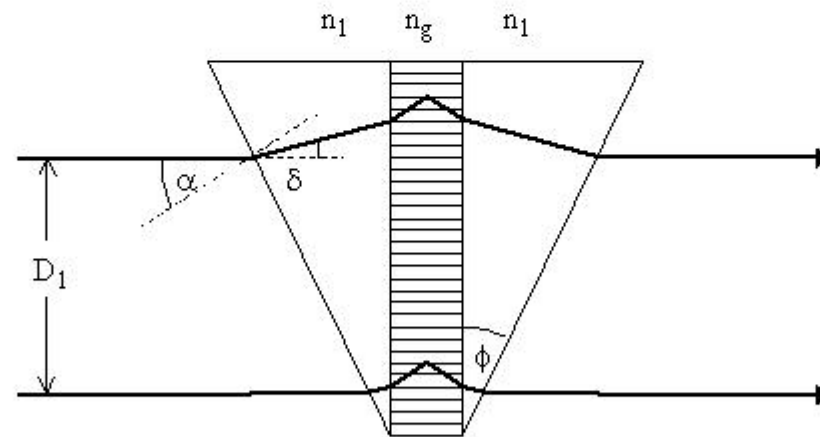
VPH configurations

- *Fringes* = planes of constant n
- Body of grating made from *Dichromated Gelatine* (DCG) which permanently adopts fringe pattern generated holographically
- Fringe orientation allows operation in transmission or reflection



VPH 'grism' = *vrism*

- Remove bent geometry, allow in-line optical layout
- Use prisms to bend input and output beams while generating required Bragg condition



$$R = \frac{m\rho\lambda W}{\chi D_T} = \frac{m\rho\lambda}{\chi D_T} D_1 (1 + \tan \delta \tan \varphi)$$

$$\delta = \varphi - \arcsin\left(\frac{\sin \varphi}{n_1}\right)$$

Limits to resolving power

- Resolving power can increase as m , ρ and W increase for a given wavelength, slit and telescope

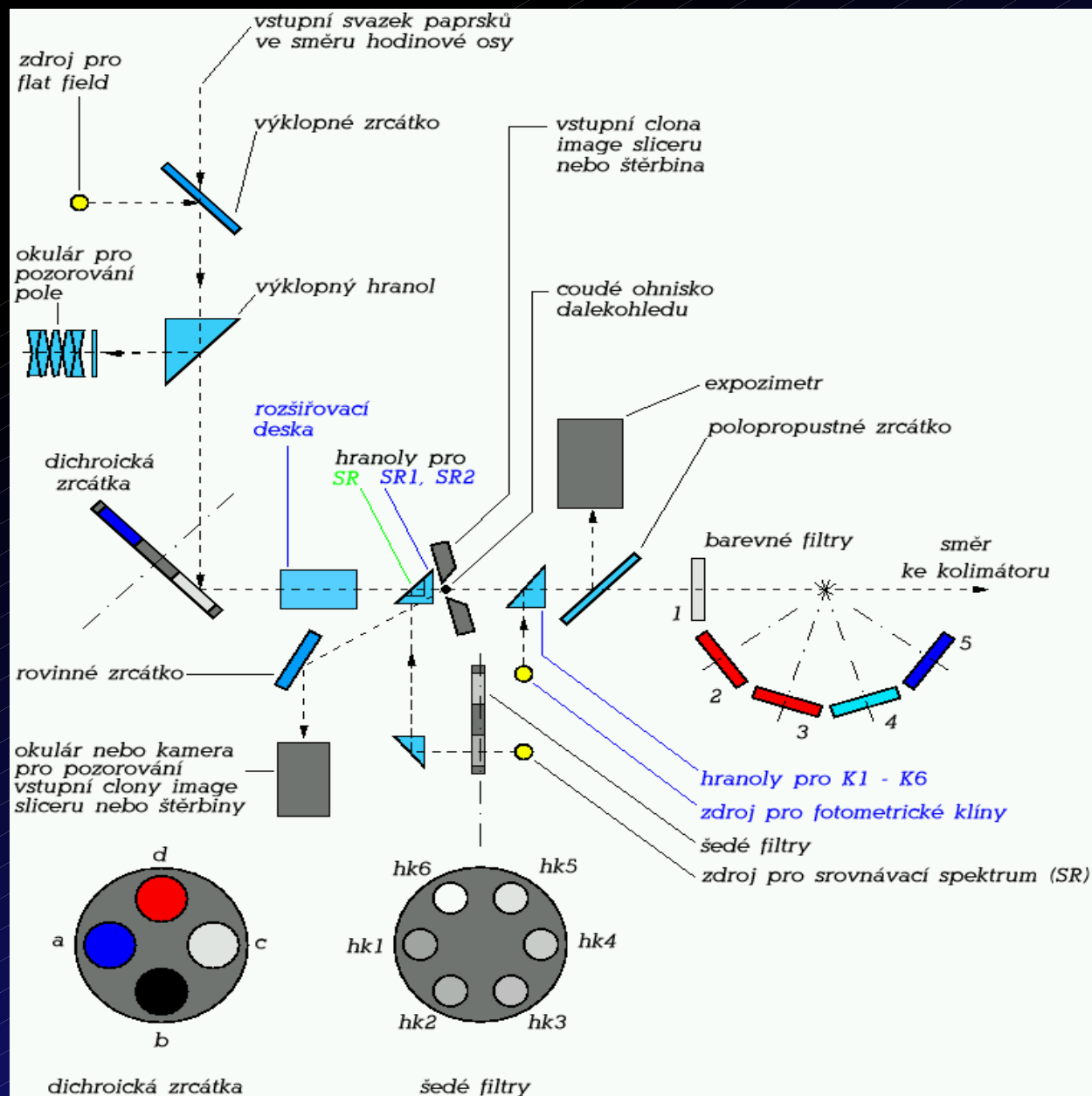
Grating
parameters

$$R = \frac{m\rho\lambda W}{\chi D_T} = \frac{D_1}{\chi D_T} \left(\frac{\sin \alpha + \sin \beta}{\cos \beta} \right)$$

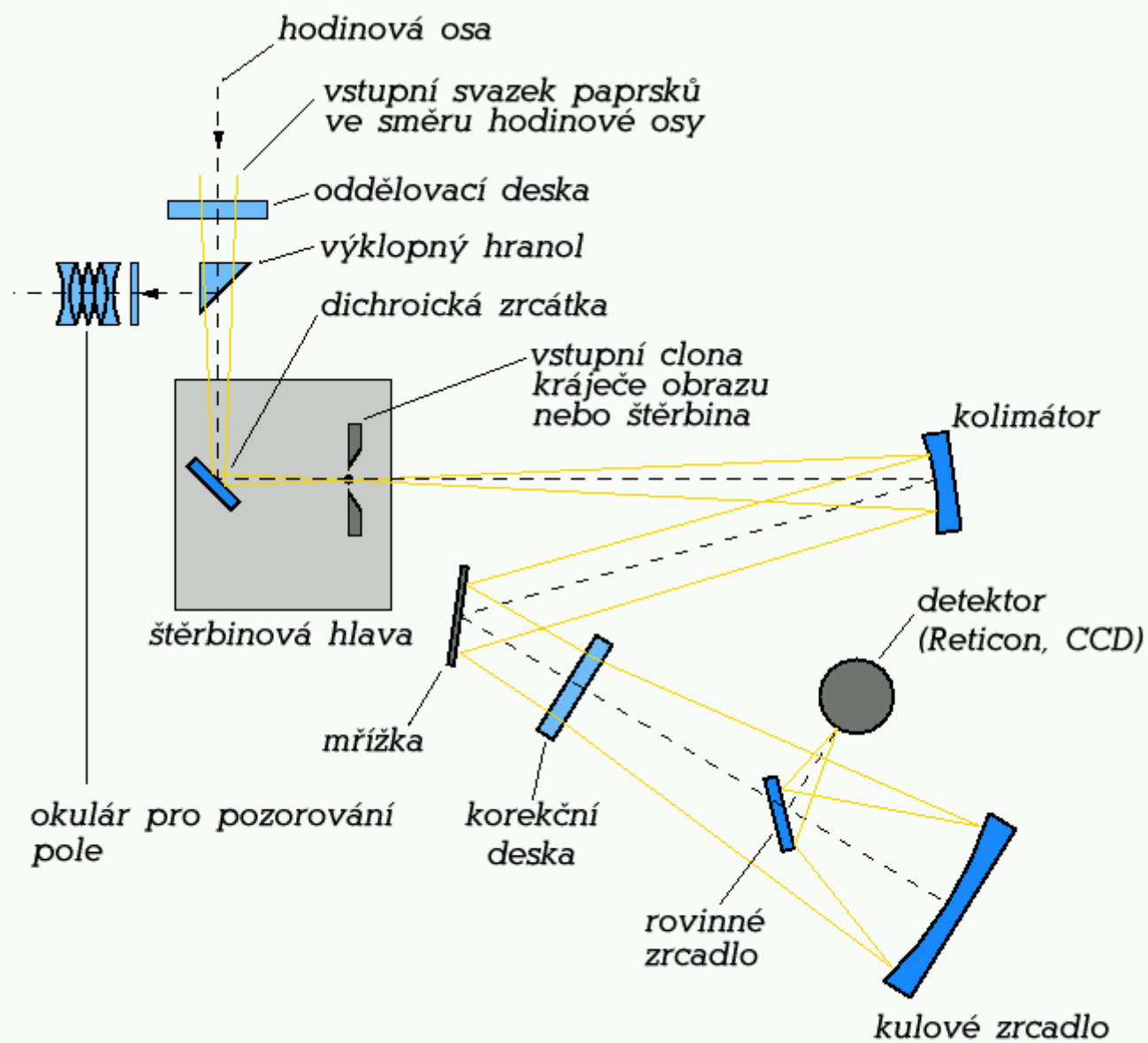
Geometrical
factors

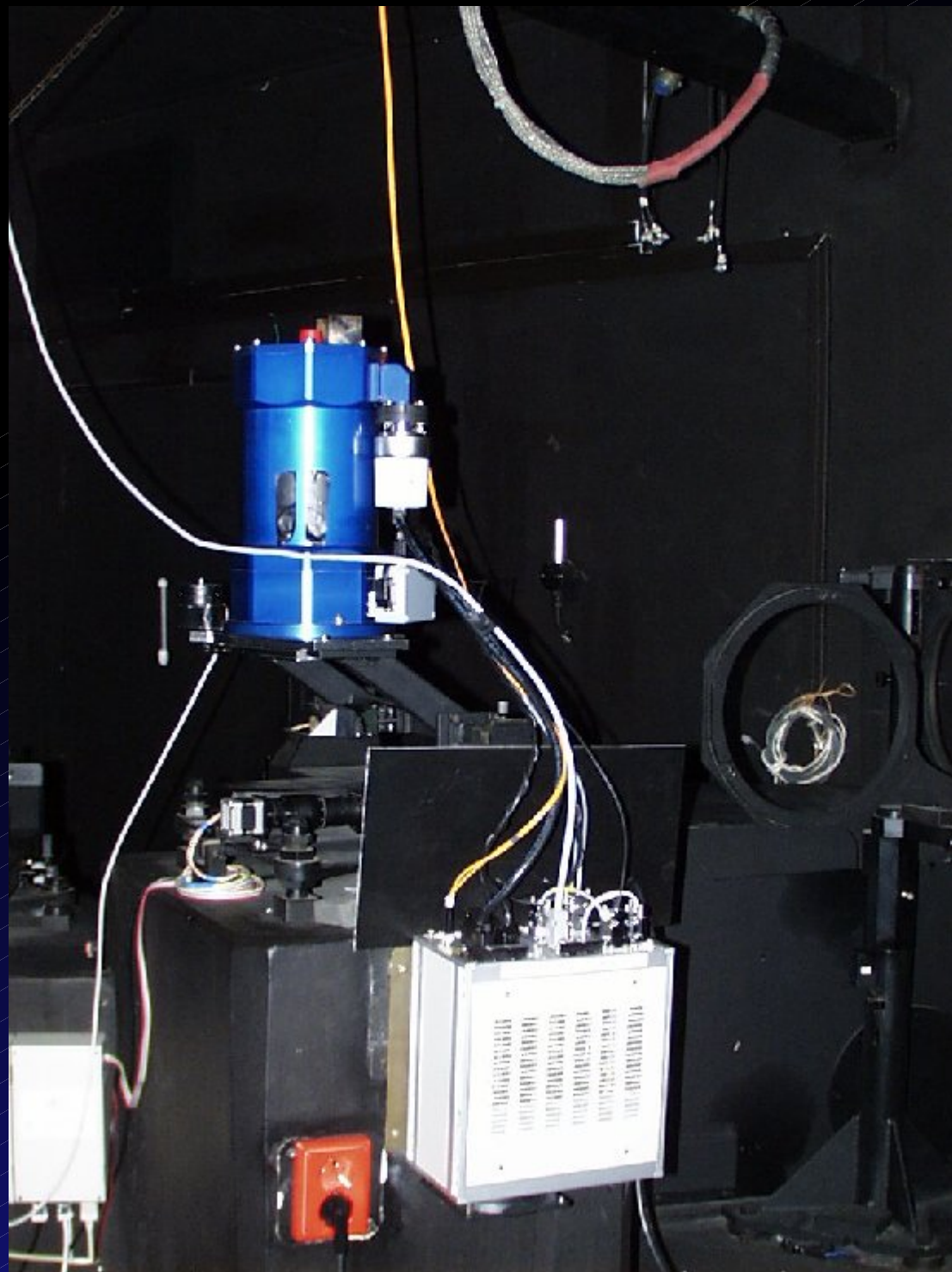
- Limit depends on geometrical factors only - increasing ρ or m will not help!
- In practice, the limit is when the output beam overfills the camera:
 - W is actually the length of the intersection between beam and grating plane - not the actual grating length
 - R will increase even if grating overfilled until diffraction-limited regime is entered ($\lambda > \chi D_T$)

Štěrbínová hlava spektrografu 2m

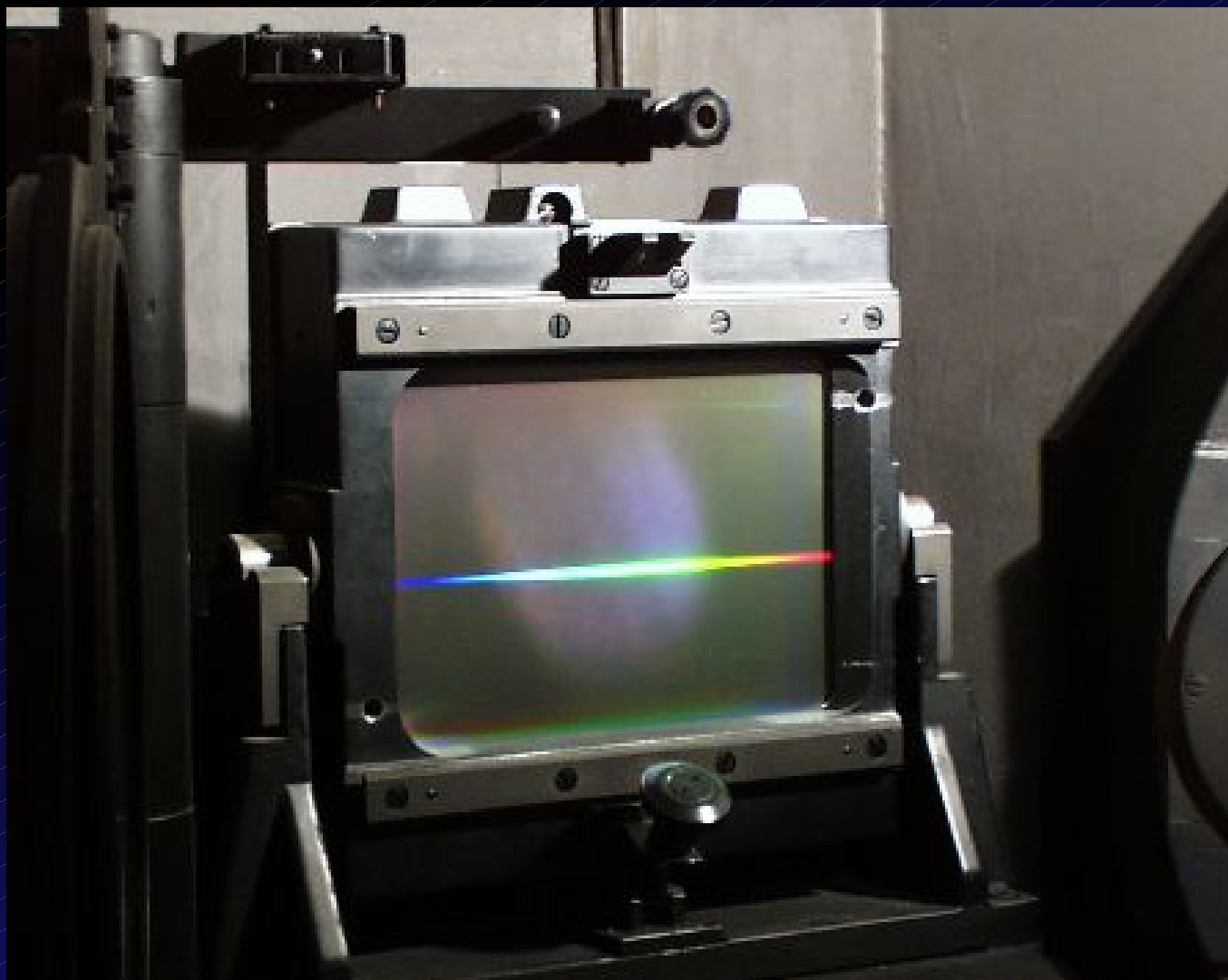


Coudé spektroskop 2m

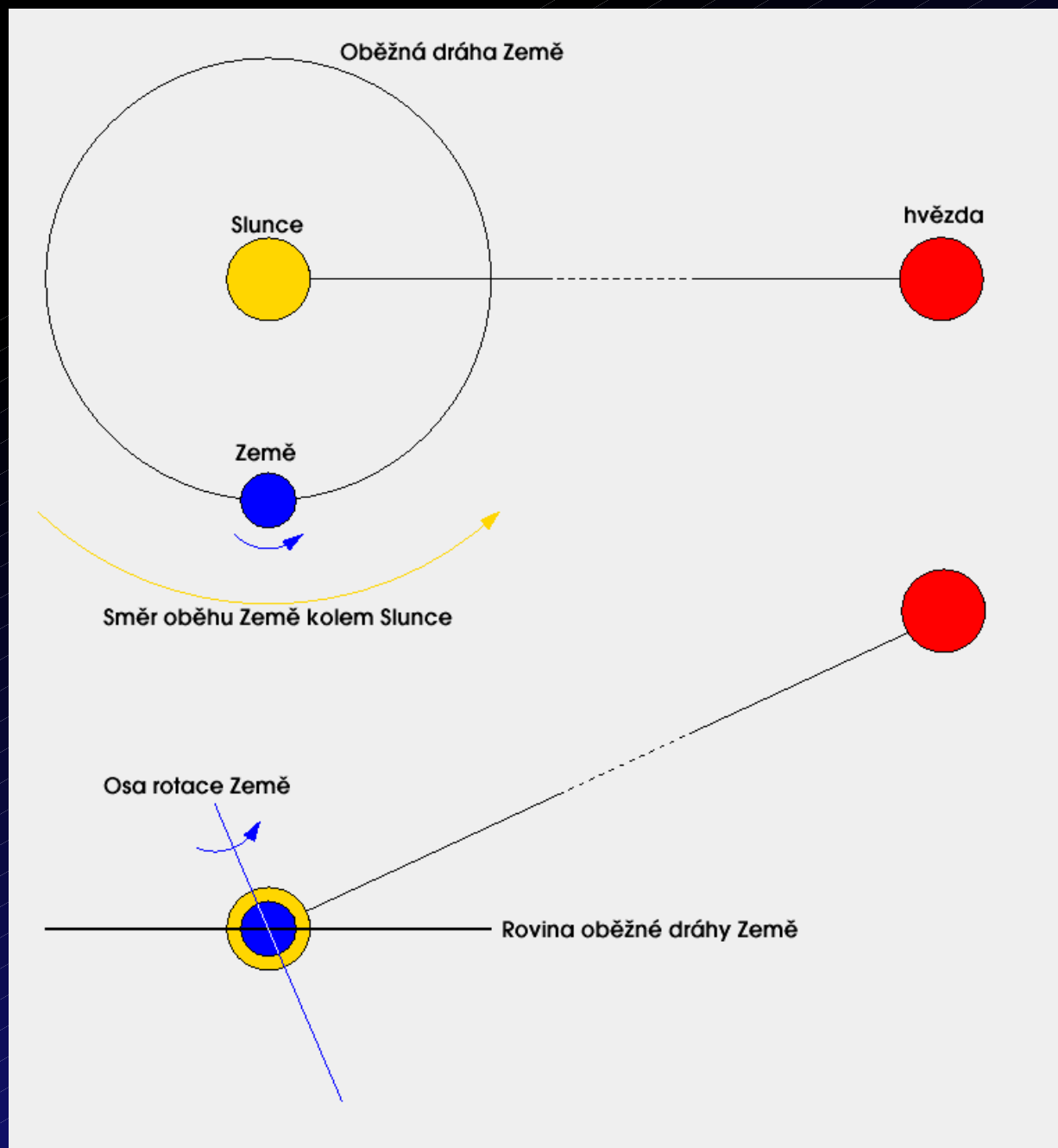




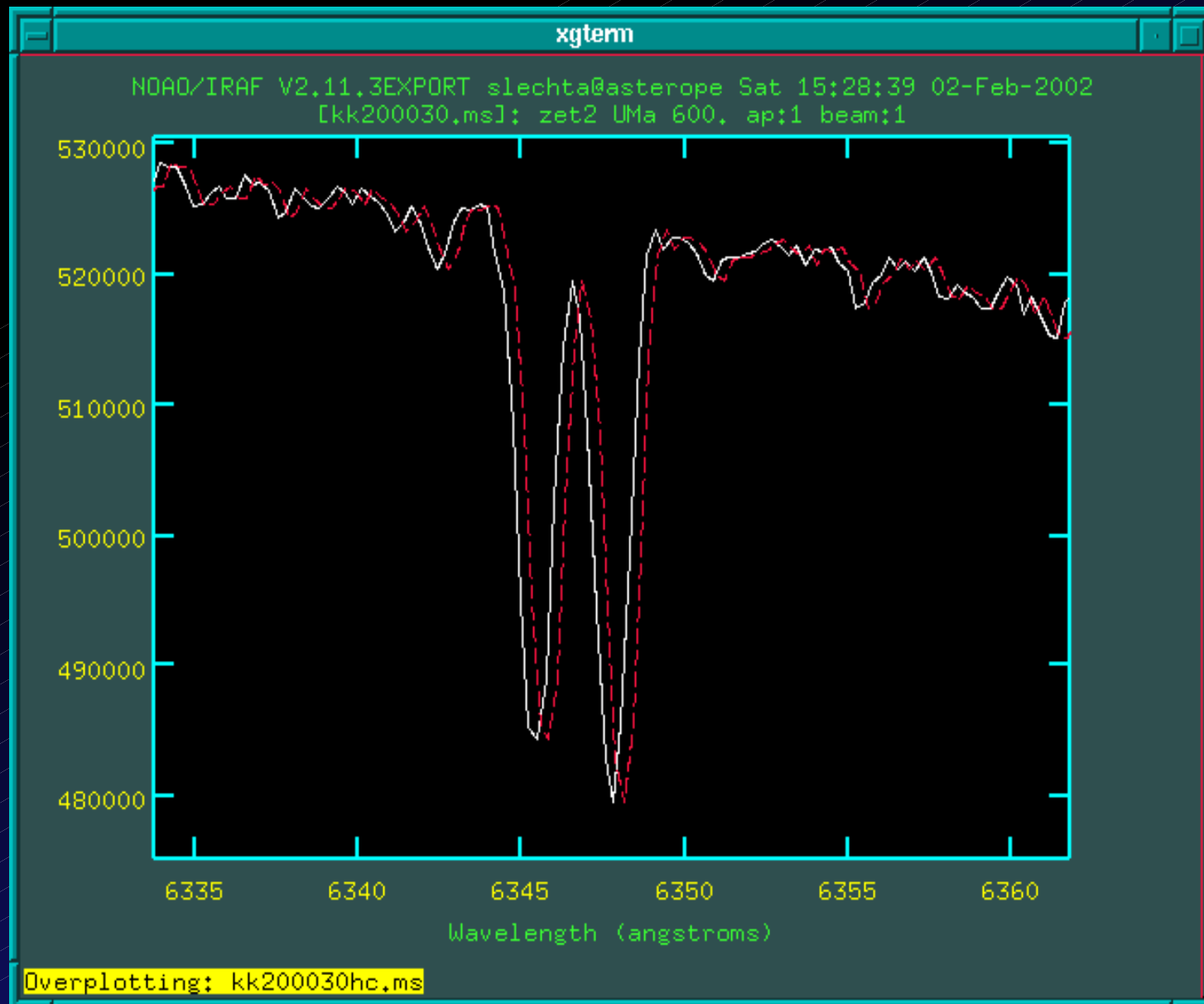
Mřížka 833/mm



Heliocentrická korekce



Posun čar díky heliocentrické korekci



Redukce – vyčtení COP

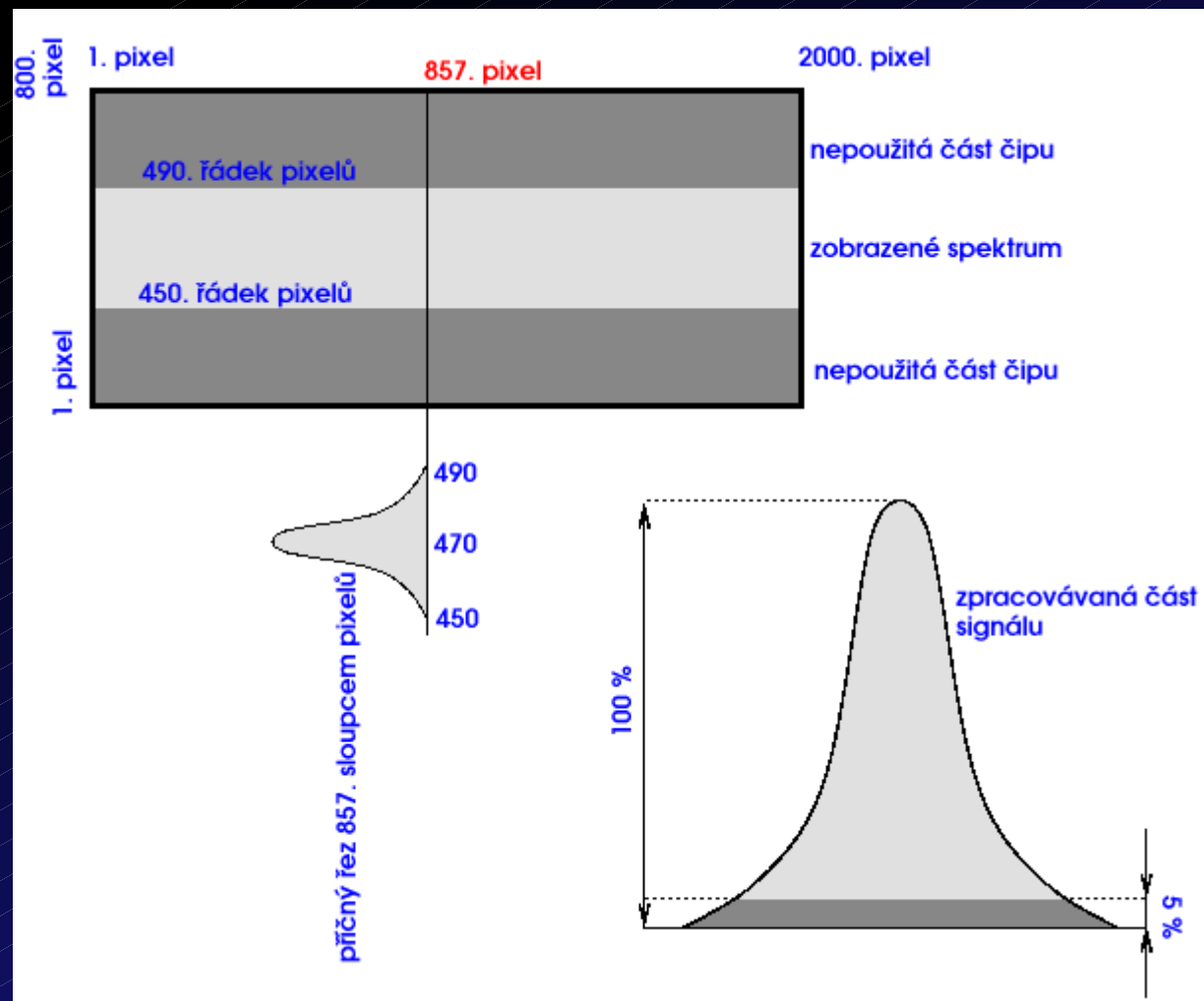
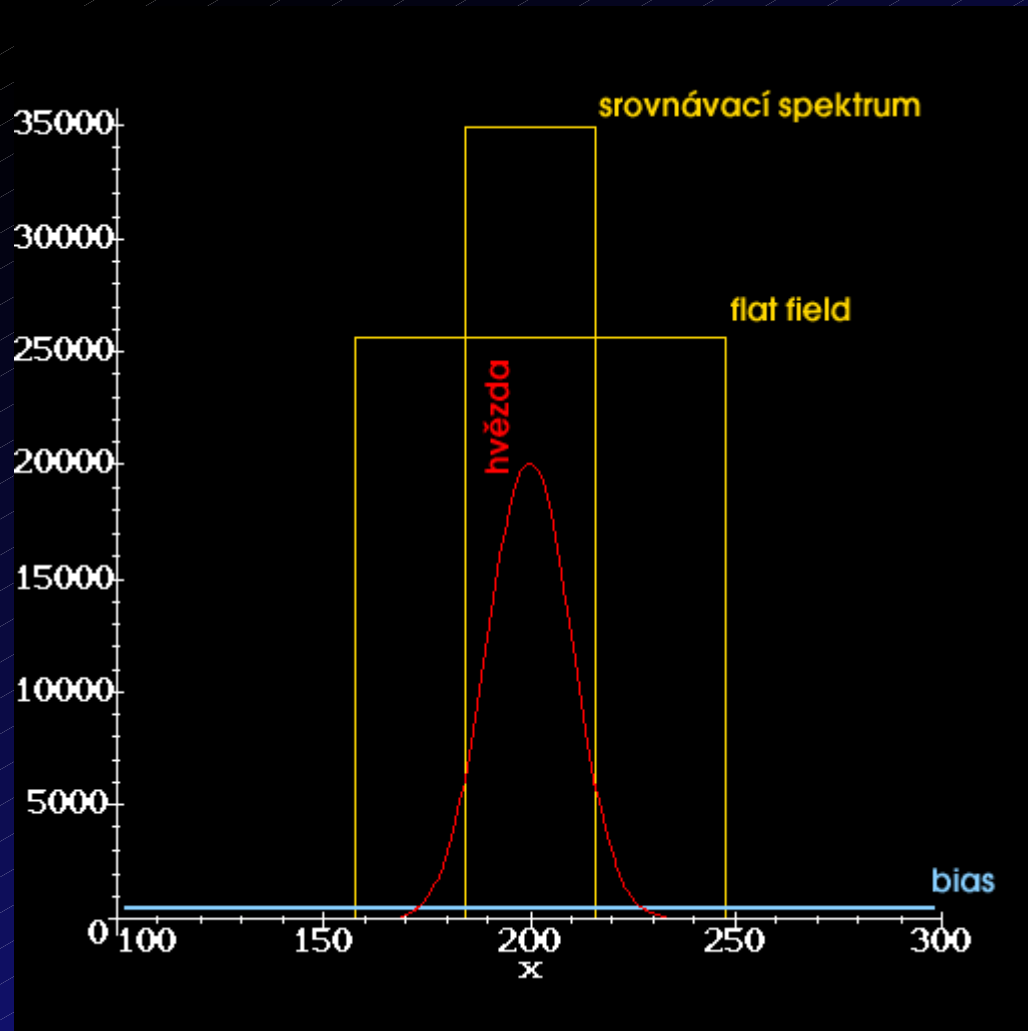
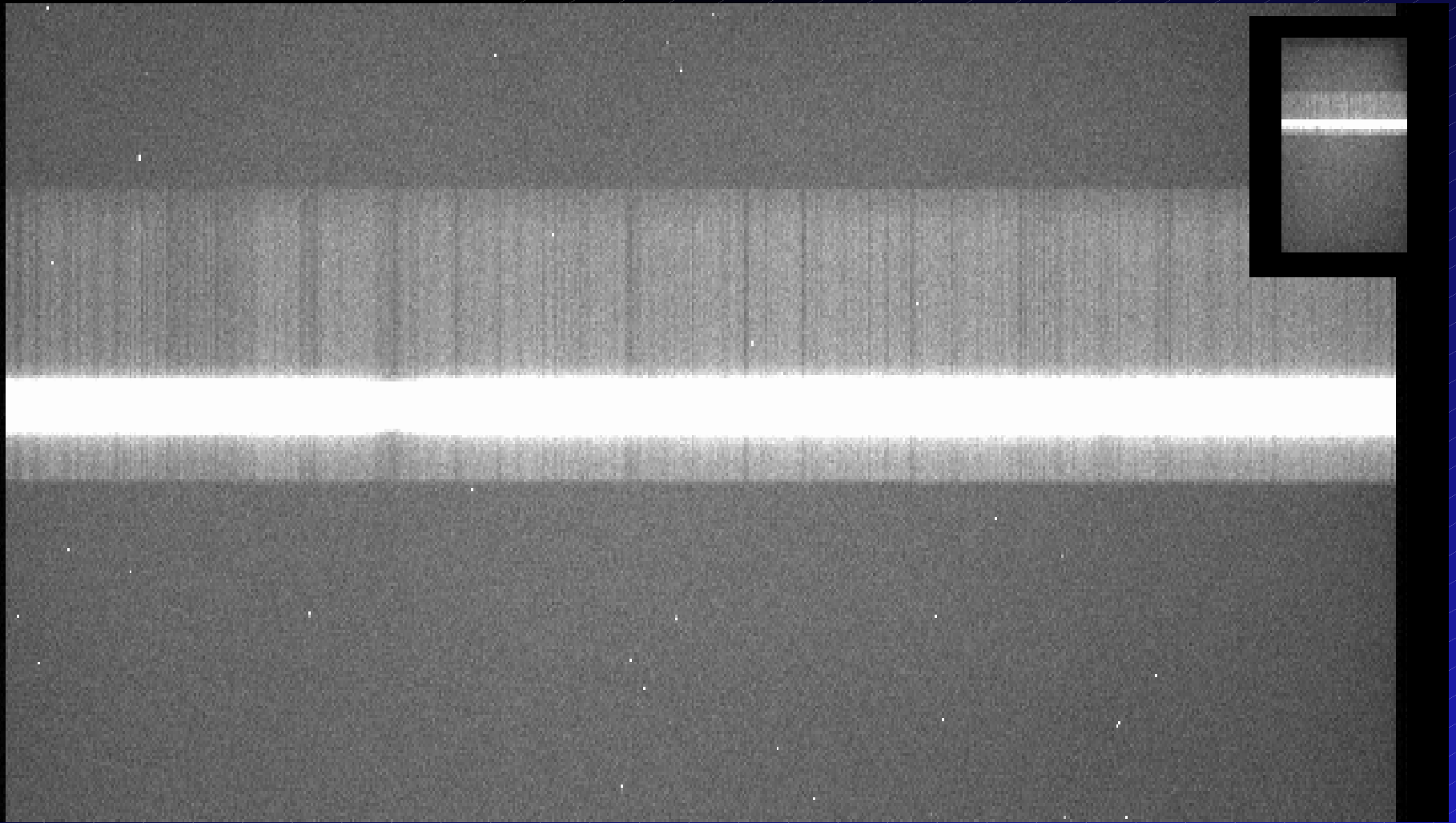


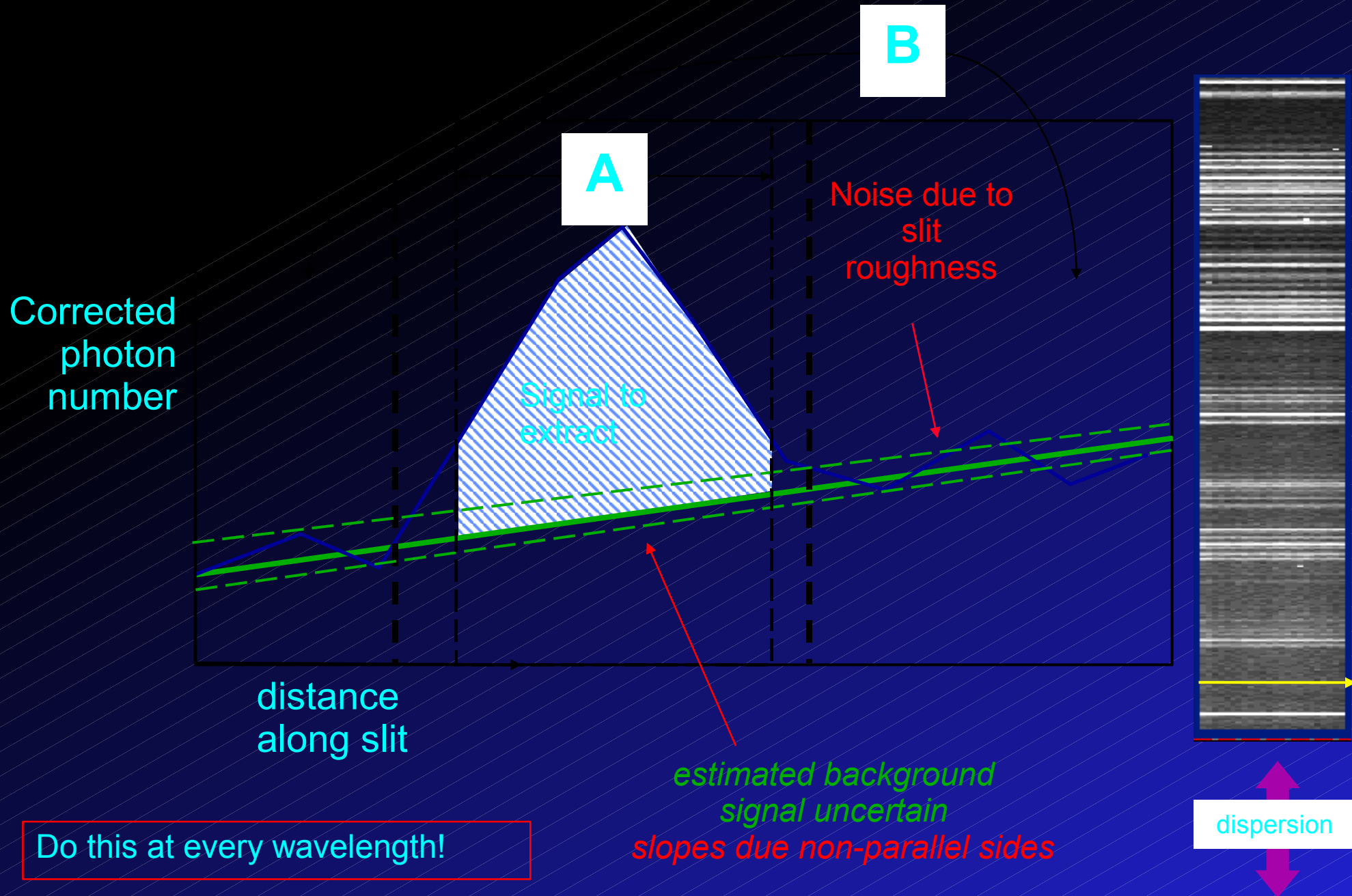
Schéma COP



Background a sky lines



Sky subtraction with slit



Sky subtraction near bright sky lines

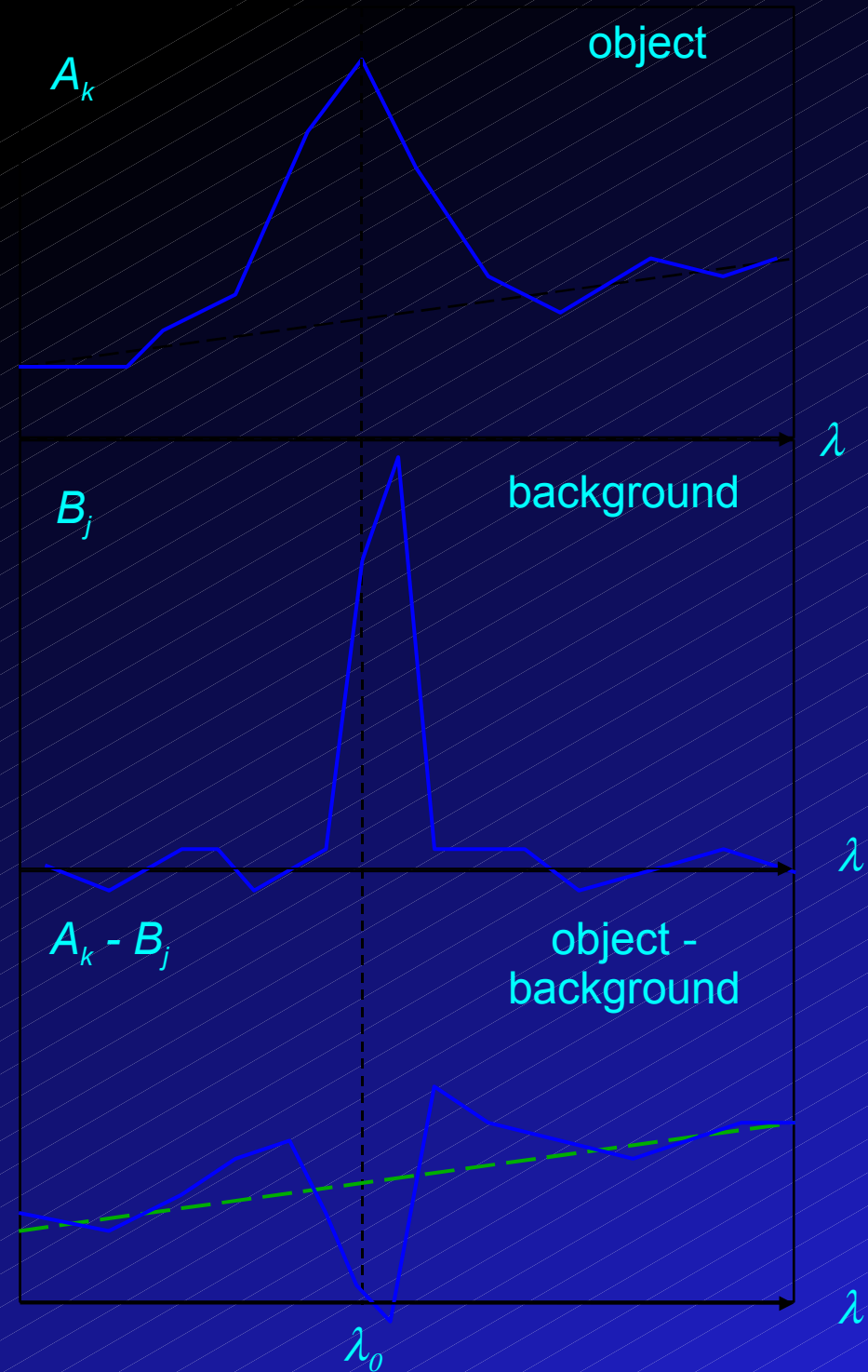
Poor cancellation of
sky line due to:

Difference in line
profile due to:

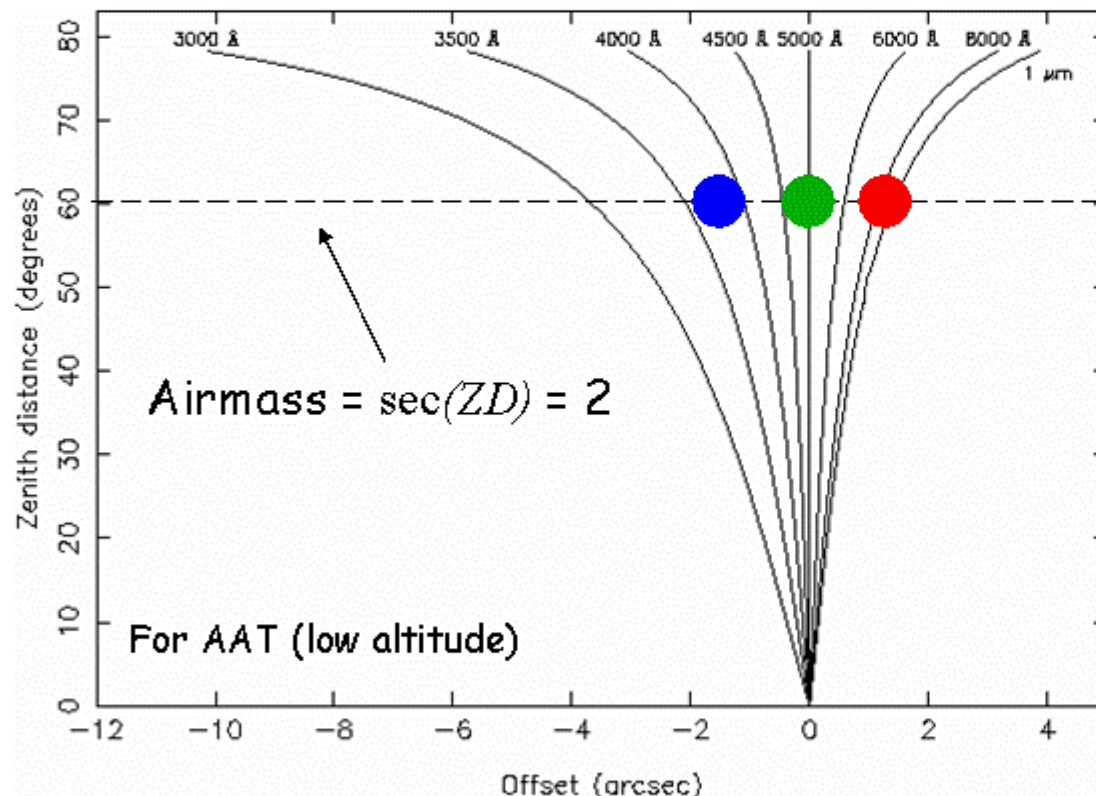
- uneven slit width
- IQ varies over field

Difference in line
location due to:

- tilt of slit
- poor wavelength
calibration/ solution/



Atmospheric dispersion (differential refraction)

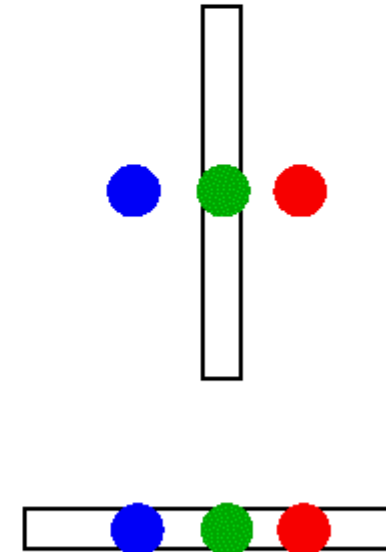


← Error in position
if guiding at 500nm →

$\propto \tan(ZD)$

To zenith

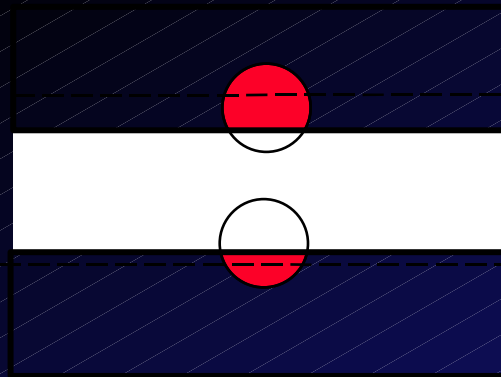
If slit is horizontal, light is lost at extreme wavelength



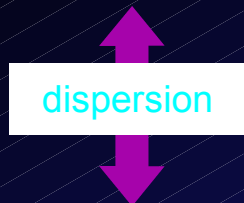
If slit is vertical, all light is in slit, but spectrum will be curved

Errors in centroid of VRE

VRE = *velocity resolution element*,
the monochromatic image of the slit as
recorded by the detector

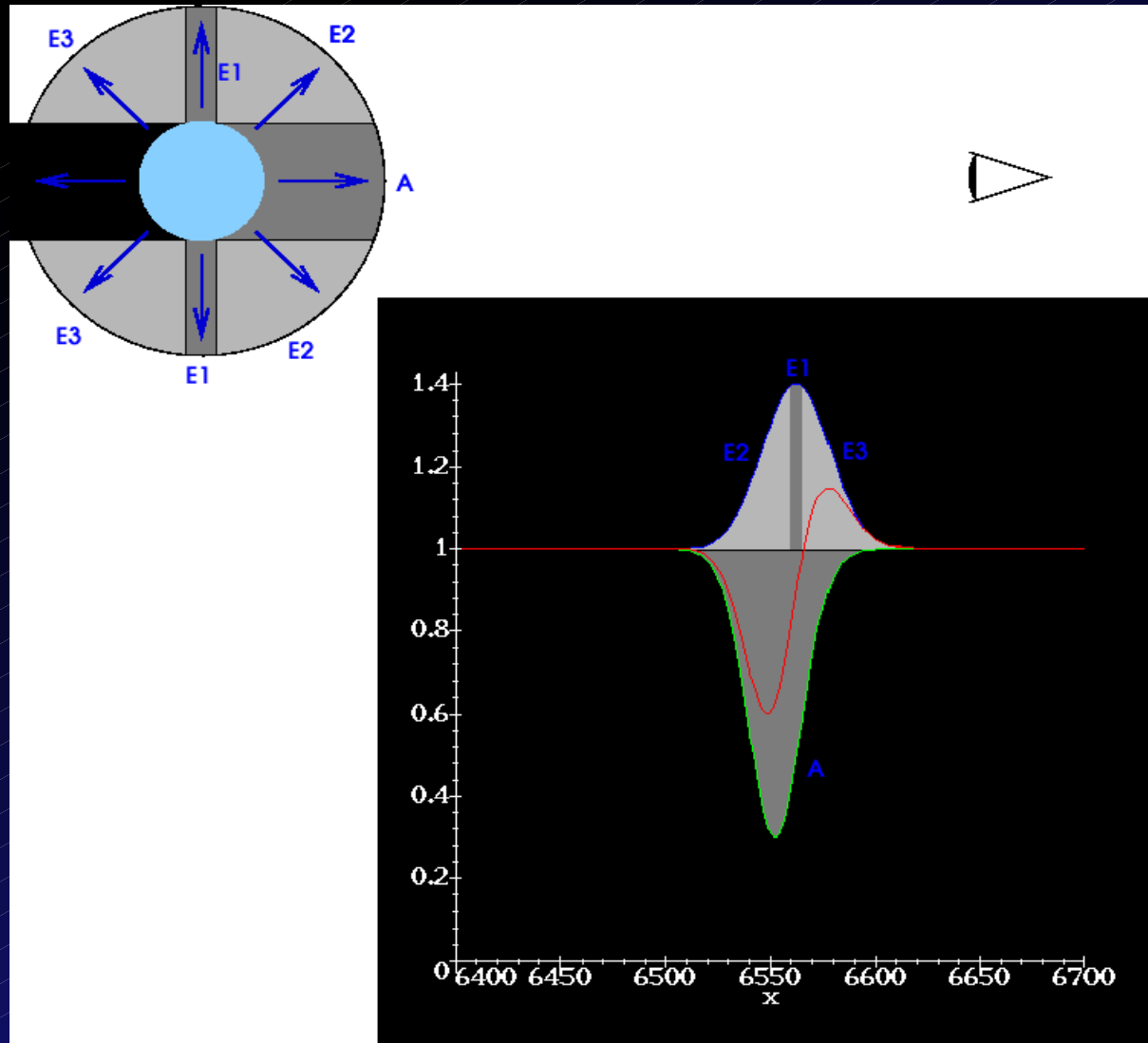


Target-slit error: Centroid varies depending on position of object with respect to slit due to guiding error or *movement between telescope and slit*

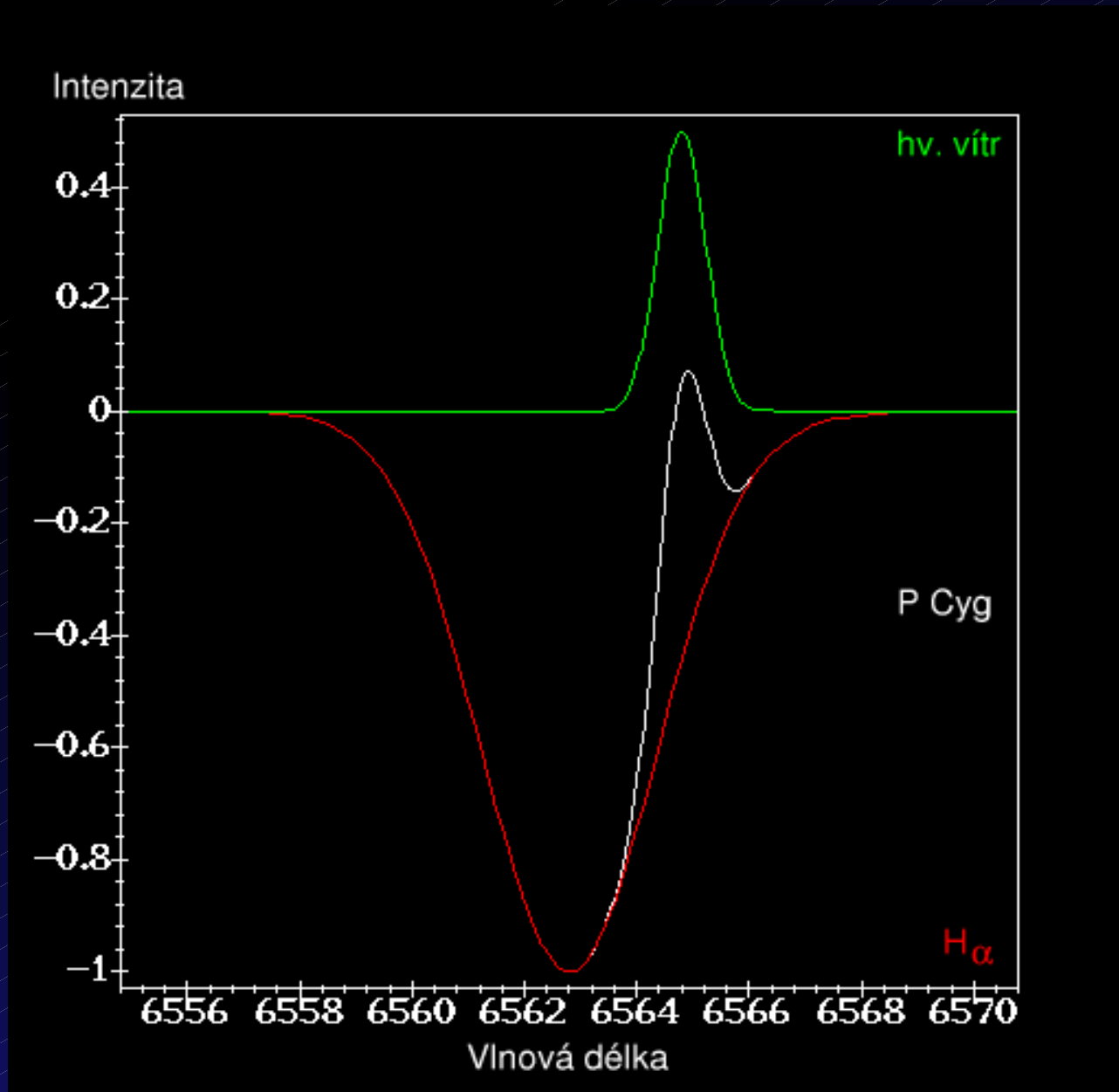


Slit-detector error: Centroid varies due to *movement between slit and detector*

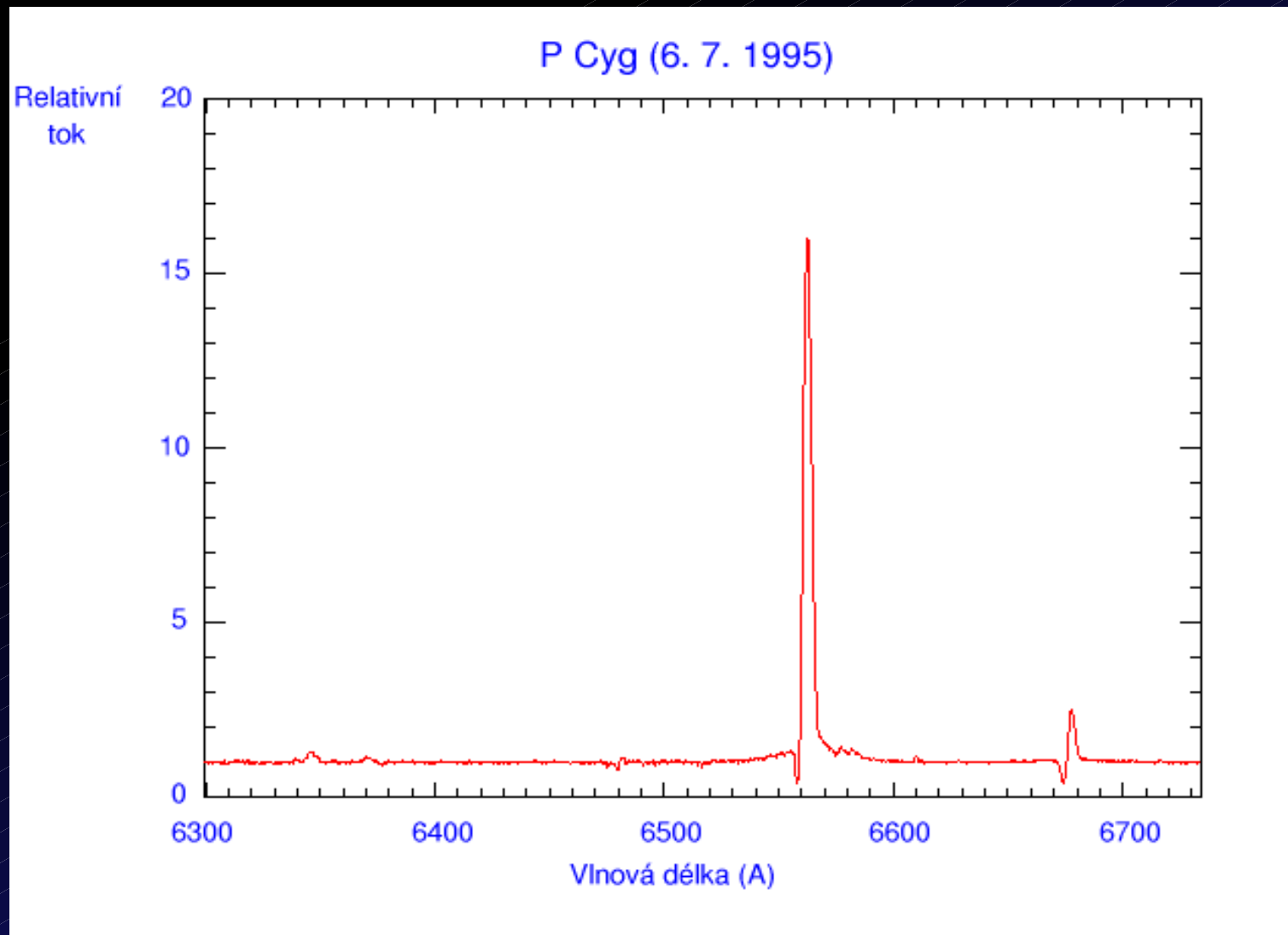
Vznik P Cyg profilu



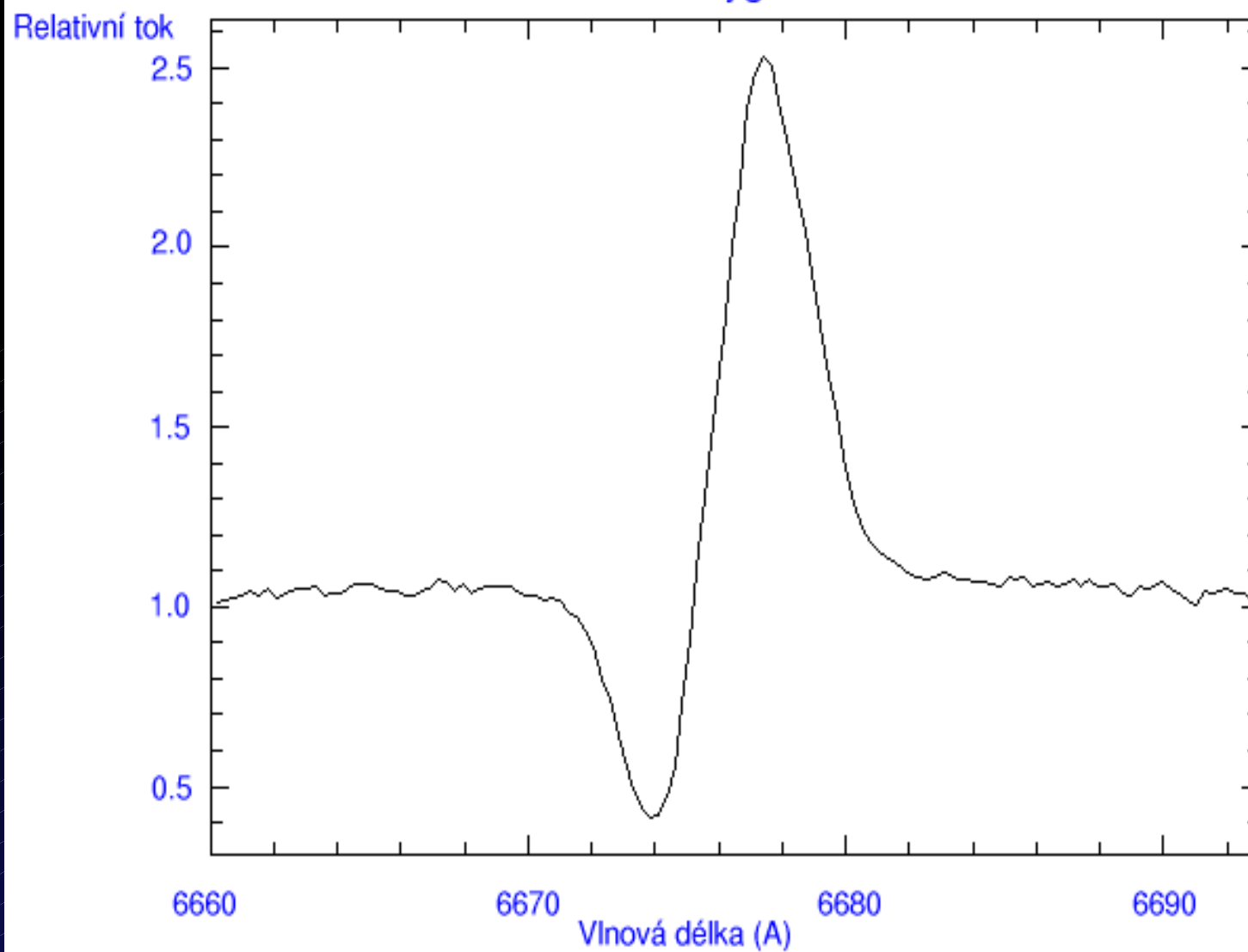
Vznik P Cyg profilu



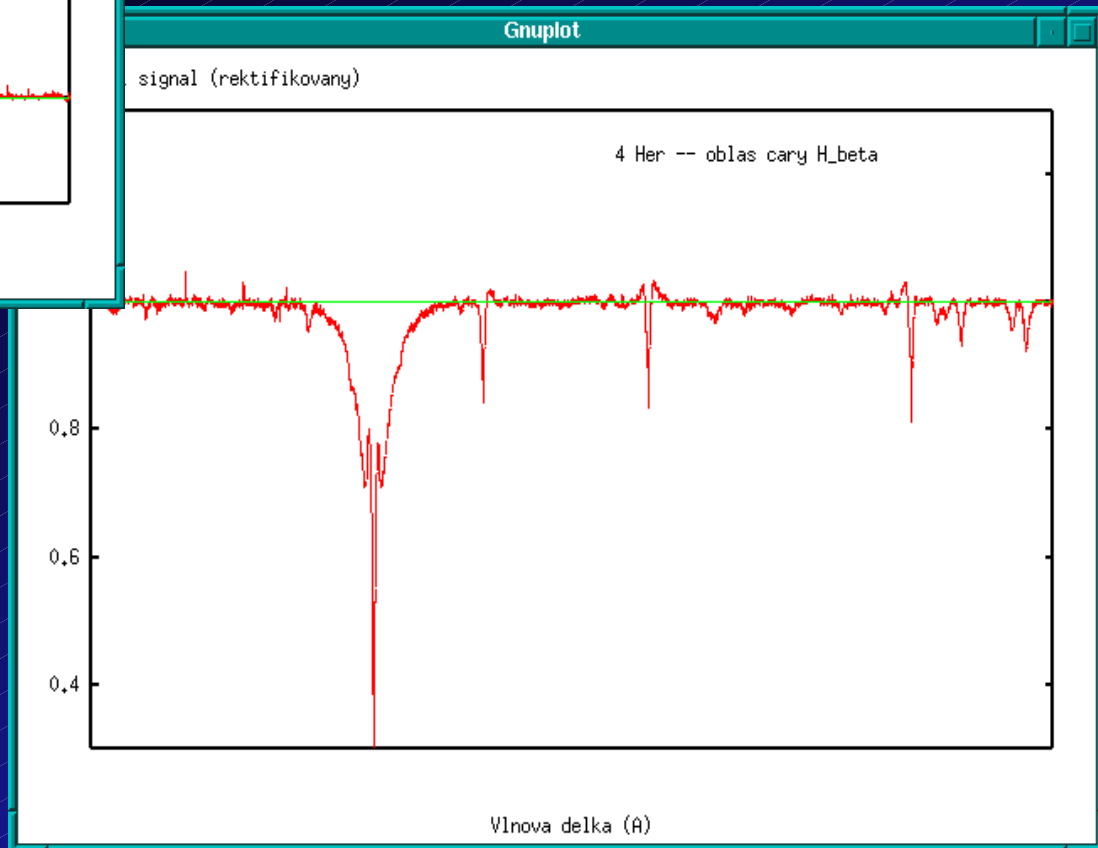
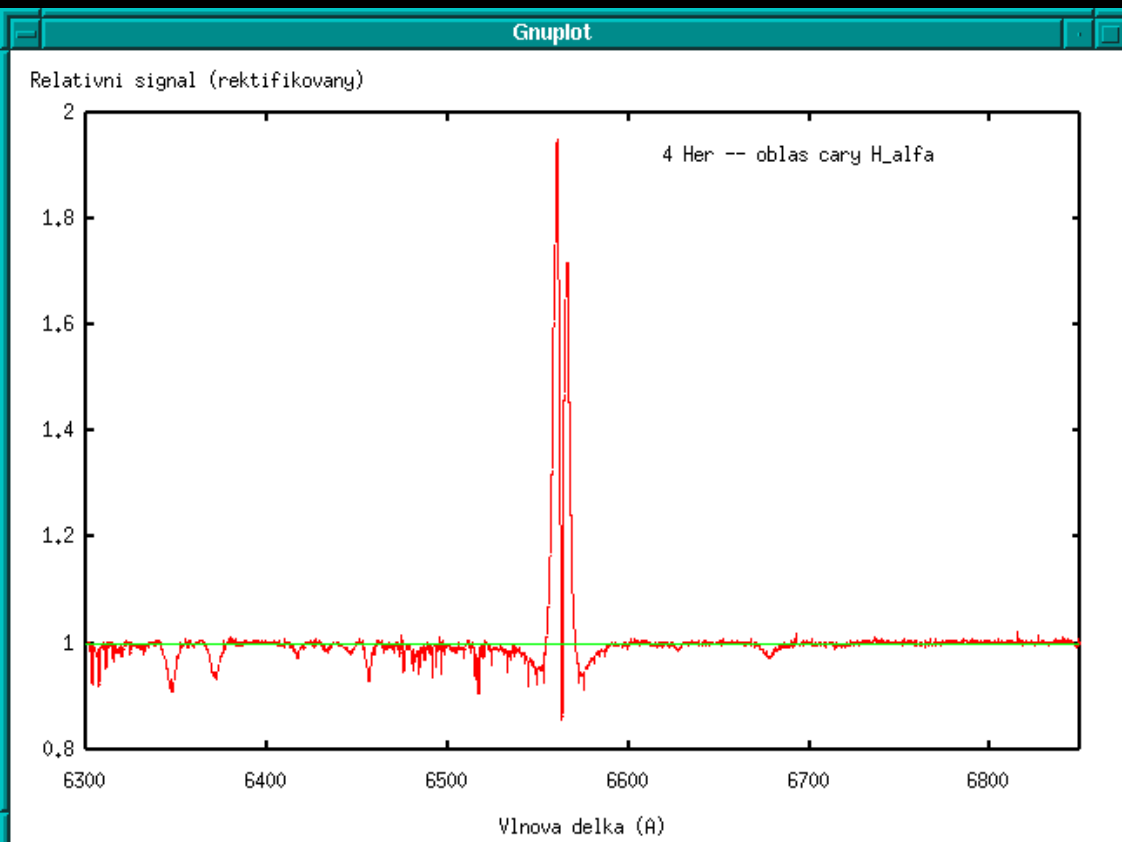
P Cyg



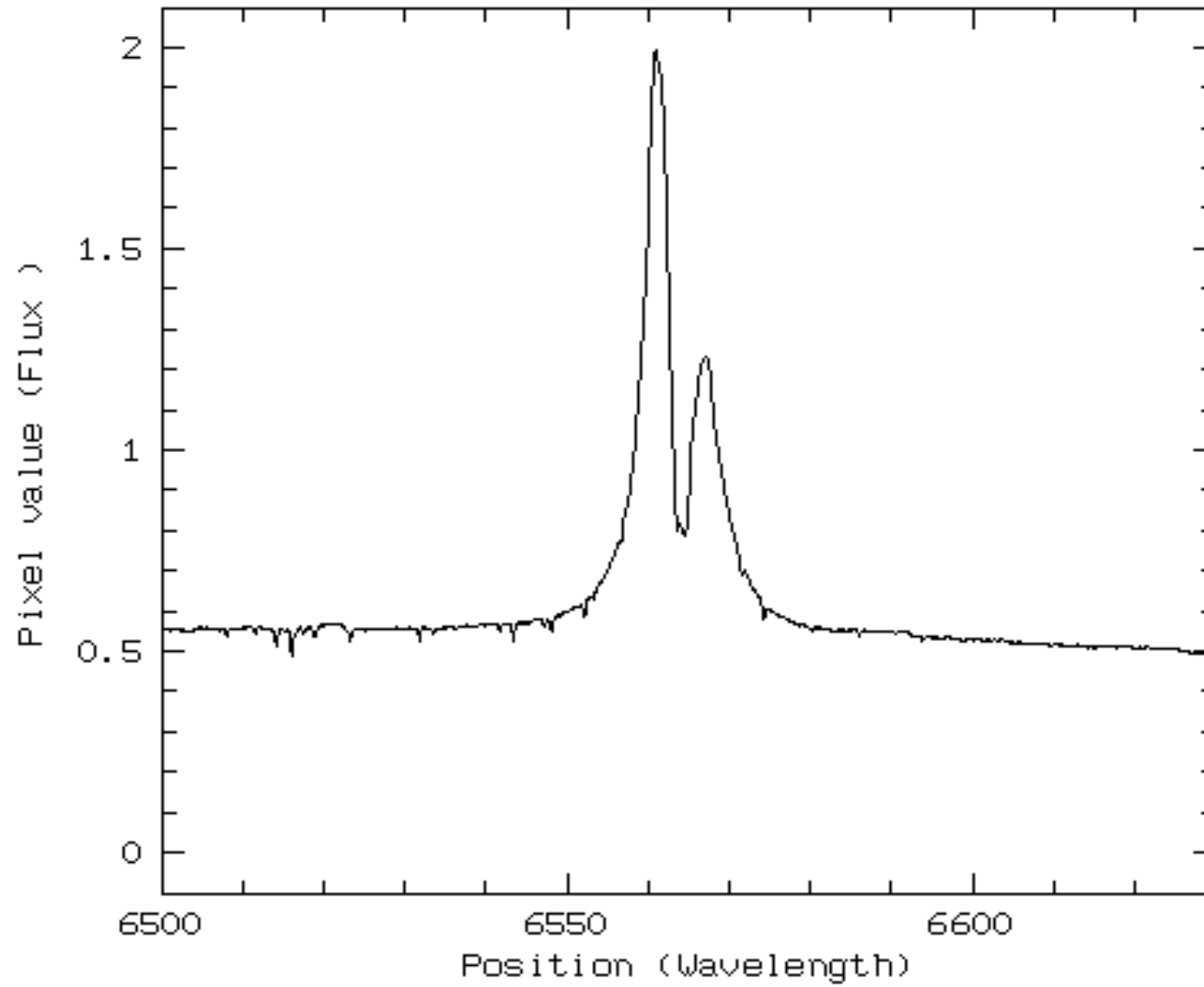
P Cyg



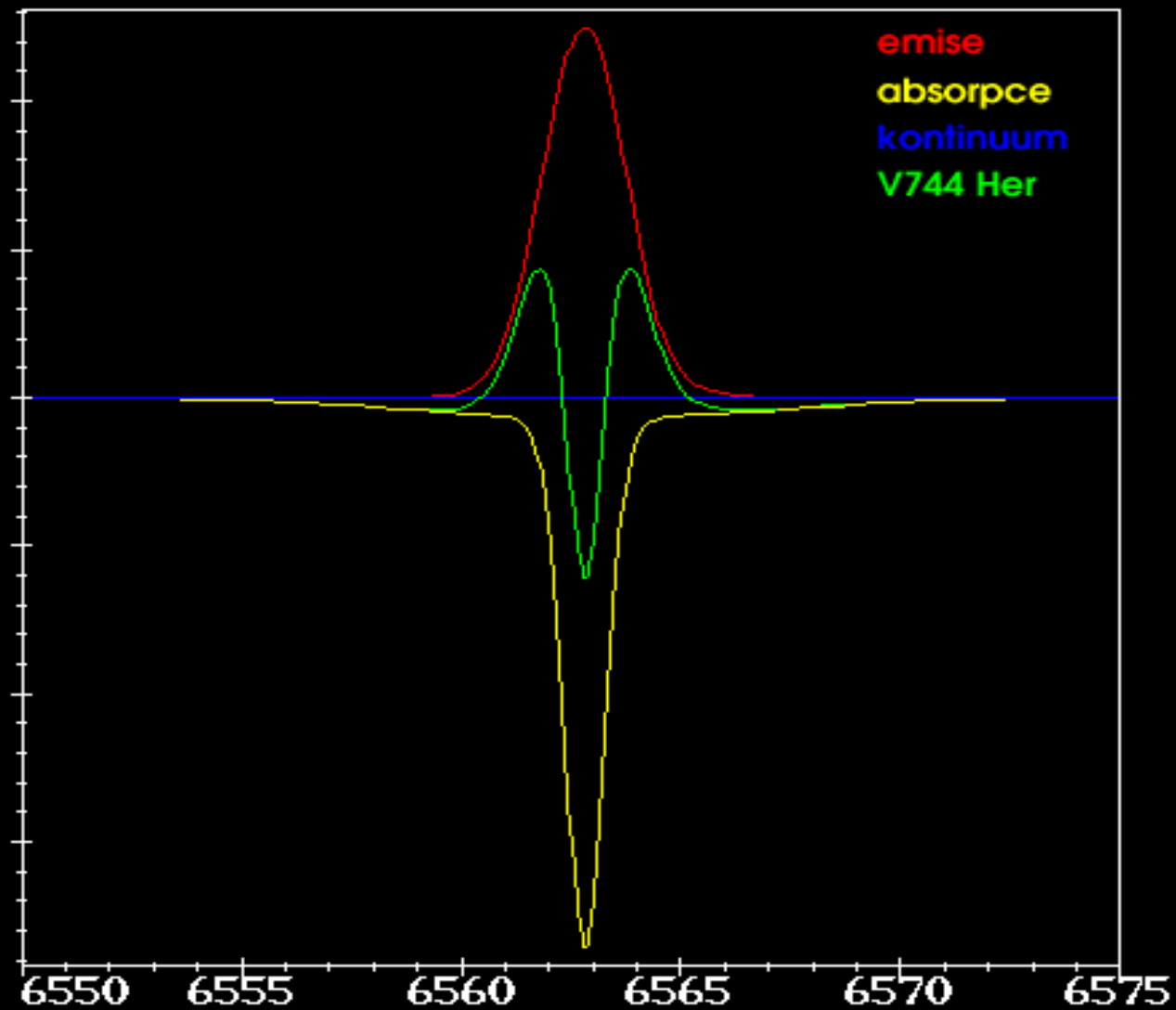
4 Her



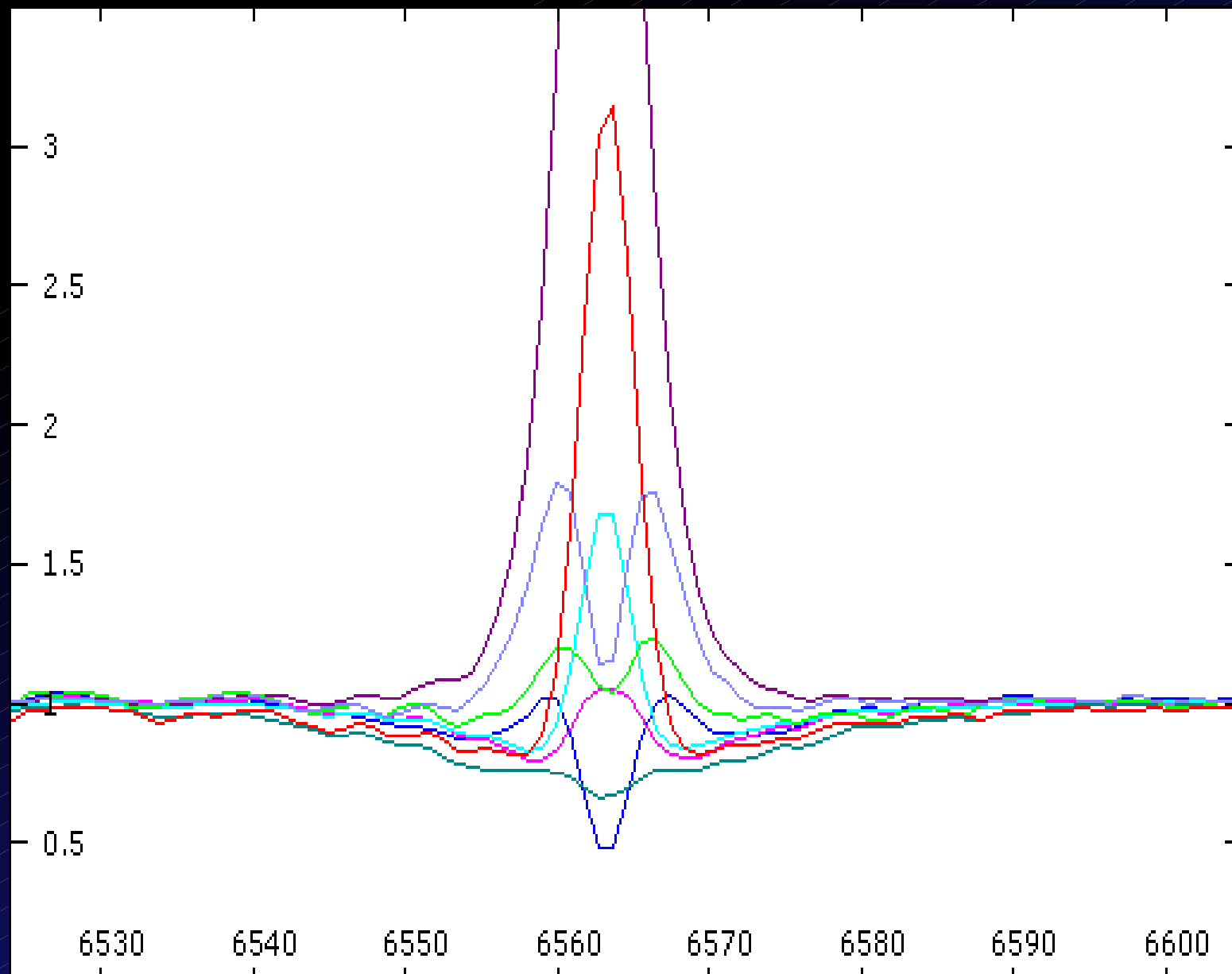
Zet Tau



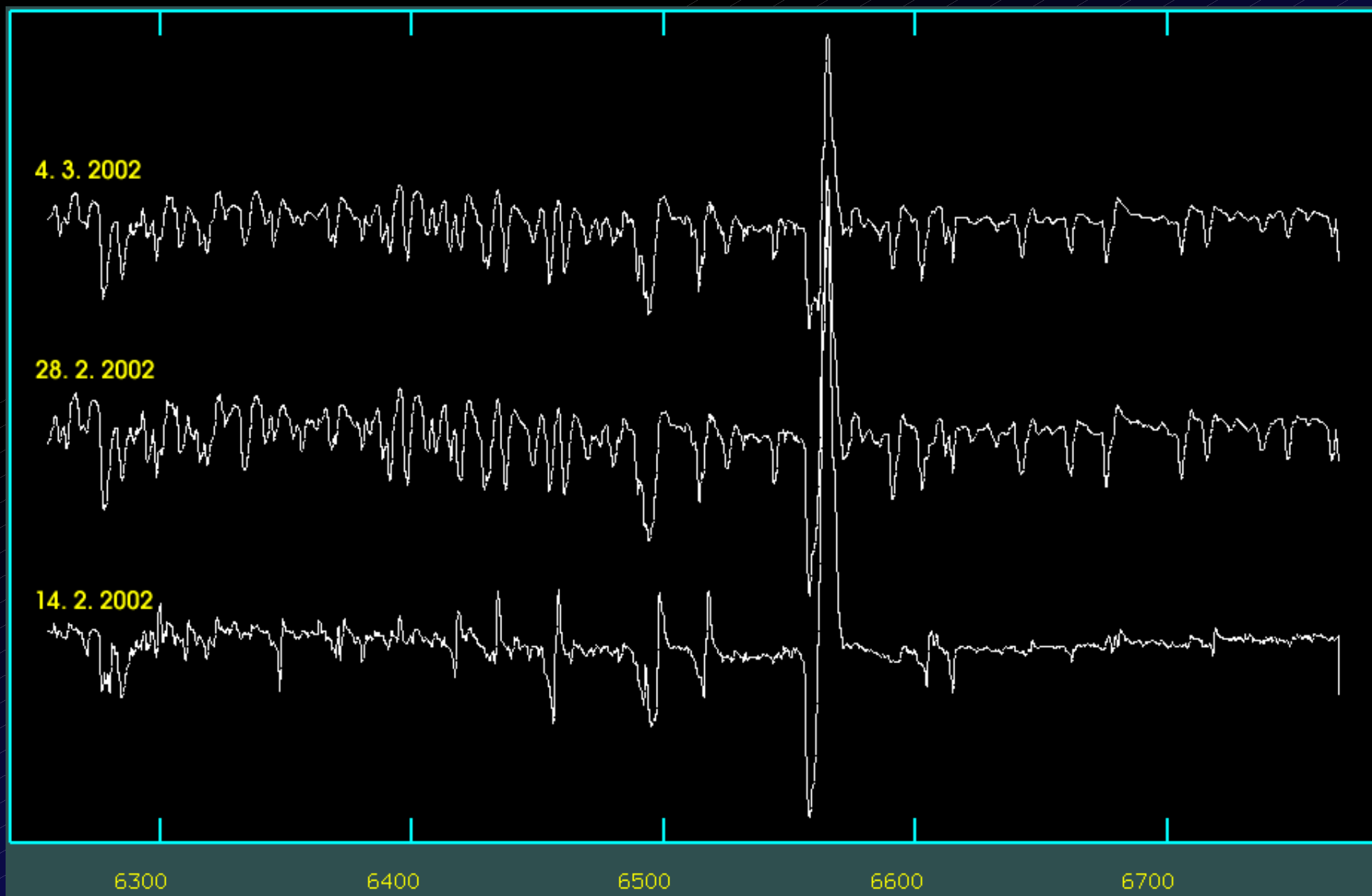
Shell čára u Be hvězd



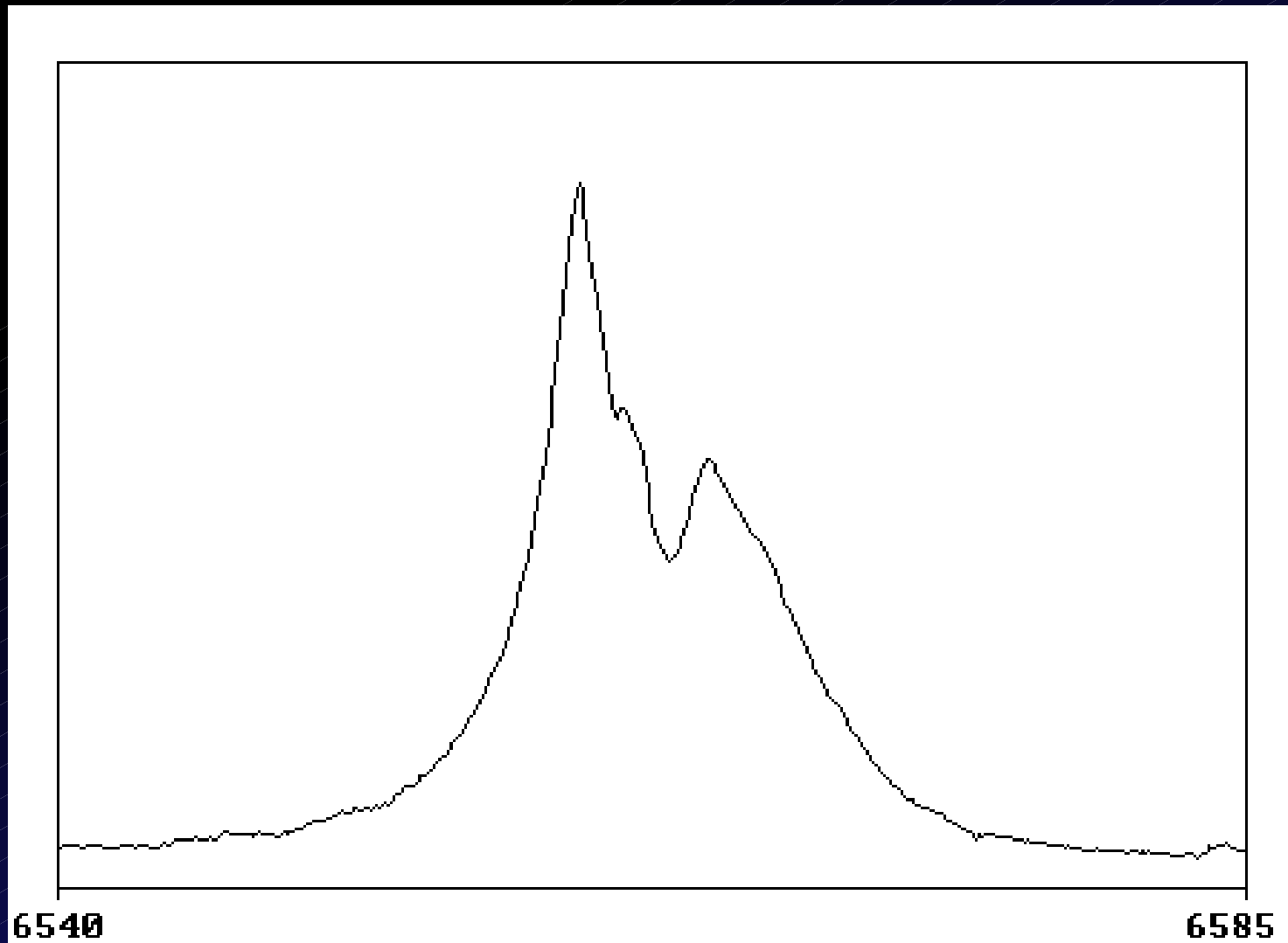
Be profily (Buill)



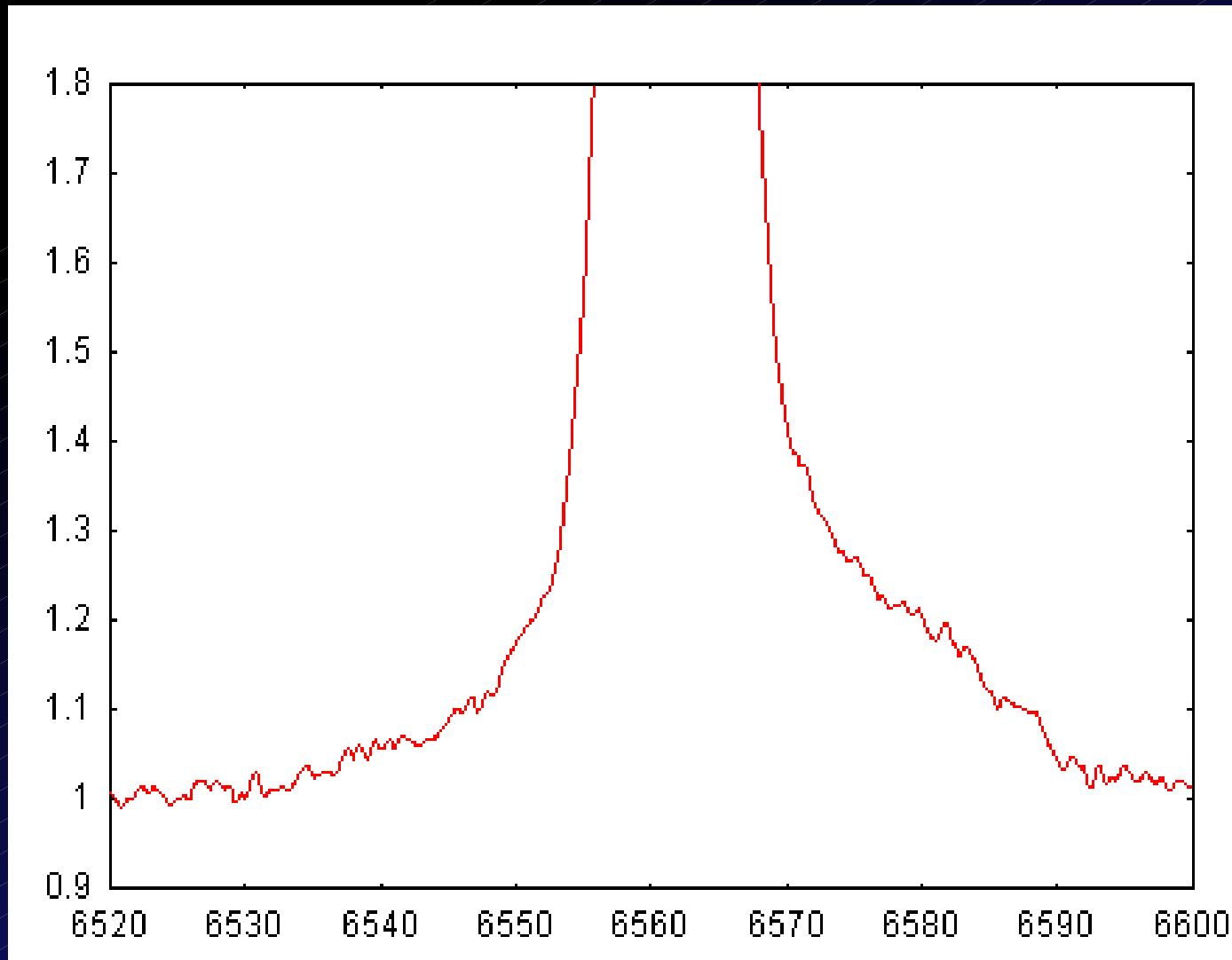
V 838 Mon



Zet Tau (Buill)

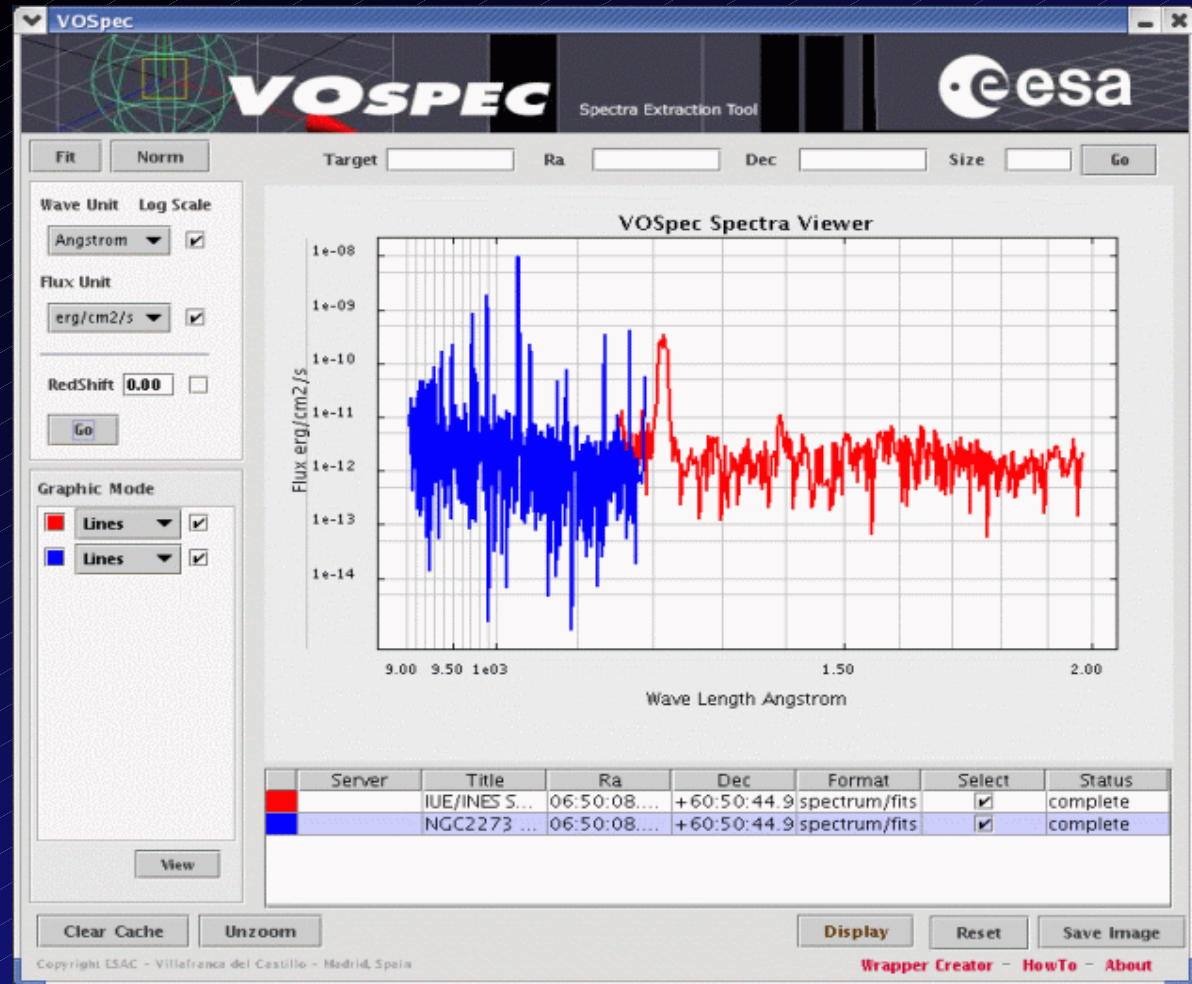


HD 206773 (Buill)

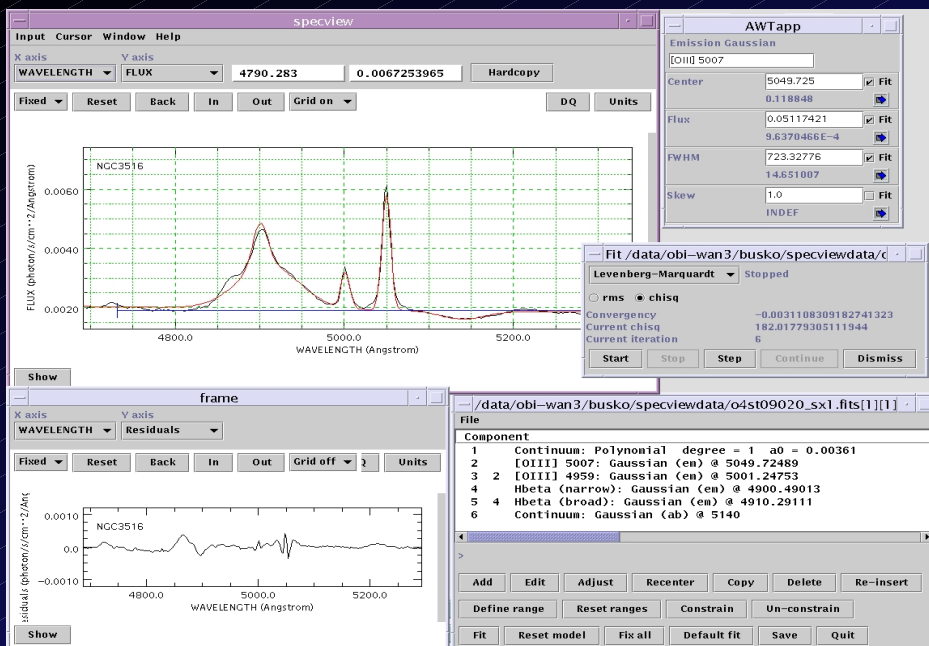
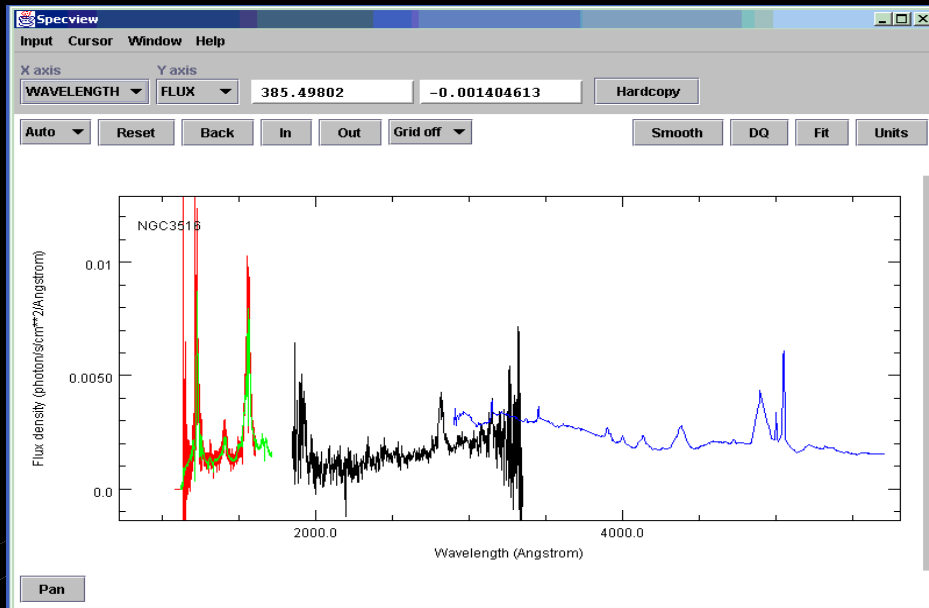


VOSpec (ESAC)

- Very simple
- Polynomial fits
- No RV measurement
- No complex operations
- In VizieR now
- Can work with SLAP!
- Theoretical VO supported
- Rapid development :-)

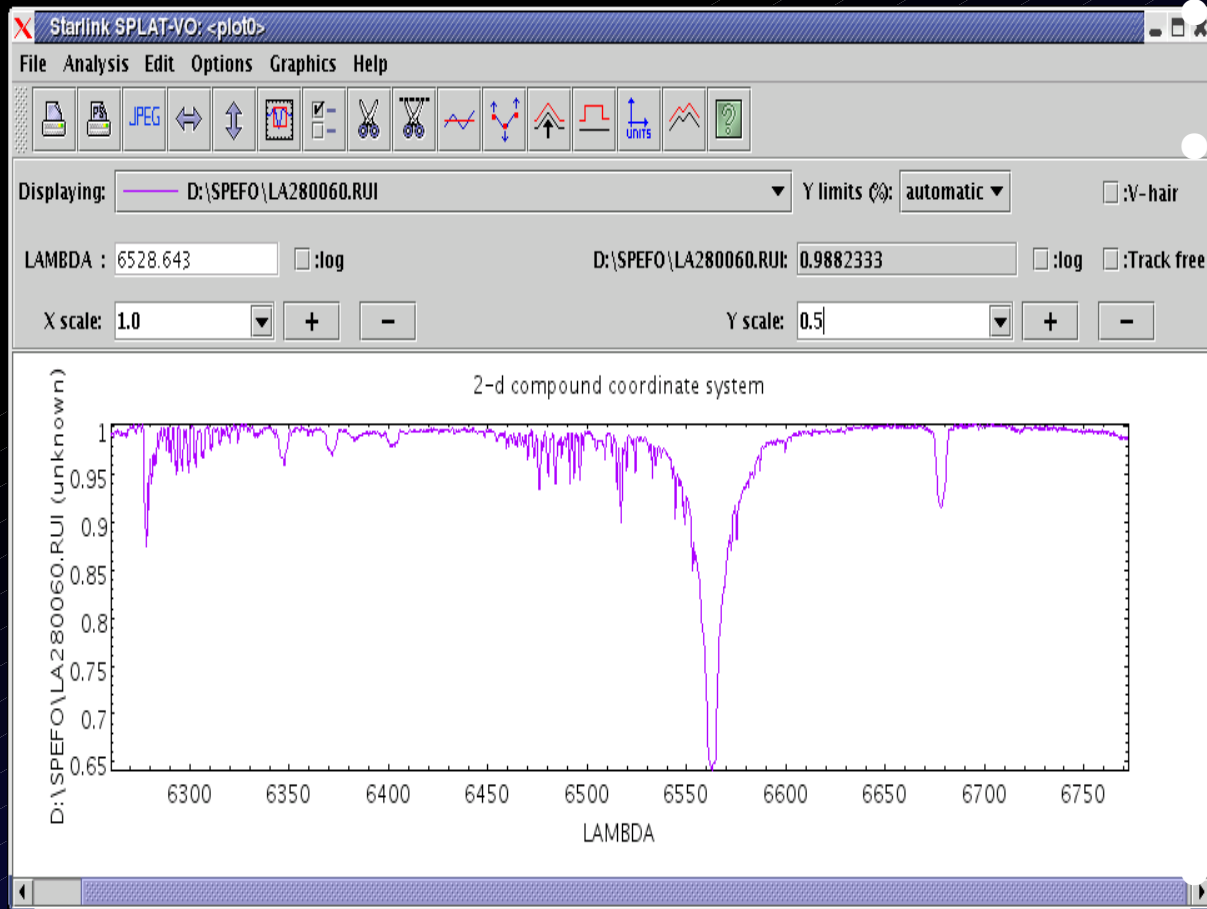


SpecView (STScI)



- Fitting profiles from models
- Simple polynomials
- Analysis strong (dereddening, CLOUDY)
- Supported !!
- Not good for IRAF WCS (1D FITS)
- BinTables+Extension

SPLAT-VO



Custom line list

Development not justified ?

Most advanced for stellar astronomy

JCMT now

Plastic

Reads 1D FITS...